

# 1 Micromechanical Modeling of Short Crack Nucleation 2 and Growth in High Cycle Fatigue of Martensitic 3 Microstructures

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## 9 **Abstract**

10 High cycle fatigue (HCF) is a frequently limiting failure mechanism of ma-  
11 chine elements and modern high strength steels. Present day design rules rely  
12 on semi-empirical methods, guidelines and utilization of macroscopic analysis  
13 means in origin, such as fracture mechanics. The resulting challenge is that  
14 short crack regime, critical for HCF in terms of lifetime of components and  
15 products, is somewhat poorly handled. This is an outcome of the fact that  
16 the present means and methodologies do not explicitly account for effects  
17 arising from material microstructure, an oversight micromechanics aims to  
18 rectify.

19 Micromechanical modeling operating on fatigue at the scale of material  
20 microstructure necessitates the introduction of suitable means to describe  
21 the mechanisms of cyclic plastic deformation and microstructural morpholo-  
22 gies, considered critical for HCF especially at the early stages of micro-crack  
23 nucleation and damage evolution towards and within the short crack regime.  
24 In current work, a crystal plasticity based approach with combined harden-  
25 ing is utilized to capture the respective deformation response utilizing full  
26 field modeling. The modeling is carried out for both simplified prior austen-  
27 ite grain like microstructures as well as complex imaging based martensitic  
28 quenched and tempered steel microstructural models. A fully coupled dam-  
29 age modeling scheme is introduced to track damage nucleation and evolution  
30 at the scale of the studied microstructures. Crack closure is included within  
31 the approach to track behavior of microstructure scale defects under, e.g.,  
32 fully reversed loading, more realistically. Model calibration is addressed and

33 application cases involving damage and crack growth both under monotonic  
34 and cyclic loading are presented.

35 The results demonstrate how the coupling of damage to crystal plasticity  
36 modeling can be utilized to identify and track the evolution of microstructure  
37 scale damage mechanisms in complex martensitic microstructures. Interac-  
38 tions between strain localization and damage accumulation are presented as  
39 well as transition from micro-cracking to short crack growth. The results  
40 show that the proposed approach can interpret the intricate dependencies  
41 and relations between complex microstructures, their (cyclic) deformation  
42 mechanisms and evolution of damage, the outcomes regarding crack forma-  
43 tion and behavior are found to be in line with similar experimental studies.

44 The proposed framework for modeling damage in polycrystalline mi-  
45 crostructures is quite general in its capabilities. By solely introducing a  
46 suitable crystal plasticity based deformation model and a damage model de-  
47 scribing nucleation and softening can plastic slip and damage interactions be  
48 studied in complex microstructures, and in principle, on any system where  
49 similar constitutive models are utilizable. The exploitation of the resulting  
50 micromechanical modeling and simulation capabilities lies both in simulation  
51 driven design of fatigue resistant components and high strength steels.

52 *Keywords:* High Cycle Fatigue, Crystal Plasticity, Micromechanics, High  
53 Strength Steel

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## 54 1. Introduction

55 High cycle fatigue (HCF) is often a limiting failure mechanism when it  
56 comes to structural and mechanical machine elements [1, 2]. Studies on HCF  
57 have indicated that the role of microstructural defects and initiation sites for  
58 microstructurally short fatigue cracks dominate the respective loading cy-  
59 cles to failure, i.e., most of the fatigue life is spent while the defect is of  
60 microstructure scale and smaller than some respective characteristic mate-  
61 rial dimension, such as prior austenite grain size (PAG). As such it comes  
62 as no surprise that ever more emphasis is being placed on understanding,  
63 and particularly, quantitatively modeling, the interactions between the mi-  
64 crostructure, different microstructural defects, and nonmetallic inclusions in  
65 modern high strength steels.

66 The present state-of-the-art as far as evaluation of HCF affiliated mate-  
67 rial properties and design guidelines is largely presented by Murakami and

68 Endo [3] who performed a thorough review of fatigue assessment methods for  
69 fatigue failures initiating from non-metallic inclusions. Their method to pre-  
70 dict fatigue strength has gained momentum and is described by the following  
71 points [4]:

- 72 • Short fatigue cracks threshold stress intensity factor range depends on  
73 the inclusion size with logarithmic slope of  $1/3$
- 74 • Short fatigue crack threshold depends linearly on the matrix hardness  
75 (up to 720 HV)
- 76 • Inclusions behave similarly in fatigue as holes produced by focused ion  
77 beam or drilling
- 78 • Fracture mechanics based size indicator of projected  $\sqrt{area}$  of the in-  
79 clusion in the direction of the maximum principal stress characterizes  
80 the maximum stress intensity factor of many differently shaped cracks  
81 (aspect ratio  $a/b < 5$ )

82 The developments in crystal plasticity modeling of cyclic plastic slip response  
83 at the scale of the material microstructure have led to ever increasing num-  
84 ber of studies attempting to establishing relations between plastic slip, stress  
85 triaxiality and fatigue failure mechanisms. The resulting micromechanical  
86 models have approached the topic in two principally different ways: firstly, by  
87 defining oftentimes cycle specific Fatigue Indicator Parameters (FIPs), and,  
88 secondly, by direct approaches where material microstructure scale plastic  
89 slip and stress state are utilized to assess material damage due to cyclic load-  
90 ing. With respect to such efforts, Sangid et al. [5, 6, 7] used persistent slip  
91 band’s energy instability in conjunction with dislocation creation in microme-  
92 chanical configuration to predict fatigue life scatter for U720. Pineau and For-  
93 est [8] performed elastic-plastic finite element calculations to assess the role  
94 of cyclic plasticity and residual stresses for very high cycle fatigue (VHCF)  
95 at the matrix-inclusion interface. Proudhon et al. [9] performed a 3D simula-  
96 tion of short fatigue crack propagation incorporating finite elements, crystal  
97 plasticity and remeshing. This is one of the few works which considers dam-  
98 age evolution beyond FIP like parameters but also models short fatigue crack  
99 propagation, albeit in a simplistic setting. Frondelius et al. [1, 10] studied  
100 inclusion to microstructure interactions in a high strength steel to develop  
101 micromechanical analysis methods for assessing the influence of defect struc-  
102 tures to nucleation and short crack regime damage evolution in the HCF

