## Observation of the doubly strange $b$ baryon $\Omega_{b}^{-}$

V.M. Abazov ${ }^{36}$, B. Abbott $^{75}$, M. Abolins ${ }^{65}$, B.S. Acharya ${ }^{29}$, M. Adams ${ }^{51}$, T. Adams ${ }^{49}$, E. Aguilo ${ }^{6}$, M. Ahsan ${ }^{59}$, G.D. Alexeev ${ }^{36}$, G. Alkhazov ${ }^{40}$, A. Alton ${ }^{64, a}$, G. Alverson ${ }^{63}$, G.A. Alves ${ }^{2}$, M. Anastasoaie ${ }^{35}$, L.S. Ancu ${ }^{35}$, T. Andeen ${ }^{53}$, B. Andrieu ${ }^{17}$, M.S. Anzelc ${ }^{53}$, M. Aoki ${ }^{50}$, Y. Arnoud ${ }^{14}$, M. Arov ${ }^{60}$, M. Arthaud ${ }^{18}$, A. Askew ${ }^{49}$, B. Åsman ${ }^{41}$, A.C.S. Assis Jesus ${ }^{3}$, O. Atramentov ${ }^{49}$, C. Avila ${ }^{8}$, F. Badaud ${ }^{13}$, L. Bagby ${ }^{50}$, B. Baldin ${ }^{50}$, D.V. Bandurin ${ }^{59}$, P. Banerjee ${ }^{29}$, S. Banerjee ${ }^{29}$, E. Barberis ${ }^{63}$, A.-F. Barfuss ${ }^{15}$, P. Bargassa ${ }^{80}$, P. Baringer ${ }^{58}$, J. Barreto ${ }^{2}$, J.F. Bartlett ${ }^{50}$, U. Bassler ${ }^{18}$, D. Bauer ${ }^{43}$, S. Beale ${ }^{6}$, A. Bean ${ }^{58}$, M. Begalli ${ }^{3}$, M. Begel ${ }^{73}$, C. Belanger-Champagne ${ }^{41}$, L. Bellantoni ${ }^{50}$, A. Bellavance ${ }^{50}$, J.A. Benitez ${ }^{65}$, S.B. Beri ${ }^{27}$, G. Bernardi ${ }^{17}$, R. Bernhard ${ }^{23}$, I. Bertram ${ }^{42}$, M. Besançon ${ }^{18}$, R. Beuselinck ${ }^{43}$, V.A. Bezzubov ${ }^{39}$, P.C. Bhat ${ }^{50}$, V. Bhatnagar ${ }^{27}$, C. Biscarat ${ }^{20}$, G. Blazey ${ }^{52}$, F. Blekman ${ }^{43}$, S. Blessing ${ }^{49}$, K. Bloom ${ }^{67}$, A. Boehnlein ${ }^{50}$, D. Boline ${ }^{62}$, T.A. Bolton ${ }^{59}$, E.E. Boos $^{38}$, G. Borissov ${ }^{42}$, T. Bose ${ }^{77}$, A. Brandt ${ }^{78}$, R. Brock ${ }^{65}$, G. Brooijmans ${ }^{70}$, A. Bross ${ }^{50}$, D. Brown ${ }^{81}$, X.B. $\mathrm{Bu}^{7}$, N.J. Buchanan ${ }^{49}$, D. Buchholz ${ }^{53}$, M. Buehler ${ }^{81}$, V. Buescher ${ }^{22}$, V. Bunichev ${ }^{38}$, S. Burdin ${ }^{42, b}$, T.H. Burnett ${ }^{82}$, C.P. Buszello ${ }^{43}$, J.M. Butler ${ }^{62}$, P. Calfayan ${ }^{25}$, S. Calvet ${ }^{16}$, J. Cammin ${ }^{71}$, E. Carrera ${ }^{49}$, W. Carvalho ${ }^{3}$, B.C.K. Casey ${ }^{50}$, H. Castilla-Valdez ${ }^{33}$, S. Chakrabarti ${ }^{18}$, D. Chakraborty ${ }^{52}$, K.M. Chan ${ }^{55}$,
A. Chandra ${ }^{48}$, E. Cheu ${ }^{45}$, F. Chevallier ${ }^{14}$, D.K. Cho ${ }^{62}$, S. Choi ${ }^{32}$, B. Choudhary ${ }^{28}$, L. Christofek ${ }^{77}$, T. Christoudias ${ }^{43}$, S. Cihangir ${ }^{50}$, D. Claes ${ }^{67}$, J. Clutter ${ }^{58}$, M. Cooke ${ }^{50}$, W.E. Cooper ${ }^{50}$, M. Corcoran ${ }^{80}$,
F. Couderc ${ }^{18}$, M.-C. Cousinou ${ }^{15}$, S. Crépé-Renaudin ${ }^{14}$, V. Cuplov ${ }^{59}$, D. Cutts ${ }^{77}$, M. Ćwiok ${ }^{30}$, H. da Motta ${ }^{2}$, A. Das ${ }^{45}$, G. Davies ${ }^{43}$, K. De ${ }^{78}$, S.J. de Jong ${ }^{35}$, E. De La Cruz-Burelo ${ }^{33}$, C. De Oliveira Martins ${ }^{3}$, K. DeVaughan ${ }^{67}$, J.D. Degenhardt ${ }^{64}$, F. Déliot ${ }^{18}$, M. Demarteau ${ }^{50}$, R. Demina ${ }^{71}$, D. Denisov ${ }^{50}$, S.P. Denisov ${ }^{39}$, S. Desai ${ }^{50}$, H.T. Diehl ${ }^{50}$, M. Diesburg ${ }^{50}$, A. Dominguez ${ }^{67}$, H. Dong ${ }^{72}$, T. Dorland ${ }^{82}$, A. Dubey ${ }^{28}$, L.V. Dudko ${ }^{38}$, L. Duflot ${ }^{16}$,
S.R. Dugad ${ }^{29}$, D. Duggan ${ }^{49}$, A. Duperrin ${ }^{15}$, J. Dyer ${ }^{65}$, A. Dyshkant ${ }^{52}$, M. Eads ${ }^{67}$, D. Edmunds ${ }^{65}$, J. Ellison ${ }^{48}$, V.D. Elvira ${ }^{50}$, Y. Enari ${ }^{77}$, S. Eno ${ }^{61}$, P. Ermolov ${ }^{38, \ddagger}$, H. Evans ${ }^{54}$, A. Evdokimov ${ }^{73}$, V.N. Evdokimov ${ }^{39}$, A.V. Ferapontov ${ }^{59}$, T. Ferbel ${ }^{71}$, F. Fiedler ${ }^{24}$, F. Filthaut ${ }^{35}$, W. Fisher ${ }^{50}$, H.E. Fisk ${ }^{50}$, M. Fortner ${ }^{52}$, H. Fox ${ }^{42}$, S. Fu ${ }^{50}$, S. Fuess ${ }^{50}$, T. Gadfort ${ }^{70}$, C.F. Galea ${ }^{35}$, C. Garcia ${ }^{71}$, A. Garcia-Bellido ${ }^{71}$, V. Gavrilov ${ }^{37}$, P. Gay ${ }^{13}$, W. Geist ${ }^{19}$, W. Geng ${ }^{15,65}$, C.E. Gerber ${ }^{51}$, Y. Gershtein ${ }^{49}$, D. Gillberg ${ }^{6}$, G. Ginther ${ }^{71}$, N. Gollub ${ }^{41}$, B. Gómez ${ }^{8}$, A. Goussiou ${ }^{82}$, P.D. Grannis ${ }^{72}$, H. Greenlee ${ }^{50}$, Z.D. Greenwood ${ }^{60}$, E.M. Gregores ${ }^{4}$, G. Grenier ${ }^{20}$, Ph. Gris ${ }^{13}$, J.-F. Grivaz ${ }^{16}$, A. Grohsjean ${ }^{25}$, S. Grünendahl ${ }^{50}$, M.W. Grünewald ${ }^{30}$, F. Guo ${ }^{72}$, J. Guo ${ }^{72}$, G. Gutierrez ${ }^{50}$, P. Gutierrez ${ }^{75}$, A. Haas $^{70}$, N.J. Hadley ${ }^{61}$, P. Haefner ${ }^{25}$, S. Hagopian ${ }^{49}$, J. Haley ${ }^{68}$, I. Hall ${ }^{65}$, R.E. Hall ${ }^{47}$,
