

THE WISE DETECTION OF AN INFRARED ECHO IN TIDAL DISRUPTION EVENT ASASSN-14li

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ABSTRACT

We report the detection of a significant infrared variability of the nearest tidal disruption event (TDE) ASASSN-14li using *Wide-field Infrared Survey Explorer* and newly released *Near-Earth Object WISE Reactivation* data. In comparison with the quiescent state, the infrared flux is brightened by 0.12 and 0.16 mag in the W1 (3.4 μ m) and W2 (4.6 μ m) bands at 36 days after the optical discovery (or ~110 days after the peak disruption date). The flux excess is still detectable ~170 days later. Assuming that the flare-like infrared emission is from the dust around the black hole, its blackbody temperature is estimated to be ~2.1 \times 10³ K, slightly higher than the dust sublimation temperature, indicating that the dust is likely located close to the dust sublimation radius. The equilibrium between the heating and radiation of the dust claims a bolometric luminosity of ~10⁴³–10⁴⁵ erg s⁻¹, comparable with the observed peak luminosity. This result has for the first time confirmed the detection of infrared emission from the dust echoes of TDEs.

Key words: galaxies: active – galaxies: individual (ASASSN-14li) – galaxies: nuclei

1. INTRODUCTION

Evidence has mounted that supermassive black holes with masses of $10^{6-10}\,M_\odot$ are present in most (possibly all) galaxies with massive bulges, and the BH mass correlates tightly with various bulge properties (see Kormendy & Ho 2013 for a review). When a star orbits too close to the black hole, the tidal force exceeds the star's self-gravity and hence the star gets tidally disrupted. In these so-called tidal disruption events (TDEs), roughly half of the mass of the star may be ejected while the rest of the stellar material is accreted on to the black hole, producing a luminous flare of electromagnetic radiation (Rees 1988). Their emission peaks in the UV or soft X-rays, with a characteristic $t^{-5/3}$ decline on the timescale of months to years (Evans & Kochanek 1989; Phinney 1989). A number of TDE candidates were discovered by flares in the X-ray, UV, optical and by peculiar superstrong coronal lines in the optical spectra (see the review by Komossa 2015). Theoretically, the event rate is estimated to be 10^{-4} – 10^{-5} galaxy⁻¹ yr⁻¹ (Rees 1988; Magorrian & Tremaine 1999) and might be highest in nucleated dwarf galaxies (Wang & Merritt 2004; Stone & Metzger 2016).

One of the most promising scopes of TDEs is to probe the circum-nuclear interstellar medium of inactive galaxies, i.e., by detecting the echoes of the gas photoionized by the flare (e.g., Komossa et al. 2008; Wang et al. 2011, 2012; Yang et al. 2013; Arcavi et al. 2014). If the medium is dusty, most of the UV–optical photons will be absorbed by dust grains in the vicinity of the BH and re-radiated in the infrared (IR). Lu et al. (2016) show that the dust emission peaks at 3–10 μ m and has typical luminosities between 10^{42} and 10^{43} erg s⁻¹, which is detectable by current-generation telescopes. Nevertheless, none of such detections has been reported as of today. Taking advantage of data from the *Wide-field Infrared Survey Explorer (WISE*; Wright et al. 2010) and newly released data from the Near-Earth Object *WISE* Reactivation mission (*NEOWISE*-R; Mainzer et al. 2014), we have studied the IR light curve of the nearby TDE ASASSN-14li and detected significant

variability that is probably the first confirmation of dust echoes from TDEs.