103 region. In Li et al. [11] the cyclic behaviour at the sub-grain level was  
 104 predicted and the effects of lath and precipitate sizes examined for elevated  
 105 temperature response of respective steels. A crystallographic, accumulated  
 106 slip (strain) parameter, modulated by triaxiality, was implemented at the  
 107 micro-scale in order to predict crack initiation in precipitate-strengthened  
 108 laths. In [12] Guan et al. replicated single and oligocrystal microstructures  
 109 and their deformation behavior with crystal plasticity finite element models  
 110 and under fatigue loading performed analyses that enabled grain-by-grain  
 111 comparison of measured and calculated slip to be carried out. Single and  
 112 multiple slip activation, slip localization and microstructure-sensitive stress  
 113 evolution were examined. The initiation of slip in the single crystals was  
 114 found to be highly localized in bands, of about  $1\text{--}2\ \mu\text{m}$  thickness and  $30\ \mu\text{m}$   
 115 separation. With increasing loading, the bands thickened to develop a  
 116 more uniform field of straining, these features of deformation behavior were  
 117 captured with crystal plasticity modelling. Bribier et al. [13] performed crys-  
 118 tal plasticity analysis on Ti-6Al-4V under high cycle fatigue. Their model  
 119 employed softening model related to dislocation slip that eventually triggers  
 120 fatigue. A simplified 3D microstructure was used to analyze performance of  
 121 the microstructure aiming to capture damage nucleation and growth in two  
 122 phase microstructure.

123 For some materials, physically material short cracks have been found to  
 124 initiate easily and the fatigue limit is defined by retardation of the growth  
 125 of these cracks [4]. The near-threshold crack growth behaviour emphasizes  
 126 roughness- and oxide-induced crack closure [14]. Much of the experimental  
 127 work studying crack closure has been focused on positive stress ratios. Pom-  
 128 mier et al. [15] and Silva [16, 17] found that for crack closure under negative  
 129 stress ratios the role of material's cyclic plastic properties, and especially  
 130 kinematic hardening, is emphasized and a clear dependence on the stress  
 131 amplitude was found. The levels of crack closure approached those of zero  
 132 stress ratio near threshold as cyclic plasticity is suppressed [15]. Develop-  
 133 ment of crack closure is also emphasized with cracks growing from notches  
 134 and could be important in the case of cracks initiating from non-metallic  
 135 inclusions. See Vaara et al. [18] for a brief review on high-cycle fatigue  
 136 with focus on non-metallic inclusions and forming. In any event, due to  
 137 the anisotropy and complex microstructural morphologies at the scale of the  
 138 material microstructure, it is envisioned that particular emphasis needs to  
 139 be placed on the capability to consider microstructure scale crack closure  
 140 associated behavior.

141 Next generation fatigue design methodologies aim to utilize microme-  
142 chanics at their core. The impact being sought is improved accuracy to be  
143 able to avoid unexpected failures and improve power density of products and  
144 components while doing so. Also, steps toward virtual testing of materials  
145 are to be expected. Simultaneously, if microstructure scale phenomena dom-  
146 inating fatigue life can be accurately modeled, it opens up the possibility to  
147 virtually design novel materials and optimize material solutions. For defect  
148 critical failure mechanisms such as HCF, this will make it possible to system-  
149 atically develop the microstructures and manufacturing processes to target  
150 performances and functionality profiles of specific products.

151 The focus on current study is in developing and demonstrating fully mi-  
152 crostructure and deformation coupled damage mechanical models for mi-  
153 cromechanical evaluation of fatigue, that is applicable also to HCF. Rather  
154 than utilize, e.g., FIP parameters to correlate to cyclic fatigue damage, a  
155 direct fully crystal plasticity coupled approach is adopted. The constitutive  
156 model is applied for different complexity containing models of martensitic  
157 microstructure, typical to those of quench and temper (Q&T) steels. A  
158 "smeared" approach tracking crack closure is integrated in the model to in-  
159 clude the effects of crack face contact. The material model is calibrated for  
160 a Q&T-steel based on digital image correlation (DIC) instrumented tensile  
161 tests. As use cases deformation and damage evolution under monotonic and  
162 cyclic loadings is investigated and compared to, e.g., typical crack propaga-  
163 tion behavior and trajectories in martensitic microstructures. The demon-  
164 strative results provide a basis for utilizing micromechanical modeling to di-  
165 rectly infer and investigate damage accumulation in complex microstructures  
166 related to HCF of high strength steels.

## 167 **2. Materials and methods**

### 168 *2.1. Material*

169 Wrought Q&T steel of 34CrNiMo6 grade with a nominal tensile strength  
170 of 1025 MPa and a nano-hardness of  $4.7 \pm 0.3$  GPa is studied in this work.  
171 The steel grade can be used in applications requiring decent strength and  
172 ductility combined with a good resistance to fatigue, for example such as  
173 in engine components. SEM (Scanning Electron Microscopy) with EBSD  
174 (Electron Backscatter Diffraction) detector was used to characterize the mi-  
175 crostructure of the material which is shown in Figure 1a. Conventional sam-  
176 ple preparation techniques was used and the sample was polished prior to

177 SEM and EBSD measurements. The microstructure is lath martensitic. De-  
 178 pending on the material’s manufacturer and process specifics, different de-  
 179 fects may exist in the material. The material can include different individual  
 180 inclusions and clusters of inclusions (typically, e.g., alumina) infused with  
 181 other slag components. Current work focuses on the modeling of plasticity  
 182 and damage in Q&T steel microstructure, and therefore the effect of local  
 183 defects is omitted from the present scope. Thus, the material is assumed  
 184 ideally coherent martensite, as is apparent in Figure 1.