L. $\operatorname{Han}^{7}$, K. Harder ${ }^{44}$, A. Harel ${ }^{71}$, J.M. Hauptman ${ }^{57}$, J. Hays ${ }^{43}$, T. Hebbeker ${ }^{21}$, D. Hedin ${ }^{52}$, J.G. Hegeman ${ }^{34}$, A.P. Heinson ${ }^{48}$, U. Heintz ${ }^{62}$, C. Hensel ${ }^{22, c}$, K. Herner ${ }^{72}$, G. Hesketh ${ }^{63}$, M.D. Hildreth ${ }^{55}$, R. Hirosky ${ }^{81}$, J.D. Hobbs ${ }^{72}$, B. Hoeneisen ${ }^{12}$, H. Hoeth ${ }^{26}$, M. Hohlfeld ${ }^{22}$, S. Hossain ${ }^{75}$, P. Houben ${ }^{34}$, Y. Hu ${ }^{72}$, Z. Hubacek ${ }^{10}$, V. Hynek ${ }^{9}$, I. Iashvili ${ }^{69}$, R. Illingworth ${ }^{50}$, A.S. Ito ${ }^{50}$, S. Jabeen ${ }^{62}$, M. Jaffré ${ }^{16}$, S. Jain ${ }^{75}$, K. Jakobs ${ }^{23}$, C. Jarvis ${ }^{61}$, R. Jesik ${ }^{43}$, K. Johns ${ }^{45}$, C. Johnson ${ }^{70}$, M. Johnson ${ }^{50}$, D. Johnston ${ }^{67}$, A. Jonckheere ${ }^{50}$, P. Jonsson ${ }^{43}$, A. Juste ${ }^{50}$, E. Kajfasz ${ }^{15}$, J.M. Kalk ${ }^{60}$, D. Karmanov ${ }^{38}$, P.A. Kasper ${ }^{50}$, I. Katsanos ${ }^{70}$, D. Kau ${ }^{49}$, V. Kaushik ${ }^{78}$, R. Kehoe ${ }^{79}$, S. Kermiche ${ }^{15}$, N. Khalatyan ${ }^{50}$, A. Khanov ${ }^{76}$, A. Kharchilava ${ }^{69}$, Y.M. Kharzheev ${ }^{36}$, D. Khatidze ${ }^{70}$, T.J. Kim ${ }^{31}$, M.H. Kirby ${ }^{53}$, M. Kirsch ${ }^{21}$, B. Klima ${ }^{50}$, J.M. Kohli ${ }^{27}$, E.V. Komissarov ${ }^{36, \ddagger}$, J.-P. Konrath ${ }^{23}$, A.V. Kozelov ${ }^{39}$, J. Kraus ${ }^{65}$, T. Kuhl ${ }^{24}$, A. Kumar ${ }^{69}$, A. Kupco ${ }^{11}$, T. Kurča ${ }^{20}$, V.A. Kuzmin ${ }^{38}$, J. Kvita ${ }^{9}$, F. Lacroix ${ }^{13}$, D. Lam ${ }^{55}$, S. Lammers ${ }^{70}$, G. Landsberg ${ }^{77}$, P. Lebrun ${ }^{20}$, W.M. Lee ${ }^{50}$, A. Leflat ${ }^{38}$, J. Lellouch ${ }^{17}$, J. Li $^{78, \ddagger}$, L. Li $^{48}$, Q.Z. Li ${ }^{50}$, S.M. Lietti ${ }^{5}$, J.K. Lim ${ }^{31}$, J.G.R. Lima ${ }^{52}$, D. Lincoln ${ }^{50}$, J. Linnemann ${ }^{65}$, V.V. Lipaev ${ }^{39}$, R. Lipton ${ }^{50}$, Y. Liu ${ }^{7}$, Z. Liu ${ }^{6}$, A. Lobodenko ${ }^{40}$, M. Lokajicek ${ }^{11}$, P. Love ${ }^{42}$, H.J. Lubatti ${ }^{82}$, R. Luna ${ }^{3}$, A.L. Lyon ${ }^{50}$, A.K.A. Maciel ${ }^{2}$, D. Mackin ${ }^{80}$, R.J. Madaras ${ }^{46}$, P. Mättig ${ }^{26}$, C. Magass $^{21}$, A. Magerkurth ${ }^{64}$, P.K. Mal ${ }^{82}$, H.B. Malbouisson ${ }^{3}$, S. Malik ${ }^{67}$, V.L. Malyshev ${ }^{36}$, Y. Maravin ${ }^{59}$, B. Martin ${ }^{14}$, R. McCarthy ${ }^{72}$, A. Melnitchouk ${ }^{66}$, L. Mendoza ${ }^{8}$, P.G. Mercadante ${ }^{5}$, Y.P. Merekov ${ }^{36}$, M. Merkin ${ }^{38}$, K.W. Merritt ${ }^{50}$, A. Meyer ${ }^{21}$, J. Meyer ${ }^{22, c}$, J. Mitrevski ${ }^{70}$, R.K. Mommsen ${ }^{44}$, N.K. Mondal ${ }^{29}$, R.W. Moore ${ }^{6}$, T. Moulik ${ }^{58}$, G.S. Muanza ${ }^{20}$, M. Mulhearn ${ }^{70}$, O. Mundal ${ }^{22}$, L. Mundim ${ }^{3}$, E. Nagy ${ }^{15}$, M. Naimuddin ${ }^{50}$, M. Narain ${ }^{77}$, N.A. Naumann ${ }^{35}$, H.A. Neal ${ }^{64}$, J.P. Negret ${ }^{8}$, P. Neustroev ${ }^{40}$, H. Nilsen ${ }^{23}$, H. Nogima ${ }^{3}$, S.F. Novaes ${ }^{5}$, T. Nunnemann ${ }^{25}$, V. O'Dell ${ }^{50}$, D.C. O’Neil ${ }^{6}$, G. Obrant ${ }^{40}$, C. Ochando ${ }^{16}$, D. Onoprienko ${ }^{59}$, J. Orduna ${ }^{33}$, N. Oshima ${ }^{50}$, N. Osman ${ }^{43}$, J. Osta ${ }^{55}$, R. Otec ${ }^{10}$, G.J. Otero y Garzón ${ }^{50}$, M. Owen ${ }^{44}$, P. Padley ${ }^{80}$, M. Pangilinan ${ }^{77}$, N. Parashar ${ }^{56}$, S.-J. Park ${ }^{22, c}$, S.K. Park ${ }^{31}$, J. Parsons ${ }^{70}$, R. Partridge ${ }^{77}$, N. Parua ${ }^{54}$, A. Patwa ${ }^{73}$, G. Pawloski ${ }^{80}$, B. Penning ${ }^{23}$, M. Perfilov ${ }^{38}$, K. Peters ${ }^{44}$, Y. Peters ${ }^{26}$, P. Pétroff ${ }^{16}$, M. Petteni ${ }^{43}$,
R. Piegaia ${ }^{1}$, J. Piper ${ }^{65}$, M.-A. Pleier ${ }^{22}$, P.L.M. Podesta-Lerma ${ }^{33, d}$, V.M. Podstavkov ${ }^{50}$, Y. Pogorelov ${ }^{55}$, M.-E. Pol ${ }^{2}$, P. Polozov ${ }^{37}$, B.G. Pope ${ }^{65}$, A.V. Popov ${ }^{39}$, C. Potter ${ }^{6}$, W.L. Prado da Silva ${ }^{3}$, H.B. Prosper ${ }^{49}$, S. Protopopescu ${ }^{73}$, J. Qian ${ }^{64}$, A. Quadt ${ }^{22, c}$, B. Quinn ${ }^{66}$, A. Rakitine ${ }^{42}$, M.S. Rangel ${ }^{2}$, K. Ranjan ${ }^{28}$, P.N. Ratoff ${ }^{42}$, P. Renkel ${ }^{79}$, P. Rich ${ }^{44}$, J. Rieger ${ }^{54}$, M. Rijssenbeek ${ }^{72}$, I. Ripp-Baudot ${ }^{19}$, F. Rizatdinova ${ }^{76}$, S. Robinson ${ }^{43}$, R.F. Rodrigues ${ }^{3}$, M. Rominsky ${ }^{75}$, C. Royon ${ }^{18}$, A. Rozhdestvenski ${ }^{36}$, P. Rubinov ${ }^{50}$, R. Ruchti ${ }^{55}$, G. Safronov ${ }^{37}$, G. Sajot ${ }^{14}$, A. Sánchez-Hernández ${ }^{33}$, M.P. Sanders ${ }^{17}$, B. Sanghi ${ }^{50}$, G. Savage ${ }^{50}$, L. Sawyer ${ }^{60}$, T. Scanlon ${ }^{43}$, D. Schaile ${ }^{25}$, R.D. Schamberger ${ }^{72}$, Y. Scheglov ${ }^{40}$, H. Schellman ${ }^{53}$, T. Schliephake ${ }^{26}$, S. Schlobohm ${ }^{82}$, C. Schwanenberger ${ }^{44}$, A. Schwartzman ${ }^{68}$, R. Schwienhorst ${ }^{65}$, J. Sekaric ${ }^{49}$, H. Severini ${ }^{75}$, E. Shabalina ${ }^{51}$, M. Shamim ${ }^{59}$, V. Shary ${ }^{18}$, A.A. Shchukin ${ }^{39}$, R.K. Shivpuri ${ }^{28}$, V. Siccardi ${ }^{19}$, V. Simak ${ }^{10}$, V. Sirotenko ${ }^{50}$, P. Skubic ${ }^{75}$, P. Slattery ${ }^{71}$, D. Smirnov ${ }^{55}$, G.R. Snow ${ }^{67}$, J. Snow ${ }^{74}$, S. Snyder ${ }^{73}$, S. Söldner-Rembold ${ }^{44}$, L. Sonnenschein ${ }^{17}$, A. Sopczak ${ }^{42}$, M. Sosebee ${ }^{78}$, K. Soustruznik ${ }^{9}$, B. Spurlock ${ }^{78}$, J. Stark ${ }^{14}$, J. Steele ${ }^{60}$, V. Stolin ${ }^{37}$, D.A. Stoyanova ${ }^{39}$, J. Strandberg ${ }^{64}$, S. Strandberg ${ }^{41}$, M.A. Strang ${ }^{69}$, E. Strauss ${ }^{72}$, M. Strauss ${ }^{75}$, R. Ströhmer ${ }^{25}$, D. Strom ${ }^{53}$, L. Stutte ${ }^{50}$, S. Sumowidagdo ${ }^{49}$, P. Svoisky ${ }^{55}$, A. Sznajder ${ }^{3}$, P. Tamburello ${ }^{45}$, A. Tanasijczuk ${ }^{1}$, W. Taylor ${ }^{6}$, B. Tiller ${ }^{25}$, F. Tissandier ${ }^{13}$, M. Titov ${ }^{18}$, V.V. Tokmenin ${ }^{36}$, I. Torchiani ${ }^{23}$, D. Tsybychev ${ }^{72}$, B. Tuchming ${ }^{18}$,
C. Tully ${ }^{68}$, P.M. Tuts ${ }^{70}$, R. Unalan ${ }^{65}$, L. Uvarov ${ }^{40}$, S. Uvarov ${ }^{40}$, S. Uzunyan ${ }^{52}$, B. Vachon ${ }^{6}$, P.J. van den Berg ${ }^{34}$, R. Van Kooten ${ }^{54}$, W.M. van Leeuwen ${ }^{34}$, N. Varelas ${ }^{51}$, E.W. Varnes ${ }^{45}$, I.A. Vasilyev ${ }^{39}$, P. Verdier ${ }^{20}$, L.S. Vertogradov ${ }^{36}$, Y. Vertogradova ${ }^{36}$, M. Verzocchi ${ }^{50}$, D. Vilanova ${ }^{18}$, F. Villeneuve-Seguier ${ }^{43}$, P. Vint ${ }^{43}$, P. Vokac ${ }^{10}$, M. Voutilainen ${ }^{67, e}$, R. Wagner ${ }^{68}$, H.D. Wahl ${ }^{49}$, M.H.L.S. Wang ${ }^{50}$, J. Warchol ${ }^{55}$, G. Watts ${ }^{82}$, M. Wayne ${ }^{55}$, G. Weber ${ }^{24}$, M. Weber ${ }^{50, f}$, L. Welty-Rieger ${ }^{54}$, A. Wenger ${ }^{23, g}$, N. Wermes ${ }^{22}$, M. Wetstein ${ }^{61}$, A. White ${ }^{78}$, D. Wicke ${ }^{26}$, M. Williams ${ }^{42}$, G.W. Wilson ${ }^{58}$, S.J. Wimpenny ${ }^{48}$, M. Wobisch ${ }^{60}$, D.R. Wood ${ }^{63}$, T.R. Wyatt ${ }^{44}$, Y. Xie ${ }^{77}$, S. Yacoob ${ }^{53}$, R. Yamada ${ }^{50}$, W.-C. Yang ${ }^{44}$, T. Yasuda ${ }^{50}$, Y.A. Yatsunenko ${ }^{36}$, H. Yin ${ }^{7}$, K. Yip ${ }^{73}$, H.D. Yoo ${ }^{77}$, S.W. Youn ${ }^{53}$, J. Yu ${ }^{78}$, C. Zeitnitz ${ }^{26}$, S. Zelitch ${ }^{81}$, T. Zhao ${ }^{82}$, B. Zhou ${ }^{64}$, J. Zhu ${ }^{72}$, M. Zielinski ${ }^{71}$, D. Zieminska ${ }^{54}$, A. Zieminski ${ }^{54, \ddagger}$, L. Zivkovic ${ }^{70}$, V. Zutshi ${ }^{52}$, and E.G. Zverev ${ }^{38}$
(The DØ Collaboration)
${ }^{1}$ Universidad de Buenos Aires, Buenos Aires, Argentina
${ }^{2}$ LAFEX, Centro Brasileiro de Pesquisas Físicas, Rio de Janeiro, Brazil
${ }^{3}$ Universidade do Estado do Rio de Janeiro, Rio de Janeiro, Brazil
${ }^{4}$ Universidade Federal do ABC, Santo André, Brazil
${ }^{5}$ Instituto de Física Teórica, Universidade Estadual Paulista, São Paulo, Brazil
${ }^{6}$ University of Alberta, Edmonton, Alberta, Canada,
Simon Fraser University, Burnaby, British Columbia,
Canada, York University, Toronto, Ontario, Canada, and McGill University, Montreal, Quebec, Canada
${ }^{7}$ University of Science and Technology of China, Hefei, People's Republic of China ${ }^{8}$ Universidad de los Andes, Bogotá, Colombia
${ }^{9}$ Center for Particle Physics, Charles University, Prague, Czech Republic
${ }^{10}$ Czech Technical University, Prague, Czech Republic
${ }^{11}$ Center for Particle Physics, Institute of Physics, Academy of Sciences of the Czech Republic, Prague, Czech Republic
${ }^{12}$ Universidad San Francisco de Quito, Quito, Ecuador
${ }^{13}$ LPC, Université Blaise Pascal, CNRS/IN2P3, Clermont, France
${ }^{14}$ LPSC, Université Joseph Fourier Grenoble 1, CNRS/IN2P3,
Institut National Polytechnique de Grenoble, Grenoble, France
${ }^{15}$ CPPM, Aix-Marseille Université, CNRS/IN2P3, Marseille, France
${ }^{16}$ LAL, Université Paris-Sud, IN2P3/CNRS, Orsay, France
${ }^{17}$ LPNHE, IN2P3/CNRS, Universités Paris VI and VII, Paris, France
${ }^{18}$ CEA, Irfu, SPP, Saclay, France
${ }^{19}$ IPHC, Université Louis Pasteur, CNRS/IN2P3, Strasbourg, France
${ }^{20}$ IPNL, Université Lyon 1, CNRS/IN2P3, Villeurbanne, France and Université de Lyon, Lyon, France
${ }^{21}$ III. Physikalisches Institut A, RWTH Aachen University, Aachen, Germany
${ }^{22}$ Physikalisches Institut, Universität Bonn, Bonn, Germany
${ }^{23}$ Physikalisches Institut, Universität Freiburg, Freiburg, Germany
${ }^{24}$ Institut für Physik, Universität Mainz, Mainz, Germany
${ }^{25}$ Ludwig-Maximilians-Universität München, München, Germany
${ }^{26}$ Fachbereich Physik, University of Wuppertal, Wuppertal, Germany
${ }^{27}$ Panjab University, Chandigarh, India
${ }^{28}$ Delhi University, Delhi, India