ASASSN-14li was originally discovered on 2014 November 22 (modified Julian day MJD 56,983) by the All-Sky Automated Survey for Supernovae (ASASSN; Jose et al. 2014). The discovery was triggered by an optical flare coincident with the center of the galaxy PGC 043234 within 0."04, prompting it to be a potential TDE flare. Later, the Swift X-ray observations established a new source at the same location. This galaxy is located at z = 0.0206, making ASASSN-14li the closest TDE discovered in over 10 years. The fitting of the light curves using a Monte Carlo approach found that the accretion rate first exceeded the Eddington ratio in 2014 July, and the peak luminosity occurred around September 6, which is two and a half months before the discovery (Alexander et al. 2016). Since its discovery, ASASSN-14li has attracted much attention and multiwavelength follow-up monitorings were carried out immediately. Light curves and spectra in the optical and near-UV (Holoien et al. 2016; Cenko et al. 2016), soft X-rays (Miller et al. 2015), and radio (Alexander et al. 2016; van Velzen et al. 2016a) have all been obtained.

However, we still lack the knowledge of the infrared properties of ASASSN-14li, which is the reprocessed emission from dust located around the BH. Fortunately, the *NEOWISE*-R data that was released on 2016 March 23^4 offers us a good opportunity to probe the IR emission associated with ASASSN-14li. The *WISE* has performed a full-sky imaging survey in four broad mid-infrared bandpasses centered at 3.4, 4.6, 12, and $22 \,\mu\text{m}$ (labeled W1–W4) from 2010 February to August. The solid hydrogen cryogen used to cool the W3 and W4 instrumentation was depleted later and it was placed in hibernation in 2011 February. *WISE* was reactivated and renamed *NEOWISE*-R since 2013 October, using only W1 and W2, to hunt for asteroids that could pose as an impact hazard to the Earth.

⁴ http://wise2.ipac.caltech.edu/docs/release/neowise/

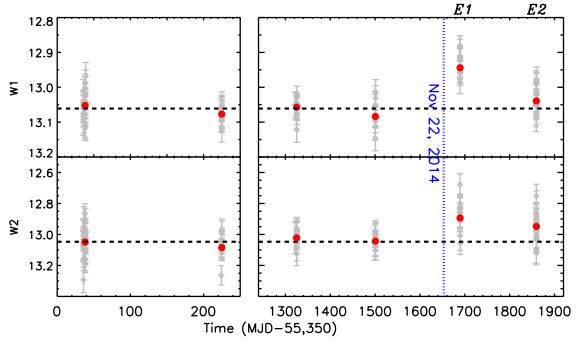


Figure 1. WISE infrared light curves of ASASSN-14li in the W1 and W2 band. Left panel: AllWISE data; right: NEOWISE-R data. The red solid circles denote the median value in each epoch while the single exposures are in gray. The uncertainties of these median magnitudes in each epoch are \lesssim 0.01 in W1 and \lesssim 0.03 in W2. The black dashed line is the median magnitude of the first four epochs, representing the quiescent state. The blue dotted line denotes the date when ASASSN-14li was discovered in ASASSN survey.

Table 1
WISE Data

Epoch (1)	MJD (2)	W1 (3)	W2 (4)	F1 (5)	F2 (6)	$T_{ m dust}$ (7)	L _{IR} (8)	L_{bol} (9)	$T_{\text{dust},2} $ (10)	$L_{\rm IR,2} $ (11)	$L_{\text{bol},2} $ (12)
Q	55365.546-56831.170	13.06	13.05								
E1	57018.881-57019.867	12.94	12.89	0.21	0.16	2091	41.6	45.6	904	41.3	42.9(44.0)
E2	57188.984–57189.968	13.04	12.95	0.04	0.10	588	41.0	45.0	444	40.8	41.9(43.0)

Note. Column (1): observational epochs. Column (2): modified Julian date. Column (3): median W1 magnitude. Column (4): median W2 magnitude. Column (5): enhanced flux in W1, in unit of mJy. Column (6): enhanced flux in W2, in unit of mJy. Column (7): dust temperature assuming blackbody radiation. Column (8): the logarithmic integrated dust IR luminosity. Column (9): bolometric luminosity required to heat the dust. Column (10): dust temperature considering the dust absorption efficiency $(Q_{abs} \propto \nu^2)$. Columns (11)–(12): same as columns (8)–(9) while with the case as described in column (10).