## 185 *2.2. Computational microstructures*

186 To investigate the deformation and damage behavior of the material, two  
 187 polycrystalline microstructures were chosen for the study. The first phase of  
 188 the simulations involves a tessellated representative volume element (RVE)  
 189 that is shown in Figure 1b. It includes a typical prior austenite grain struc-  
 190 ture with 100 grains all assigned to different orientations. The microstruc-  
 191 ture is used for crystal plasticity model parameter identification in Section  
 192 3.2 to capture polycrystalline effects with arbitrarily shaped grains and grain  
 193 boundaries, which can affect crack/damage evolution. The same simplified  
 194 microstructure RVE was utilized also in the simulation of cyclic loading con-  
 195 ditions in Section 3.4. A more detailed and relevant QT microstructure in  
 196 Figure 1c representing lath martensite was extracted as a computationally  
 197 efficient subset from the large EBSD scan in 1a. The finite element mesh  
 198 was based on the EBSD imaging of the microstructure segmented based on  
 199 a presumed discrete orientation distribution capturing the misorientation re-  
 200 lationships as common in martensite. The finite element meshes and models  
 201 were obtained using methodologies described in more detail in [19, 20, 21].  
 202 Advanced characterization techniques exist to track crack growth in 3D [22],  
 203 which could provide more detailed information for the material identification.  
 204 However, the present work focuses on the demonstration of the modeling  
 205 framework.

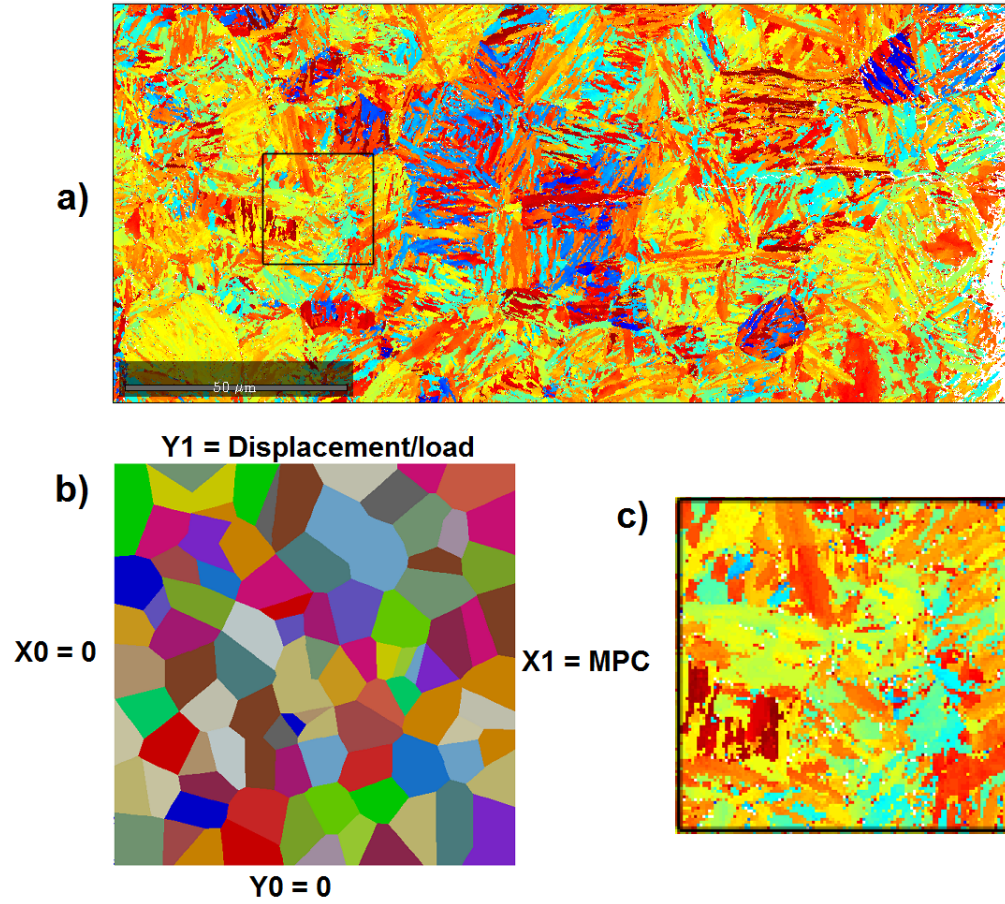


Figure 1: a) Scanning electron microscope image with a orientation map of the Q&T-steel, b) polycrystalline microstructure and c) subtracted computational martensitic microstructure used in the simulations.

206 The computational domains were constrained to retain uniform deformation  
207 at the boundaries (kinematic uniform boundary conditions). The RVEs  
208 were constrained by a direct displacement boundary condition at one side  
209 of the RVE and by a multipoint displacement constraint at the other side.  
210 Displacement controlled loading was applied throughout this work.

### 211 2.3. *Experimental setup*

212 In order to analyze the strain hardening/softening behavior of the Q&T  
213 steel (forged, quenched and tempered), a uniaxial tensile test program with  
214 frequent unloadings to investigate accumulation of damage was established.  
215 The aim of the test sequence is to increase the strain of the sample in incre-  
216 ments up to failure and record the decrease in elastic modulus, the "effective  
217 damage" in a scalar valued form. Approximately one percent increments in  
218 strain were used by first loading the material and then unloading it to almost  
219 zero load, before following with next loading increment. The test was strain  
220 controlled up to two percent of strain and displacement controlled thereafter  
221 with a rate of 1.5 percent per minute. The loading and unloading rate was  
222 carried out at approximately 20 MPa per second and a minimum threshold  
223 of 10 MPa was used at the end of unloading phase.

224 An attempt to evaluate effect of coupled plasticity and damage from the  
225 experiments and to support model parameterization was performed. Besides  
226 only obtaining tensile stress-strain curves with a failure point, the degra-  
227 dation of elasticity could be approximated from the unloading curves as a  
228 function of macroscopic strain. Hence, a classical Lemaitre-type of damage  
229 effect of the elasticity of metals is sought-after in the experiments and there-  
230 fore we denote this experimental procedure as Lemaitre-type tensile test in  
231 the present context (see Lemaitre [23] for background). A 50 mm exten-  
232 someter was used to measure full length deformation of the samples with  
233 12.5 mm range. In addition, DIC system was used to analyze local defor-  
234 mation. A DIC framerate of 2 Hz was chosen. The sample size was 10 mm  
235 in diameter. A spatial resolution of 21.5 pixels per millimeter was chosen.  
236 The speckle size pattern was considered rather macroscopic with 5-20 pixels.  
237 Hence, the DIC system was employed to measure the in-situ diameter of the  
238 test samples. The data was used to correlate between engineering stress and  
239 apparent true stress of the samples by correcting the cross-sectional data  
240 during deformation.