${ }^{29}$ Tata Institute of Fundamental Research, Mumbai, India<br>${ }^{30}$ University College Dublin, Dublin, Ireland<br>${ }^{31}$ Korea Detector Laboratory, Korea University, Seoul, Korea<br>${ }^{32}$ SungKyunKwan University, Suwon, Korea<br>${ }^{33}$ CInVESTAV, Mexico City, Mexico<br>${ }^{34}$ FOM-Institute NIKHEF and University of Amsterdam/NIKHEF, Amsterdam, The Netherlands<br>${ }^{35}$ Radboud University Nijmegen/NIKHEF, Nijmegen, The Netherlands<br>${ }^{36}$ Joint Institute for Nuclear Research, Dubna, Russia<br>${ }^{37}$ Institute for Theoretical and Experimental Physics, Moscow, Russia<br>${ }^{38}$ Moscow State University, Moscow, Russia<br>${ }^{39}$ Institute for High Energy Physics, Protvino, Russia<br>${ }^{40}$ Petersburg Nuclear Physics Institute, St. Petersburg, Russia<br>${ }^{41}$ Lund University, Lund, Sweden, Royal Institute of Technology and Stockholm University,<br>Stockholm, Sweden, and Uppsala University, Uppsala, Sweden<br>${ }^{42}$ Lancaster University, Lancaster, United Kingdom<br>${ }^{43}$ Imperial College, London, United Kingdom<br>${ }^{44}$ University of Manchester, Manchester, United Kingdom<br>${ }^{45}$ University of Arizona, Tucson, Arizona 85721, USA<br>${ }^{46}$ Lawrence Berkeley National Laboratory and University of California, Berkeley, California 94720, USA<br>${ }^{47}$ California State University, Fresno, California 93740, USA<br>${ }^{48}$ University of California, Riverside, California 92521, USA<br>${ }^{49}$ Florida State University, Tallahassee, Florida 32306, USA<br>${ }^{50}$ Fermi National Accelerator Laboratory, Batavia, Illinois 60510, USA<br>${ }^{51}$ University of Illinois at Chicago, Chicago, Illinois 60607, USA<br>${ }^{52}$ Northern Illinois University, DeKalb, Illinois 60115, USA<br>${ }^{53}$ Northwestern University, Evanston, Illinois 60208, USA<br>${ }^{54}$ Indiana University, Bloomington, Indiana 47405, USA<br>${ }^{55}$ University of Notre Dame, Notre Dame, Indiana 46556, USA<br>${ }^{56}$ Purdue University Calumet, Hammond, Indiana 46323, USA<br>${ }^{57}$ Iowa State University, Ames, Iowa 50011, USA<br>${ }^{58}$ University of Kansas, Lawrence, Kansas 66045, USA<br>${ }^{59}$ Kansas State University, Manhattan, Kansas 66506, USA<br>${ }^{60}$ Louisiana Tech University, Ruston, Louisiana 71272, USA<br>${ }^{61}$ University of Maryland, College Park, Maryland 20742, USA<br>${ }^{62}$ Boston University, Boston, Massachusetts 02215, USA<br>${ }^{63}$ Northeastern University, Boston, Massachusetts 02115, USA<br>${ }^{64}$ University of Michigan, Ann Arbor, Michigan 48109, USA<br>${ }^{65}$ Michigan State University, East Lansing, Michigan 48824, USA<br>${ }^{66}$ University of Mississippi, University, Mississippi 38677, USA<br>${ }^{67}$ University of Nebraska, Lincoln, Nebraska 68588, USA<br>${ }^{68}$ Princeton University, Princeton, New Jersey 08544, USA<br>${ }^{69}$ State University of New York, Buffalo, New York 14260, USA<br>${ }^{70}$ Columbia University, New York, New York 10027, USA<br>${ }^{71}$ University of Rochester, Rochester, New York 14627, USA<br>${ }^{72}$ State University of New York, Stony Brook, New York 11794, USA<br>${ }^{73}$ Brookhaven National Laboratory, Upton, New York 11973, USA<br>${ }^{74}$ Langston University, Langston, Oklahoma 73050, USA<br>${ }^{75}$ University of Oklahoma, Norman, Oklahoma 73019, USA<br>${ }^{76}$ Oklahoma State University, Stillwater, Oklahoma 74078, USA<br>${ }^{77}$ Brown University, Providence, Rhode Island 02912, USA ${ }^{78}$ University of Texas, Arlington, Texas 76019, USA<br>${ }^{79}$ Southern Methodist University, Dallas, Texas 75275, USA<br>${ }^{80}$ Rice University, Houston, Texas 77005, USA<br>${ }^{81}$ University of Virginia, Charlottesville, Virginia 22901, USA and<br>${ }^{82}$ University of Washington, Seattle, Washington 98195, USA