We assume a cosmology with $H_0 = 70 \,\mathrm{km \, s^{-1} \, Mpc^{-1}}$, $\Omega_m = 0.3$, and $\Omega_{\Lambda} = 0.7$. At a redshift of z = 0.0206, ASASSN-14li has a luminosity distance of 89.5 Mpc.

2. WISE LIGHT CURVE OF ASASSN-14li

ASASSN-14li is bright enough to have been detected in single exposures throughout the *WISE* mission. We have collected both cryogenic and post-cryogenic multiepoch photometry from the AllWISE⁵ and the most recent *NEO-WISE*-R data releases. The light curves of ASASSN-14li in the W1 and W2 bands have a baseline of 5 years (MJD 55,365-57,190) and are presented in Figure 1. The data points are mainly distributed over six epochs at intervals of about six months with the latest data taken in 2015 June. There is a large gap ($\sim 10^3$ days) between epoch 2 and epoch 3 due to the hibernation of *WISE* between 2011 February and 2013 October. Each epoch has 10-24 single-exposure observations, provided by the unique-designed observational cadence of *WISE*, which

is well suited for studying rapid variability (e.g., Jiang et al. 2012). We do not include W3 and W4 data that are restricted to the cryogenic period of the *WISE* mission.

First, no reliable variability is detected within each epoch. Second, we have not found any obvious flux variations between the first four epochs (MJD-55350 < 1600). However, we can see a significant flux increase during the fifth epoch (MJD \approx 57,019) that is 36 days after the TDE discovery time (Jose et al. 2014). We take the former four epochs as the quiescent state of the IR emission from the host in light of the fact that the magnitudes remain almost constant. For convenience, we denote the quiescent state as "Q" and the two epochs after the discovery of the TDE as "E1" and "E2" (see Figure 1). The median magnitudes in the W1 and the W2 bands at E1 are increased by 0.12 and 0.16 mag, respectively. The sudden flare-like characteristic, coordinated with the optical/UV/X-ray flare, makes it probably the first detection of TDE in the infrared. The flux excess at E2 becomes weaker, vet still detectable in comparison with the quiescent state (see Table 1).

⁵ http://wise2.ipac.caltech.edu/docs/release/allwise/

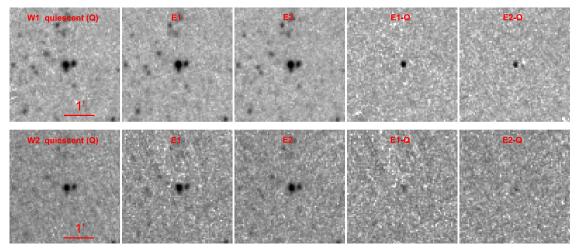


Figure 2. W1 and W2 images in different epochs. From left to right: the coadded image of the quiescent state (Q), E1, E2, the residual image of E1-Q; the residual image of E2-Q.

To exclude any possible systematic biases that make a spurious variability signal, we have examined the photometric zeropoint offsets between the different frames by employing nearby stars as "standards." The zeropoint errors are found to be less than 0.02 mag, which is much lower than the variation between E1 and Q (see also Jiang et al. 2012). We have also constructed coadded images in the three states (Q, E1, E2), respectively. After subtracting the quiescent state, clear flux excesses are shown in the W1 images, at both E1 and E2; the excesses in the W2 bands are also visible, albeit weaker due to the noisy background (see Figure 2).

3. DISCUSSIONS AND CONCLUSIONS

The increased IR emission associated with ASASSN-14li can be naturally attributed to the dust-reprocessed light of the optical–UV photons produced by the TDE. Because the dust is located at some distance from the BH, the IR variability is delayed with respect to the UV–optical emission by the light-travel time from the BH to the dust. As a rough estimation, the time delay after the maximum luminosity ($\Delta t \sim 110$ days) leads to a distance of ~ 0.09 pc, or $\sim 3 \times 10^5$ Schwarzschild radius (r_g) given a black hole mass of $10^{6.5} M_{\odot}$.