241 It is noteworthy that the error of the technique increases towards higher  
242 localization as the cross-sectional area of the tensile bar does not necessarily



243 retain its symmetricity, which can cause uncertainties at the end of the test  
 244 curve. In spite of this typical possibility, the cross-section of the samples were  
 245 assumed circular throughout the engineering to true stress correction proce-  
 246 dure and provide data the very least from the early stages of necking that can  
 247 be considered valid. Both engineering stress data and corrected true stress  
 248 data were used in the comparison between experiments and simulations.

#### 249 2.4. *Crystal plasticity model*

250 The crystal plasticity model is extended to couple single crystal plasticity  
 251 and micromechanical description of damage, inspired by studies focusing on  
 252 this direct formulation [24, 25, 26]. The following single crystal model was  
 253 implemented in Zset-software and it as used for finite element simulations  
 254 in this work. Deformation gradient is decomposed into elastic and inelastic  
 255 parts:

$$\underline{F} = \underline{F}^e \cdot \underline{F}^{in} \quad (1)$$

256 Inelastic deformation gradient  $\underline{F}^{in}$  features both plastic deformation and  
 257 damage contribution. In the present model, a concept of inelastic damage  
 258 strain is adopted [24, 25], to distinct contributions of plasticity and damage  
 259 in single crystal formulation. It also follows that inelastic strain rate  $\dot{\underline{F}} \cdot \dot{\underline{F}}^{-1}$   
 260 is not always traceless because occurrence of volume change that is associated  
 261 with damage.

$$\dot{\underline{\epsilon}}^{in} = \dot{\underline{\epsilon}}^p + \dot{\underline{\epsilon}}^d \quad (2)$$

262 The model has three important features. Firstly, this approach allows for  
 263 the damage to operate on crystallographic planes and to produce damage  
 264 strain according to specific assigned mechanism, i.e., cleavage damage strain  
 265 operating on the cleavage planes of the single crystals. The formulation also  
 266 quantifies that damage is not exactly bounded to be a variable between zero  
 267 and one in contrast to conventional continuum damage approaches. Secondly,  
 268 more recent development of similar strain based single crystal damage model  
 269 [25] demonstrated that besides the model describes damage evolution and  
 270 related material softening, it also aims to describe crack closure. Thirdly,  
 271 the model features a two-way coupling between damage and plasticity, which  
 272 means that damage allows plastic localization and on the other hand plastic  
 273 localization makes it easier to nucleate and grow damage (cracks).

274 The plastic deformation of the polycrystalline material, without any dam-  
 275 age, is described with a rate dependent crystal plasticity model. Dislocation  
 276 slip is considered to carry out plasticity in martensitic steels. The slip rate  
 277 is described by phenomenological slip rate formulation.

$$\dot{\gamma}^s = \dot{\nu} \text{sign}(\tau^s) = \left\langle \frac{|\tau^s - x^s| - r^s - \tau_0}{K} \right\rangle^n \text{sign}(\tau^s - x^s) \quad (3)$$

278 where material parameters  $K$  and  $n$  characterize the viscosity,  $\tau^s$  is the  
 279 current resolved shear stress in a system  $s$ ,  $\tau_0$  is the initial shear resistance of  
 280 slip system families  $\{110\} < 111 >$  and  $\{112\} < 111 >$ , and  $r^s$  is the isotropic  
 281 hardening variable. For simplicity, the slip resistance is assumed the same  
 282 for both slip families. The resolved shear stress is computed utilizing the  
 283 Mandel stress  $\underline{M}$  in the intermediate configuration, so that:

$$\tau^s = \underline{M} : \underline{m}^s = (\underline{C}^e \cdot \underline{S}^e) : \underline{m}^s = (\underline{C}^e \cdot (\underline{\Lambda} : \underline{E}_{gl})) : \underline{m}^s \quad (4)$$

284 where  $\underline{C}^e$  is the Cauchy-Green tensor,  $\underline{E}_{gl}$  is the Green-Lagrange strain  
 285 tensor, and  $\underline{\Lambda}$  is the elastic stiffness tensor, and  $\cdot$  The second Piola-Kirchhoff  
 286 stress tensor  $\underline{S}^e$  and the Mandel stress tensor  $\underline{M}$  are connected in the usual  
 287 manner by  $\underline{M} = \underline{C}^e \cdot \underline{S}^e = (\underline{F}_e^T \cdot \underline{F}_e) \cdot (\underline{\Lambda} : \underline{E}_{gl})$ , where  $\underline{F}_e$  is the elastic part  
 288 of the deformation gradient.

289 The evolution of the kinematic hardening parameter  $x^s$  is computed with  
 290 the following evolution equation.

$$x^s = c\alpha_s ; \quad \dot{\alpha}_s = (\text{sign}(\tau^s - x^s) - d\alpha^s)\dot{\nu}^s ; \quad \nu^s = \int_0^t |\dot{\gamma}^s| \quad (5)$$

291 where the rate of cumulative plastic slip  $\dot{\nu}^s$  controls the evolution of the  
 292 kinematic hardening, and coefficients  $c$  and  $d$  are related to its intensity.

293 Both of the included slip families  $\{110\} < 111 >$  and  $\{112\} < 111 >$  share  
 294 the same flow rule. However, the viscous parameters  $K$  and  $n$  can be different  
 295 in addition to dissimilar critical resolved shear stresses for the activation, if  
 296 necessary, and the parameters can be identified from the experimental data  
 297 and characterization if such are available. All of the slip systems in a slip  
 298 family are assumed to have the same critical value for activation for simplicity.