(Dated: August 29, 2008)
We report the observation of the doubly strange $b$ baryon $\Omega_{b}^{-}$in the decay channel $\Omega_{b}^{-} \rightarrow J / \psi \Omega^{-}$, with $J / \psi \rightarrow \mu^{+} \mu^{-}$and $\Omega^{-} \rightarrow \Lambda K^{-} \rightarrow\left(p \pi^{-}\right) K^{-}$, in $p \bar{p}$ collisions at $\sqrt{s}=1.96 \mathrm{TeV}$. Using approximately $1.3 \mathrm{fb}^{-1}$ of data collected with the D 0 detector at the Fermilab Tevatron Collider, we observe $17.8 \pm 4.9$ (stat.) $\pm 0.8$ (syst.) $\Omega_{b}^{-}$signal events at a mass of $6.165 \pm 0.010$ (stat.) $\pm 0.013$ (syst.) GeV . The significance of the observed signal is $5.4 \sigma$, corresponding to a probability of $6.7 \times 10^{-8}$ of it arising from a background fluctuation.

The $\Omega^{-}$baryon, composed of three strange quarks, played an important historical role in our understanding of the basic structure of matter. Its discovery in 1964 [1] at a mass predicted from $\mathrm{SU}(3)$ symmetry breaking was a great success for the theory [2]. The $\Omega_{b}^{-}$(bss) (charge conjugate states are assumed throughout this Letter) is a predicted heavy cousin of the $\Omega^{-}$with a $b$ quark replacing one of the three strange quarks. While the $\Omega^{-}$ has $J^{P}=3 / 2^{+}$, the ground state $\Omega_{b}^{-}$is expected to have $J^{P}=1 / 2^{+}$, a mass between $5.94-6.12 \mathrm{GeV}$ and a lifetime such that $0.55<\tau\left(\Omega_{b}^{-}\right) / \tau\left(B^{0}\right)<1.10$ [3].

In this Letter, we report the first observation of the $\Omega_{b}^{-}$ baryon, fully reconstructed from its decay $\Omega_{b}^{-} \rightarrow J / \psi \Omega^{-}$, with $J / \psi \rightarrow \mu^{+} \mu^{-}, \Omega^{-} \rightarrow \Lambda K^{-}$and $\Lambda \rightarrow p \pi^{-}$. The analysis is based on a data sample of $1.3 \mathrm{fb}^{-1}$ collected in $p \bar{p}$ collisions at $\sqrt{s}=1.96 \mathrm{TeV}$ with the D 0 detector (4] at the Fermilab Tevatron Collider. The detector components most relevant to this analysis are the central tracking system and the muon spectrometer. The central tracking system consists of a silicon microstrip tracker (SMT) and a central fiber tracker (CFT) inside a 2 Tesla superconducting solenoid. The SMT is optimized for tracking and vertexing over the pseudorapidity region $|\eta|<3$ while the CFT has coverage for $|\eta|<2$. A liquid argon and uranium calorimeter provides coverage up to $|\eta|<4.2$. The muon spectrometer covers $|\eta|<2$.