Having been monitored in multiple wavelengths (soft X-rays, UV, optical, and radio) after its discovery, ASASSN-14li possesses an exceptionally rich data set up to now. We have collected all available data reported in the literature, which are quasi-simultaneous with the *WISE* observations. With the newly added *WISE* IR photometry, we have constructed a host-subtracted spectral energy distribution (SED) with the most complete wavelength coverage among all known TDEs. Figure 3 shows the broadband SED of ASASSN-14li at the time of E1.

A double blackbody fitting to the UV-optical–IR SED at E1 yields a dust temperature ($T_{\rm dust}$) of 2091 \pm 172 K, contributing an integrated IR luminosity of 2.5 \times 10⁴¹ erg s⁻¹. The X-ray (\sim 7.7 \times 10⁵ K) and UV-optical (\sim 2.6 \times 10⁴ K) temperatures were nearly constant over a period of several months (Holoien et al. 2016; Cenko et al. 2016; van Velzen et al. 2016a), although their luminosities were declining. The dust is quite hot, with a temperature even slightly higher than the typical sublimation temperature ($T_{\rm sub}$; e.g., 1400 K for silicate, 1800 K for graphite). The peak luminosity (\sim 10⁴⁴ erg s⁻¹) of

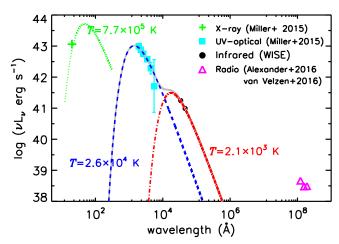


Figure 3. Broadband SED of ASASSN-14li at E1 (MJD \approx 57019; see Figure 1). The X-ray, UV, and optical data are adapted from Figure 1 of Miller et al. (2015); the radio data are drawn from Table S1 of van Velzen et al. (2016a) and Table 1 of Alexander et al. (2016). The infrared data are converted from our measurement. The blue dashed line and red dotted–dashed line are the two-blackbody fitting to the UV–optical–infrared data, which have been already subtracted from the host emission before the TDE.

ASASSN-14li determines a characteristic dust sublimation radius ($R_{\rm sub}$) of ~0.04 pc⁶, comparable to the estimation dust location from the light-travel distance. After being heated by the initial continuum pulse, only grains outside the sublimation radius at peak luminosity can survive. As a result, the high $T_{\rm dust}$ can be reasonably understood as the IR emission we observed originates from the dust around the $R_{\rm sub}$. The same approach gives a $T_{\rm dust}$ of 588 K at E2, which is at ~210 days after its discovery.

The dust mass associated with the IR emission can be also derived. We assume that the dust grains follow an MRN size distribution (Mathis et al. 1977; see also Draine & Lee 1984) as $dn/da \propto a^{-3.5}$ with $a_{\rm min}=0.005\,\mu{\rm m}$, $a_{\rm max}=0.3\,\mu{\rm m}$ and an average density of $\rho=2.5\,{\rm g\,cm^{-3}}$. The total dust mass is estimated to be $1.2\times10^{27}\,{\rm g}$, that is, $\sim\!10^{-6}\,M_{\odot}$. Assuming a dust to gas ratio of 0.01, the amount of corresponding gas is $\sim\!10^{-4}\,M_{\odot}$.

⁶ The formula we used to calculate is that of footnote 1 in Namekata & Umemura (2016).

For a dust grain with a radius a at a distance of R from the heating source, its equilibrium temperature T is determined by the balance between the radiative heating by the absorption of UV–optical photons (proportional to the angular average of radiation intensity) and the thermal re-emission in the IR. Then the bolometric luminosity ($L_{\rm bol}$) can be calculated by the equation below:

$$\frac{L_{\text{bol}}}{4\pi R^2} \pi a^2 = 4\pi a^2 \sigma T^4. \tag{1}$$

We adopt $R=c\Delta t$, in which $\Delta t=110$ days and c is the light speed, then $L_{\rm bol}=4.4\times10^{45}\,{\rm erg\,s^{-1}}$, which is one order of magnitude higher than the observed peak luminosity. The blackbody assumption for the dust emission could be oversimplistic. In a more practical calculation taking into consideration the absorption efficiency ($Q_{\rm abs}\propto \nu^{\beta}$; see formula 8.16 in Kruegel 2003), the fitted $T_{\rm dust}=904\,{\rm K}$ (if we adopt $\beta=2$) and the required $L_{\rm bol}$ can be estimated as follows:

$$\frac{L_{\text{bol}}}{16\pi^2 R^2} \simeq 1.47 \times 10^{-6} aT^6. \tag{2}$$

The result is very sensitive to the grain size distribution: $L_{\rm bol}=7.8\times 10^{42}\,{\rm erg\,s^{-1}}$ assuming an MRN grain size distribution; if we adopt a uniform grain size of $0.1\,\mu{\rm m},\,L_{\rm bol}$ turns out to be $1.0\times 10^{44}\,{\rm erg\,s^{-1}}$ (see relevant results in Table 1). In this case, the estimated dust mass is ~ 10 times higher ($\sim 10^{-5}\,M_{\odot}$) than that assuming the simple blackbody radiation.

During our calculation above, we assumed that the medium is optically thin, which is consistent with the fact that the infrared luminosity ($\sim 10^{41} \, {\rm erg \ s^{-1}}$; see Table 1) is only a small fraction ($\sim 1\%$) of the predicted luminosity if the dust intercept all of the available flux ($\sim 10^{43} \, {\rm erg \ s^{-1}}$; Lu et al. 2016). The derived $L_{\rm bol}$ above ($10^{43} - 10^{45} \, {\rm erg \ s^{-1}}$) is basically comparable to the observed peak luminosity of ASASSN-14li (Holoien et al. 2016), providing a substantial evidence of the IR echo associated with the TDE. Moreover, the temperature at E2 (444 K) is roughly consistent with the value predicted by the $t^{0.5}$ decaying trend (494 K) under the equilibrium assumption. Previous works show that [O III] ($L_{\rm O \ III} \sim 10^{39} \, {\rm erg \ s^{-1}}$) and

radio emission do exist in PGC 043234 before the TDE, suggesting that the galaxy may host a low-luminosity AGN. However, the lack of X-ray emission and the W1-W2 color of WISE imply that the nuclear activity is rather weak (Holoien et al. 2016). The gas and dust we found from the dust echo of TDE indicate that they likely are located quite close to the black hole prior to the disruption. It is interesting to investigate why the gas is there but is not accreted efficiently by the black hole. Given the gas mass estimated above, a back of the envelope calculation suggests that Bondi accretion would yield an accretion rate of only $\sim 10^{-7} M_{\odot} \text{ yr}^{-1}$, comparable to the amount required to feed the black hole as revealed by predisruption radio emission ($\sim 10^{-4}$ Eddington ratio). The other possible origin of the dust is from the unbound debris from the disrupted star itself either in the form of a wind, which is also thought to be responsible for the radio emission (Alexander et al. 2016), or in the form of the unbound tail (Guillochon et al. 2016). Assuming a velocity of 0.3c (Alexander et al. 2016), the distance that the unbound debris can travel is ~ 0.04 pc, close to the R_{sub} .

This work demonstrates that the WISE and NEOWISE-R data afford us an excellent opportunity to study the dust properties in the vicinity of black holes residing in dormant (or weakly active) galaxies, by taking advantage of the dust echo light of TDEs (see also van Velzen et al. 2016b; Dou et al. 2016). Besides the studies of known TDEs, we can also attempt to discover new TDEs at the infrared band by making full use of the released WISE and NEOWISE-R data, which have an allsky covering, a wide time span (~6 years up until now), and a unique survey cadence. Our preliminary try shows that this method is quite efficient and some TDE candidates with similar IR echo signals can be chosen in a blind search. In the future, we hope to characterize the nuclear environments of inactive galaxies statistically by comparing the luminosity functions of TDEs found in the infrared and other bands (e.g., optical, X-ray).

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