299 Sabnis et al. [25] suggested a coupling between plasticity and damage, in  
 300 which damage is driven by plasticity. Accordingly, apart from conventional  
 301 treatments of hardening formulations, we consider that strain accumulation  
 302 decreases the local cleavage resistance, and thus producing a link between

303 plastic deformation and susceptibility to damage. Free energy density func-  
 304 tion is formulated to consist elastic, plastic and damage contributions, re-  
 305 spectively.

$$\rho\Psi(\underline{\epsilon}^e, \nu, d) = \rho\Psi^e(\underline{\epsilon}^e) + \rho\Psi^p(\nu) + \rho\Psi^d(\nu, d) \quad (6)$$

306 The additional contribution of damage in the free energy potential con-  
 307 tribution takes a form:

$$\rho\Psi^d = Y_0 d + \frac{H}{2}(d + \beta\nu)^2 \quad (7)$$

308 where  $Y_0$  is the initial cleavage resistance, described in detail below.  $H$   
 309 is a softening parameter and  $\beta$  plasticity-damage coupling parameter, and  $d$   
 310 is the cumulative damage, the rate of which is defined by crack opening and  
 311 two shear mechanisms as:

$$\dot{d} = \sum_{s=1}^{N_{damage}} |\dot{\delta}_c^s| + |\dot{\delta}_1^s| + |\dot{\delta}_2^s| \quad (8)$$

312 Isotropic hardening is produced for all systems, including contributions  
 313 from self and latent hardening. The present expression, in Equation 9, con-  
 314 sideres the interactions between the two slip families. The coupling between  
 315 plasticity and damage relates isotropic hardening and softening arising from  
 316 the contribution of damage, which is obtained by partial derivation of the  
 317 free energy functional:

$$r^s = \rho \frac{\partial \Psi}{\partial \nu} = \tau_0 + Q \sum_r H_{rs} \{1 - \exp(-bv^r)\} + H\beta^2\nu + H\beta d \quad (9)$$

318 The effect of damage does not diminish even if the crack becomes closed,  
 319 i.e., crack opening strain is zero, because plastic softening coming from dam-  
 320 age is controlled by cumulative damage form presented in Equation 8. A  
 321 case where high damage strains can introduce excessive softening in the crack  
 322 closed state can realize after severe plasticity and damage. Therefore reg-  
 323 ularization of damage softening may be required that simply saturates the  
 324 damage softening contribution to a maximum softening value. In general, the  
 325 softening-coupling modulus  $H$  describes the softening effect of micro-cracking  
 326 in the material, which accelerates both damage and plasticity.

327 The interaction matrix  $H_{rs}$  is adopted from the work of Hoc and Forest  
 328 [27] on a BCC material. The form of the interaction matrix showed good

Table 1: The coefficients for the interaction matrix in BCC crystals

Plane	$\{110\} \cap \{110\}$	$\{110\} \cap \{112\}$	$\{112\} \cap \{112\}$
Same	$a_0 (h_8)$		$k_{s0}a_0 (h_1)$
Collinear	$k_1a_0 (h_2)$	$k_{p1}a_0 (h_3)$	$k_{s0}k_1a_0 (h_6)$
No collinearity	$k_2k_1a_0 (h_4)$	$k_{p2}k_{p1}a_0 (h_5)$	$k_{s0}k_2k_1a_0 (h_7)$

agreement with experiments involving complex loading paths that may occur in loading various loading conditions including multi-axial fatigue. The number of independent coefficients is reduced to eight within the 24x24 interaction matrix, i.e., classifying the slip systems belonging to the same slip family, to a collinear system or to a non-collinear system. These interactions are presented in Table 1. Kinematic hardening parameterization is obtained from earlier cyclic tensile tests, while calibration of the isotropic response is describe in detail within the results section.

We retain the assumption that interaction between planes  $\{110\}$  is smaller than between planes  $\{112\}$  [27].

The rate of damage process (crack growth) includes crack opening rate and modes II/III crack shearing rates on all given cleavage planes:

$$\dot{\epsilon}^d = \sum_{s=1}^{N_{damage}} \dot{\delta}_c^s \underline{n}_d^s \otimes \underline{n}_d^s + \dot{\delta}_1^s \underline{n}_d^s \otimes \underline{l}_{d1}^s + \dot{\delta}_2^s \underline{n}_d^s \otimes \underline{l}_{d2}^s \quad (10)$$

where  $\dot{\delta}_c^s$ ,  $\dot{\delta}_1^s$ ,  $\dot{\delta}_2^s$  are the strain rates for mode I, mode II, and mode III crack growth, respectively.  $N_{damage}$  denotes the number of damage planes, which are fixed crystallographic cleavage planes for given crystal structure. Cleavage damage is represented by the opening  $\delta_c$  of crystallographic cleavage planes with the normal vector  $\underline{n}^s$  and other damage systems must be introduced for the in-plane accommodation along orthogonal directions  $\underline{l}_{d1}^s$  and  $\underline{l}_{d2}^s$  once cleavage opening has initiated. The model considers the BCC cleavage planes of type  $\{100\}$  for martensitic steels, which have been experimentally observed to have significant effect in the crack initiation and propagation for martensitic and ferritic steels [28, 29]. The evolution of the opening rate is given by:

$$\dot{\delta}_c^s = \left\langle \frac{|\sigma_{dc}| - Y_c^s}{K_d} \right\rangle^{n_d} \text{sign}(\tau_{dc}) \quad \text{with} \quad \sigma_{dc} = \underline{n}_d^s \cdot \underline{M} \cdot \underline{n}_d^s \quad (11)$$

352 The opening rate  $\dot{\delta}_c^s$  is driven by the normal component acting on cleavage  
 353 planes whenever normal stress  $\sigma_{dc}$  exceeds the cleavage crack resistance  $Y_c^s$ .  
 354 Initially only positive stress  $\sigma_{dc}$  can initiate cleavage. After crack opening,  
 355 the crack is allowed to close if the opening stress is negative with a constraint  
 356 that  $\delta_c^s \geq 0$ , which renders a crack closure scheme possible for the model.