The $\Omega_{b}^{-} \rightarrow J / \psi \Omega^{-} \rightarrow J / \psi \Lambda K^{-} \rightarrow J / \psi p \pi^{-} K^{-}$decay topology is similar to that of the $\Xi_{b}^{-} \rightarrow J / \psi \Xi^{-} \rightarrow$ $J / \psi \Lambda \pi^{-} \rightarrow J / \psi p \pi^{-} \pi^{-}$decay with $\Omega^{-}$in place of $\Xi^{-}$ and $K^{-}$in place of $\pi^{-}$. Consequently, the reconstruction of the $J / \psi$ and $\Lambda$ and their selection discussed below follow closely the analysis that led to the first direct observation of the $\Xi_{b}^{-}$baryon [5]. However, in this analysis we use a multivariate selection for the $\Omega^{-}$owing to the smaller signal to background ratio compared to that for the $\Xi^{-}$in the $\Xi_{b}^{-}$analysis. We use the pythia Monte Carlo (MC) program [6] to generate $\Omega_{b}^{-}$and the EVTGEN program [7] to simulate $\Omega_{b}^{-}$decays. The $\Omega_{b}^{-}$mass and lifetime are set to be 6.052 GeV and 1.54 ps respectively. The generated events are subjected to a GEANT [8] based D0 detector simulation, and to the same reconstruction and selection programs as the data. We optimize the $\Omega^{-}$ selection using MC $\Omega_{b}^{-}$events for the signal and a sample of $J / \psi\left(\Lambda K^{+}\right)$data (referred to below as wrong-sign events) for the background, while keeping the $J / \psi \Omega^{-}$ data blinded. Once all selection criteria have been determined, we apply them to the $J / \psi \Omega^{-}$data.

We begin the event selection by reconstructing $J / \psi \rightarrow$ $\mu^{+} \mu^{-}$candidates from two oppositely charged muons with transverse momentum $\left(p_{T}\right)$ greater than 1.5 GeV that are compatible with being from a common vertex. Muons are identified by tracks reconstructed in the central tracking system that are matched with either track
segments in the muon spectrometer or calorimeter energy deposits consistent with a minimum ionizing particle. Events must have a well-reconstructed $p \bar{p}$ interaction point that we take to be the $\Omega_{b}^{-}$production vertex and a $J / \psi \rightarrow \mu^{+} \mu^{-}$candidate in the mass window $2.75<M_{\mu \mu}<3.40 \mathrm{GeV}$. Events with $J / \psi$ candidates are reprocessed with a version of the track reconstruction algorithm that increases the efficiency for tracks with low $p_{T}$ and high impact parameters.

We form $\Lambda \rightarrow p \pi^{-}$candidates from two oppositely charged particles, each with $p_{T}>0.2 \mathrm{GeV}$, that are consistent with having originated from a common vertex. The two tracks are required to have a total of no more than two hits in the tracking system before the reconstructed $p \pi^{-}$vertex. The impact parameter significance (the impact parameter with respect to the $p \bar{p}$ interaction point divided by its uncertainty) must exceed four for at least one of the tracks and three for the other. The track with the higher $p_{T}$ is assumed to be the proton. MC studies show that this assignment leads to the correct combination nearly $100 \%$ of the time. Furthermore, we require the $\Lambda$ transverse decay length to be greater than four times its uncertainty and the proper decay length to exceed ten times its uncertainty, where the transverse decay length is the distance between the production and decay vertices in the transverse plane while the proper decay length is the transverse decay length corrected by the Lorentz boost calculated from $p_{T}(\Lambda) . \Lambda$ candidates must have reconstructed masses between 1.108 and 1.126 GeV .


FIG. 1: The invariant mass distribution of the $\Lambda K$ pair before (a) and after (b) the BDT selection. Filled circles are from the right-sign $\Lambda K^{-}$events while the histogram is from the wrongsign $\Lambda K^{+}$events without any additional normalization.

We combine $\Lambda$ candidates with negatively charged particles (assumed to be kaons) to form $\Omega^{-} \rightarrow \Lambda K^{-}$decay candidates. The $\Lambda$ and the kaon are required to have a common vertex. The $\Omega^{-}$candidates must have transverse decay length significances greater than four and the uncertainties of the proper decay lengths less than 0.5 cm . These two requirements reduce backgrounds from combinatorics and mismeasured tracks. The $\Xi^{-}$baryon has
a mass of 1.322 GeV [9] and decays into $\Lambda \pi^{-}$. If the kaon mass is assigned to the pion, this decay could be a major background for $\Omega^{-} \rightarrow \Lambda K^{-}$. To eliminate this background, we remove candidates with $\Lambda \pi^{-}$mass less than 1.34 GeV . Figure 1 (a) shows the mass distribution of the reconstructed $\Omega^{-} \rightarrow \Lambda K^{-}$candidates after these selections. The distribution of wrong-sign $\Lambda K^{+}$events is also shown. An excess of events above the background around the expected $\Omega^{-}$mass of 1.672 GeV [9] is visible in the distribution of the right-sign $\Lambda K^{-}$events.

To further enhance the $\Omega^{-}$signal over the combinatorial background, kinematic variables associated with daughter particle momenta, vertices, and track qualities are combined using Boosted Decision Trees (BDT) 10, 11]. The $\Lambda K^{-}$mass distribution after the BDT selection is shown in Fig. [1(b). The BDT selection retains $87 \%$ of the $\Omega^{-}$signal while rejecting $89 \%$ of the background. The enhanced $\Omega^{-}$mass peak is evident in the distribution. A $\Lambda K^{-}$pair is considered to be a $\Omega^{-}$candidate if its mass is in the range of $1.662-1.682 \mathrm{GeV}$.

To select $\Omega_{b}^{-} \rightarrow J / \psi \Omega^{-}$candidates, we develop selection criteria using the MC $\Omega_{b}^{-}$events as the signal and the data wrong-sign events as the background. The background events are formed by combining $J / \psi$ candidates with $\Lambda K^{+}$pairs with mass between 1.662 and 1.682 GeV . We form $\Omega_{b}^{-} \rightarrow J / \psi \Omega^{-}$decay candidates from $J / \psi$ and $\Omega^{-}$pairs that are consistent with being from a common vertex. We require the uncertainty of the $\Omega_{b}^{-}$proper decay length to be less than 0.03 cm and impose a minimum $p_{T}$ cut of 6 GeV on the $\Omega_{b}^{-}$candidates. Finally, $J / \psi$ and $\Omega^{-}$daughters from the $\Omega_{b}^{-}$decays are expected to be boosted in the direction of the $\Omega_{b}^{-}$; therefore, we require the opening angle in the transverse plane between the $J / \psi$ and the $\Omega^{-}$to be less than $\pi / 2$.

We then apply the above selections to the right-sign events in the data to search for the $\Omega_{b}^{-}$baryon in the mass window between 5.6 and 7.0 GeV . This range is chosen since 5.624 GeV is the mass of the lightest $b$ baryon, the $\Lambda_{b}$, and the upper limit of 7.0 GeV is nearly 1 GeV higher than the predicted $\Omega_{b}^{-}$mass [3]. We calculate the $\Omega_{b}^{-}$candidate mass using the formula $M\left(\Omega_{b}^{-}\right)=$ $M\left(J / \psi \Omega^{-}\right)-M\left(\mu^{+} \mu^{-}\right)-M\left(\Lambda K^{-}\right)+\hat{M}(J / \psi)+\hat{M}\left(\Omega^{-}\right)$. Here $M\left(J / \psi \Omega^{-}\right), M\left(\mu^{+} \mu^{-}\right)$, and $M\left(\Lambda K^{-}\right)$are the reconstructed masses while $\hat{M}(J / \psi)$ and $\hat{M}\left(\Omega^{-}\right)$are taken from Ref. [9]. This calculation improves the mass resolution of the $\mathrm{MC} \Omega_{b}^{-}$events from 0.080 GeV to 0.034 GeV . In the mass search window, we observe 79 candidates in the data with the mass distribution shown in Fig. 2(a). An excess of events near 6.2 GeV is apparent. No such structure, however, is seen in the corresponding mass distribution (Fig. 2(b)) of the 30 wrong-sign events.