357 The rate of Mode II and III shear mechanisms operating in conjunction  
 358 with cleavage crack opening mode I are given by a similar rate dependent  
 359 formulation:

$$\dot{\delta}_{1,2}^s = \left\langle \frac{|\tau_{di}| - Y_i^s}{K_d} \right\rangle^{n_d} \text{sign}(\tau_{di}) \quad \text{with} \quad \tau_{di} = \underline{n}_d \cdot \underline{M} \cdot \underline{l}_{di}^s \quad (12)$$

360 where shear stress  $\tau_{di}$  drives the damage shear mechanisms when shear  
 361 resistance  $Y_i^s$  is overcome. In both crack opening and shearing mechanisms,  
 362  $K_d$  and  $n_d$  are viscous parameters.

363 The damage resistance of the material consists material's initial cleavage  
 364 resistance  $Y_0$  that can vary within the microstructure depending on the de-  
 365 fect population and other nanostructural material characteristics. The initial  
 366 resistance can be different for crack opening and shear mechanisms, but it is  
 367 assumed to be the same in the present context. Obtaining a partial deriva-  
 368 tive with respect to damage of the free energy function yields opening/shear  
 369 resistance for cleavage:

$$Y = \rho \frac{\partial \Psi}{\partial d} = Y_0 + Hd + H\beta\nu \quad (13)$$

370 Cumulative damage  $d$  has direct softening effect on the residual cleavage  
 371 resistance, while plastic localization softens the material via coupling factor  
 372  $\beta$ . It is worth noting that the final cleavage resistance is not exactly zero  
 373 because of viscous components arising from visco-plastic and visco-damage  
 374 formulation via parameters  $K$  and  $n$  as well as residual component of cleavage  
 375 resistance can be assigned to improve convergence.

376 A difference exists between previous developments of this type of damage  
 377 framework in small deformations [24, 25]. The present model omits mi-  
 378 cromorphic damage regularization in the formulation to drastically decrease  
 379 overhead of the computations of large microstructures in the absence of new  
 380 degree(s) of freedom and complexity of the microdamage concept, when fi-  
 381 nite strains are considered. This simplification makes the damage zones more  
 382 diffuse in the present model and somewhat mesh size dependent in terms of  
 383 damage width, which is inherent characteristics of finite elements. A better

384 description of mesh-independent damage is achievable by introducing damage  
385 size dependent regularization.

### 386 **3. Results**

#### 387 *3.1. Deformation and damage response*

388 The model response is demonstrated in cyclic uniaxial test in Figure 2 per-  
389 formed on a single element and single crystal case with a displacement control  
390 loading, demonstrating also the behavior of the smeared contact approach.  
391 The material is first loaded in tensile direction to cause severe plasticity and  
392 strain hardening, and then reversed to compression to further allow plastic-  
393 ity to develop. This cycle decreases the cleavage resistance of the material  
394 through damage-plasticity coupling and reversion back to tension shows initi-  
395 ation of cleavage. The crack opening occurs and the material quickly softens  
396 further promoting crack opening. The loading is again reversed to compres-  
397 sion and crack closure begins with diminishing stress resistance until crack  
398 opening strain reaches zero and crack is considered closed. At this point,  
399 the material regains some of its original strength, which itself is degraded  
400 according to Equation 9. Finally if the loading is reversed to tension, the  
401 material has very small resistance against cleavage when being already highly  
402 damaged and the crack opens with ease.

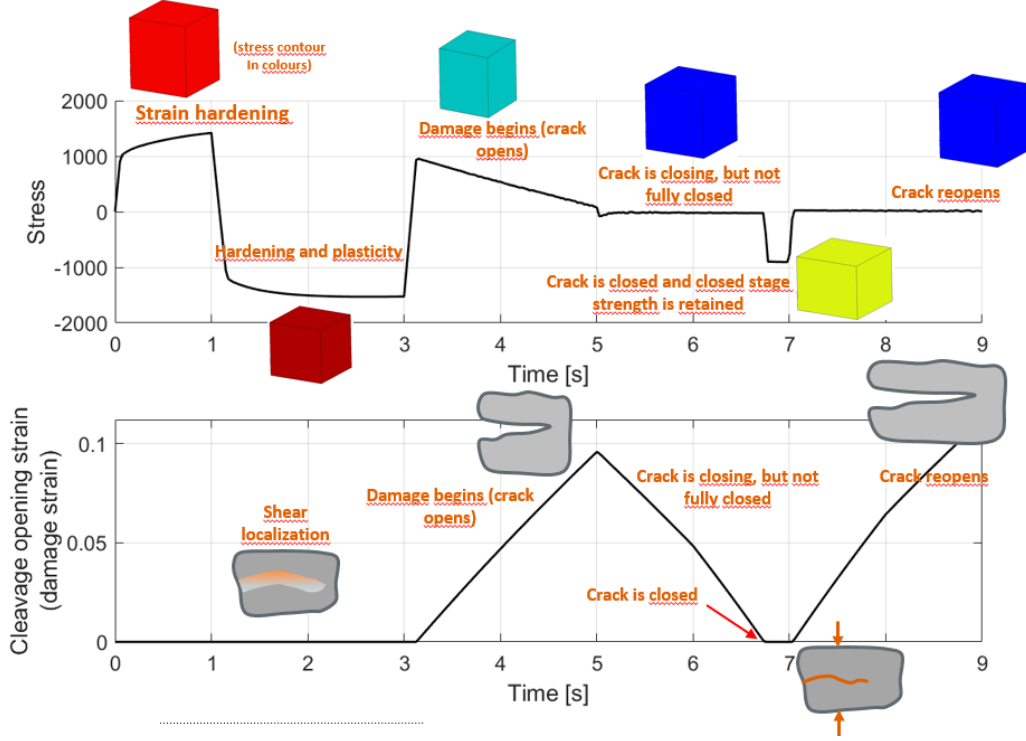


Figure 2: Plasticity-damage model response and crack opening-closure behavior. Element colors describe stress intensity with red being high stress and blue zero stress.

### 3.2. Damage model parameterization

The stress-strain response of the constitutive model is affected by the coupling between plasticity and damage, which means that hardening part of the stress-strain curve is greatly influenced by the softening and coupling parameterization of the model. Therefore the following presents some cases where damage softening and coupling parameterization is altered, while dislocation slip hardening parameters were kept constant. Figure 3 shows the effect of softening parameter  $H$  to the stress-strain and damage responses during an uniaxial tensile test. A polycrystal microstructure with 100 grains as presented in Figure 1 was used in the analyses. The stress-strain, damage and equivalent plastic strain values were obtained by averaging over the whole computational microstructure domain. Low softening modulus value retards the occurrence of damage to relatively high strains and hardening of the material prevails with this parametrization. High softening values, such

417 as -3000 MPa, accelerate the degradation of material strength and cleavage  
 418 resistance, leading to early damage initiation and spontaneous spreading of  
 419 the damage. The growth rate ( $\partial d / \partial \epsilon$ ) of damage can be controlled also with  
 420 the softening modulus, which is evident when comparing the -3000 MPa and  
 421 -1750 MPa cases. The coupling parameter was kept constant at 0.265 during  
 422 the simulations.

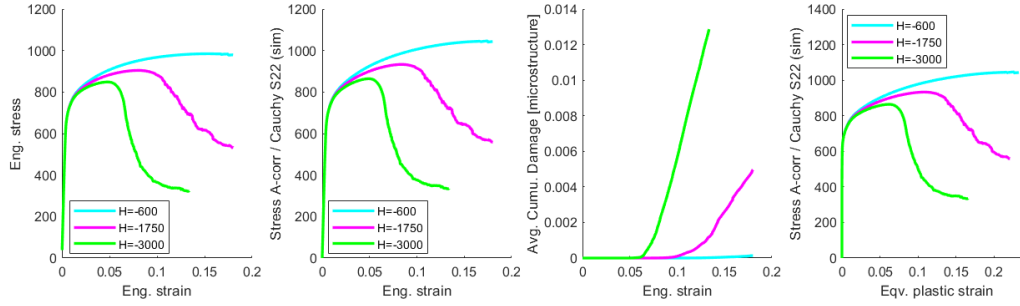


Figure 3: Effect of material softening to stress-strain-damage response

423 The coupling parameter has similar retarding or accelerating effect to the  
 424 damage initiation as well as on the softening of stress-strain response. The  
 425 effect of coupling factor is shown in Figure 4. The softening parameter was  
 426 kept constant at -1750 MPa during the simulations. A physical interpretation  
 427 of softening parameter could be considered as a ductility parameter  
 428 which represents material's capability to continue plastic deformation even  
 429 at high localization strains and in the presence of damage, while it also has  
 430 great influence on the effective rate how rapidly damage (cracking) occurs.  
 431 The coupling parameter then is a similar measure of how easily and quickly  
 432 existing damage leads to further plastic strain localization and on the other  
 433 hand what is the contribution of localized plastic strain to damage nucleation  
 434 and growth processes.



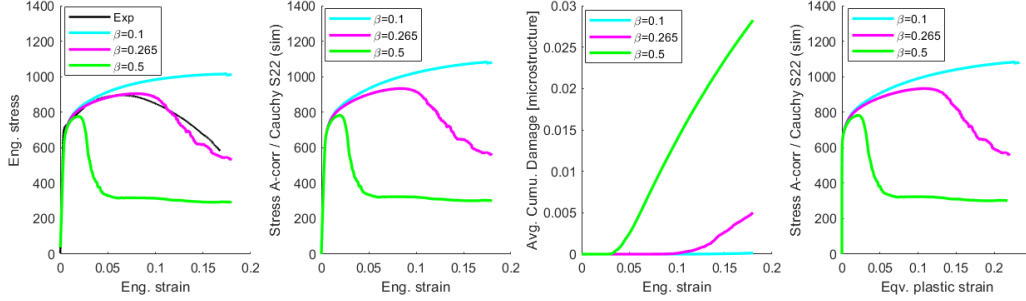


Figure 4: Effect of damage-plasticity coupling.

435 In both of the previous sensitivity analysis cases, it is clear that the bal-  
 436 ance between each parameter and their effect on stress-strain response is not  
 437 straightforwardly identifiable. A principal assumption is made that coupling  
 438 parameter would lie in the moderate region, e.g., 0.2-0.4. Larger values would  
 439 lead to very early damage initiation and material would experience brittle  
 440 like behavior, which would be more typical for untempered martensite. Low  
 441 values favor plasticity driven damage process, where extensive strain local-  
 442 ization would be required, making it more suitable for ductile materials. A  
 443 follow up to current work is to further investigate the dependencies utilizing  
 444 molecular and dislocation dynamics modeling as well as in-situ microstruc-  
 445 ture resolution tensile testing. Although these methodologies have challenges  
 446 with respect to spatial and temporal scales with respect to modeling poly-  
 447 crystals, they are still foreseen to yield additional data to support solely  
 448 experimentally driven calibration approaches.

449 Figure 5 compares the Lemaitre-type of uniaxial tensile loading-unloading  
 450 experimental results with the simulated response of the model. Figure 5a  
 451 shows the engineering values. Figure 5b shows an expression of true stress  
 452 state against engineering strain. The true stress for experiments was obtained  
 453 by correcting engineering stress with the instantaneous cross-sectional area  
 454 that was defined as equivalent area in the necking region by measuring the  
 455 diameter of the tensile bar with DIC system in the two dimensional analysis,  
 456 as explained in Section 2.3. The corrected experimental true stress values  
 457 were thus compared against an axial component of the Cauchy stress tensor  
 458 of the simulations, which is the stress component in the loading direction.  
 459 The results have good agreement throughout the strain range. The dam-  
 460 age dominated region beyond 0.1 of strain shows largest difference because  
 461 of the uncertainties related to damage evolution of the material as well as

462 in the definition of the cross-sectional area used to compute the experimen-  
463 tal true stress. Furthermore, it is also evident that the material's softening  
464 behavior is not constant in the three experimental cases and the scatter is  
465 relatively large when compared to the more constant hardening part of the  
466 stress-strain curve. A parametric simulation study revealed that it is pos-  
467 sible to adjust the shape of the curve in the damage dominated region by  
468 varying the softening and coupling parameters ( $\pm 200$  MPa, and  $\pm 0.02$ , re-  
469 spectively). Yet, given experimentally observed scatter the parameterization  
470 can be considered acceptable due to the results residing within the experi-  
471 mentally observed scatter. Figure 5c shows the true stress as a function of  
472 inelastic strain. The inelastic strain was defined as the plastic part of total  
473 strain measured during the experiment, i.e., total strain comprises elastic  
474 and inelastic strain contributions. No direct distinction between plasticity  
475 and damage could be made in the experiments. The value from simulations  
476 corresponds to the average of inelastic strain over the microstructure. The  
477 actual local strains are much larger than the average strain of the sample,  
478 which explains the difference between the experiment and simulations.