Assuming the excess is due to the $\Omega_{b}^{-}$production, we fit $\Omega_{b}^{-}$candidate masses with the hypothesis of a Gaussian signal plus a flat background using an unbinned likelihood method. We fix the Gaussian width to 0.034 GeV , the width of the MC $\Omega_{b}^{-}$signal. The fit gives an $\Omega_{b}^{-}$mass
of $6.165 \pm 0.010$ (stat.) GeV and a yield of $17.8 \pm 4.9$ (stat.) signal events. To assess the significance of the excess, we first determine the likelihood $\mathcal{L}_{s+b}$ of the signal plus background fit above and then repeat the fit with only the background contribution to find a new likelihood $\mathcal{L}_{b}$. The logarithmic likelihood ratio $\sqrt{2 \ln \left(\mathcal{L}_{s+b} / \mathcal{L}_{b}\right)}$ yields a statistical significance of $5.4 \sigma$, equivalent to a probability of $6.7 \times 10^{-8}$ that the background could fluctuate with a significance equal to or greater than what is observed. Fitted yields for positively and negatively charged candidates are $6.2 \pm 3.1$ (stat.) $\Omega_{b}^{+}$and $12.0 \pm 3.9$ (stat.) $\Omega_{b}^{-}$, respectively.


FIG. 2: (a) The $M\left(\Omega_{b}^{-}\right)$distribution of the $\Omega_{b}^{-}$candidates after all selection criteria. The dotted curve is an unbinned likelihood fit to the model of a constant background plus a Gaussian signal. The mass distributions for the wrong-sign background (b), the $\Omega^{-}$sideband events (c), and the $\Lambda$ sideband events (d).

Various checks have been performed to ensure that the observed resonance is genuine. (1) We apply the event selection to data events in the sidebands of the reconstructed $\Omega^{-}$and $\Lambda$ resonances separately. The $J / \psi\left(p \pi^{-}\right) K^{-}$mass distributions of these sideband events are shown in Figs. 2(c) and 2(d). No significant excess is present in either distribution. (2) We investigate the possibility of a false signal due to residual $b$ hadron backgrounds by applying the final $\Omega_{b}^{-}$selection to $\mathrm{MC} B^{-} \rightarrow J / \psi K^{*-} \rightarrow J / \psi K_{S}^{0} \pi^{-}, \Xi_{b}^{-} \rightarrow J / \psi \Xi^{-}$, and $\Lambda_{b} \rightarrow J / \psi \Lambda$ samples with equivalent integrated luminosities significantly greater than that of the analyzed data. No indication of any resonance is observed in the reconstructed $J / \psi \Omega^{-}$mass distribution. (3) We check the mass distributions of the $\Omega_{b}^{-}$decay products. For $\Omega_{b}^{-}$candidates within a $\pm 3 \sigma$ mass window around the observed peak, we relax the mass requirements on the $\Omega^{-}$and $\Lambda$ candidates and perform a fit to each mass distribution. The numbers of the $\Omega^{-}$and $\Lambda$ candidates from the fits are consistent with the observed number of $\Omega_{b}^{-}$ signal events. (4) We replace the BDT selection with individual cuts on the most important variables according to the BDT optimization and confirm the existence of a
peak with a comparable event yield but a higher background at a mass consistent with that observed using the BDT. (5) We test the robustness of the peak by varying selections such as the $\Xi^{-}$veto, $\Lambda$ and $\Omega^{-}$mass windows, $\Lambda$ transverse decay requirements, BDT selection, and the requirement on $p_{T}\left(\Omega_{b}^{-}\right)$. All the above studies confirm the existence of the resonance.

Potential sources of systematic uncertainties on the measured $\Omega_{b}^{-}$mass include event selection, signal and background models, and momentum scale. Varying the selection criteria and applying a set of cuts on individual kinematic variables lead to a maximum change of 0.012 GeV in the fitted mass. Using a linear function as the background model results in negligible change in the mass. Varying the Gaussian width in the signal model between 0.028 and 0.040 GeV changes the fitted mass by at most 0.003 GeV . When a tighter selection is applied to enhance signal over background, we can float the width of the signal model in the fit. This leads to a mass shift of 0.002 GeV and a fitted signal width of $0.033 \pm 0.010 \mathrm{GeV}$, consistent with the MC expectation. To study the effect of the track momentum scale uncertainty on the measured $\Omega_{b}^{-}$mass, we reconstruct the higher statistics $\Lambda_{b} \rightarrow J / \psi \Lambda$ decays and measure the $\Lambda_{b}$ mass for different minimum $p_{T}$ requirements on the $\Lambda_{b}$ daughter particles. We compare these measurements to the world average value of the $\Lambda_{b}$ mass [9] and take the maximum deviation of 0.004 GeV as a systematic uncertainty. Adding in quadrature, we get a total systematic uncertainty of 0.013 GeV to obtain a measured $\Omega_{b}^{-}$mass: $6.165 \pm 0.010$ (stat.) $\pm 0.013$ (syst.) GeV. Similarly, we estimate the systematic uncertainty on the $\Omega_{b}^{-}$yield by varying the signal and background models in the fit. The observed maximum change of 0.8 is assigned as the systematic uncertainty on the yield: $17.8 \pm 4.9$ (stat.) $\pm 0.8$ (syst.). In all these studies, the signal significance remains above $5 \sigma$.


FIG. 3: The distribution of the proper decay length of the $\Omega_{b}^{-}$candidates in the $\pm 3 \sigma$ mass window around the observed peak along with the expected distribution from the MC $\Omega_{b}^{-}$ signal with a lifetime 1.54 ps plus the data background events.

Figure 3 shows the distribution of the proper decay length of the $\Omega_{b}^{-}$candidates observed in the $\pm 3 \sigma$ mass window around the fitted $\Omega_{b}^{-}$mass. The distribution of the MC $\Omega_{b}^{-}$signal plus the data background events is also
shown. The background distribution is modeled using events in the $\Omega_{b}^{-}$sidebands of $5.8-6.0$ and $6.4-6.6 \mathrm{GeV}$. Despite the low statistics, the data distribution contains significantly more positive than negative decay lengths as expected and consistent with a weakly decaying $b$ hadron.

We calculate the $\Omega_{b}^{-}$production rate relative to that of the $\Xi_{b}^{-}[5]$. The selection efficiency ratio $\epsilon\left(\Omega_{b}^{-} \rightarrow\right.$ $\left.J / \psi \Omega^{-}\right) / \epsilon\left(\Xi_{b}^{-} \rightarrow J / \psi \Xi^{-}\right)$is found to be $1.5 \pm 0.2$ (stat.) assuming inclusive $\Omega^{-}$and $\Xi^{-}$decays. The higher efficiency for the $\Omega_{b}^{-}$is due primarily to a harder $p_{T}$ spectrum of the kaon from the $\Omega^{-}$decay than that of the pion from the $\Xi^{-}$decay and a shorter lifetime of the $\Omega^{-}$ compared to the $\Xi^{-}$. By using the reported $\Xi_{b}^{-}$events [5] and the observed $\Omega_{b}^{-}$yield here, we estimate

$$
\mathcal{R}=\frac{f\left(b \rightarrow \Omega_{b}^{-}\right)}{f\left(b \rightarrow \Xi_{b}^{-}\right)} \cdot \frac{\operatorname{Br}\left(\Omega_{b}^{-} \rightarrow J / \psi \Omega^{-}\right)}{\operatorname{Br}\left(\Xi_{b}^{-} \rightarrow J / \psi \Xi^{-}\right)}
$$

to be $\mathcal{R}=0.80 \pm 0.32$ (stat.) $)_{-0.22}^{+0.14}$ (syst.). Here $f\left(b \rightarrow \Omega_{b}^{-}\right)$ and $f\left(b \rightarrow \Xi_{b}^{-}\right)$are the fractions of $b$ quarks that hadronize to form $\Omega_{b}^{-}$and $\Xi_{b}^{-}$, respectively. The systematic uncertainty includes contributions from the signal yields as well as the efficiency ratio. Using $\Gamma\left(\Omega_{b}^{-} \rightarrow\right.$ $\left.J / \psi \Omega^{-}\right) / \Gamma\left(\Xi_{b}^{-} \rightarrow J / \psi \Xi^{-}\right)=9.8$ [12], the central values of $\tau\left(\Xi_{b}^{-}\right)=1.42_{-0.24}^{+0.28} \mathrm{ps}$ [9], the $\mathcal{R}$ value above, and $\tau\left(\Omega_{b}^{-}\right)$in the range of $0.83-1.67 \mathrm{ps}$ [3], we obtain $f\left(b \rightarrow \Omega_{b}^{-}\right) / f\left(b \rightarrow \Xi_{b}^{-}\right) \approx 0.07-0.14$.

In summary, by analyzing $1.3 \mathrm{fb}^{-1}$ of data collected by the D0 experiment in $p \bar{p}$ collisions at $\sqrt{s}=1.96 \mathrm{TeV}$ at the Fermilab Tevatron Collider, we have made the first observation of the doubly strange $b$ baryon $\Omega_{b}^{-}$in the fully reconstructed decay mode $\Omega_{b}^{-} \rightarrow J / \psi \Omega^{-}$with $J / \psi \rightarrow$ $\mu^{+} \mu^{-}, \Omega^{-} \rightarrow \Lambda K^{-}$and $\Lambda \rightarrow p \pi^{-}$. We measure the $\Omega_{b}^{-}$ mass to be $6.165 \pm 0.010$ (stat.) $\pm 0.013$ (syst.) GeV. The significance of the observed signal is greater than $5 \sigma$.

We thank the staffs at Fermilab and collaborating institutions, and acknowledge support from the DOE and NSF (USA); CEA and CNRS/IN2P3 (France); FASI, Rosatom and RFBR (Russia); CNPq, FAPERJ, FAPESP and FUNDUNESP (Brazil); DAE and DST (India); Colciencias (Colombia); CONACyT (Mexico); KRF and KOSEF (Korea); CONICET and UBACyT (Argentina); FOM (The Netherlands); STFC (United Kingdom); MSMT and GACR (Czech Republic); CRC Program, CFI, NSERC and WestGrid Project (Canada); BMBF and DFG (Germany); SFI (Ireland); The Swedish Research Council (Sweden); CAS and CNSF (China); and the Alexander von Humboldt Foundation (Germany).

[^0][d] Visitor from ECFM, Universidad Autonoma de Sinaloa, Culiacán, Mexico.
[e] Visitor from Helsinki Institute of Physics, Helsinki, Finland.
[f] Visitor from Universität Bern, Bern, Switzerland.
[g] Visitor from Universität Zürich, Zürich, Switzerland.
[ $\ddagger$ ] Deceased.
[1] V.E. Barnes et al., Phys. Rev. Lett. 12, 204 (1964).
[2] M. Gell-Mann and Y. Ne'eman, The Eightfold Way, (W.A. Benjamin, New York, 1964).
[3] X. Liu et al., Phys. Rev. D 77, 014031 (2008); M. Karliner et al., arXiv:0804.1575 (2008); E.E. Jenkins, Phys. Rev. D 77, 034012 (2008); R. Roncaglia, D.B. Lichtenberg, and E. Predazzi, Phys. Rev. D 52, 1722 (1995); N. Mathur, R. Lewis, and R.M. Woloshyn, Phys. Rev. D 66, 014502 (2002); D. Ebert, R.N. Faustov, and V.O. Galkin, Phys. Rev. D 72, 034026 (2005); T. Ito, M. Matsuda, and
Y. Matsui, Prog. Theor. Phys. 99, 271 (1998).
[4] V.M. Abazov et al. (D0 Collaboration), Nucl. Instrum. Methods Phys. Res. A 565, 463 (2006).
[5] V.M. Abazov et al. (D0 Collaboration), Phys. Rev. Lett. 99, 052001 (2007).
[6] T. Sjöstrand et al., Comput. Phys. Commun. 135, 238 (2001).
[7] D.J. Lange, Nucl. Instrum. Methods Phys. Res. A 462, 152 (2001).
[8] R. Brun and F. Carminati, CERN Program Library Writeup W5013, 1993 (unpublished).
[9] C. Amsler et al., Physics Letters B 667, 1 (2008).
[10] L. Breiman et al., Classification and Regression Trees, (Wadsworth, Stamford, 1984).
[11] A. Höcker et al., arXiv:physics/0703039 [physics.data-an] (2007).
[12] H.-Y. Cheng, Phys. Rev. D 56, 2799 (1997).


[^0]:    [a] Visitor from Augustana College, Sioux Falls, SD, USA.
    [b] Visitor from The University of Liverpool, Liverpool, UK.
    [c] Visitor from II. Physikalisches Institut, Georg-AugustUniversity, Göttingen, Germany.