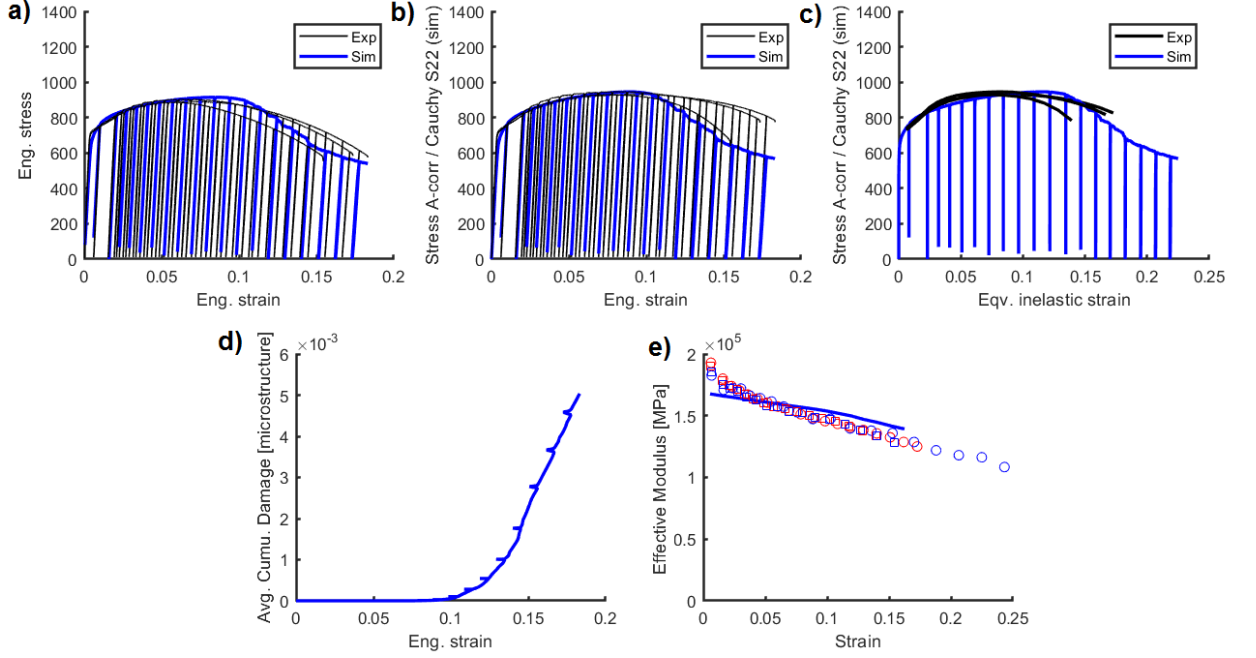


Figure 5: Experimental loading-unloading tensile experiment used for model parameter fitting. a) Engineering stress and engineering strain, b) simulated true axial stress and experiment based area corrected stress c) true stress and inelastic strain (simulated and exp. extensometer), d) cumulative damage over the RVE microstructure, e) effective elasticity measured from unloading curves of experiments (extensometer and DIC) and simulations.

479 The elastic stiffness of the material is not explicitly degraded in the model  
 480 by bounding damage variable to instantaneous elastic moduli as a function  
 481 of microcracking. It is assumed that plasticity dominates over elasticity during  
 482 hardening and also during the rapidly developing cracking process. One  
 483 fundamental aim of the Lemaitre-type damage experiment is to approximate  
 484 the change in the effective stiffness of the material, i.e., damage degrading  
 485 the elasticity of the material. In the experiments, the effective elastic modulus  
 486 was determined from the unloading curves after each loading increment  
 487 using extensometer data and DIC measurements. Figure 5d shows the accumulation  
 488 of damage in the microstructure as a function of cumulatively  
 489 increasing strain during the loading-unloading sequence. Once damage is  
 490 initiated, its effect quickly increases. During unloading, cumulative damage  
 491 increases slightly when the cracks start to close under compression. Figure

492 5e plots the degradation of elasticity as a function of equivalent inelastic  
493 strain. It can be seen that the material experiences notable loss of stiffness  
494 when inelastic strain increases. The simulation model can correlate similar  
495 degradation in elasticity to a certain degree, but because of lacking full cou-  
496 pling between damage and elastic moduli, a deviation in the results is to be  
497 expected. Simulations showed that damage itself can largely accurately ac-  
498 count for the loss of stiffness, but in such cases the hardening of the material  
499 must to be set to unrealistic low.

500 Figure 5e shows the evolution of damage during the simulation range.  
501 The shape of the curve shows that the repeated loading sequences first ini-  
502 tiate and grow damage slowly, i.e., interpretable as small microcrack growth  
503 dominated region in the damage process. Figure 6a visualizes total inelastic  
504 strain contributed by plasticity and damage at 8 % of macroscopic strain.  
505 After the initiation, the rate of damage and associated plasticity rapidly in-  
506 crease between 0.11 and 0.18 of macroscopic strain. Figure 6b,c. shows these  
507 two stages of inelastic strain localization. Damage and plasticity continues  
508 to localize between two nucleation sites, which both are in the vicinity of  
509 intersection of multiple grains in the present simulation case. The growth  
510 mechanism of the cracks include crack/damage growth at the grain bound-  
511 aries of different grains (orientations) as well as the interior regions of the  
512 grains, finally extending from short local intra-grain cracks to inter-grain  
513 polycrystal scale cracks.

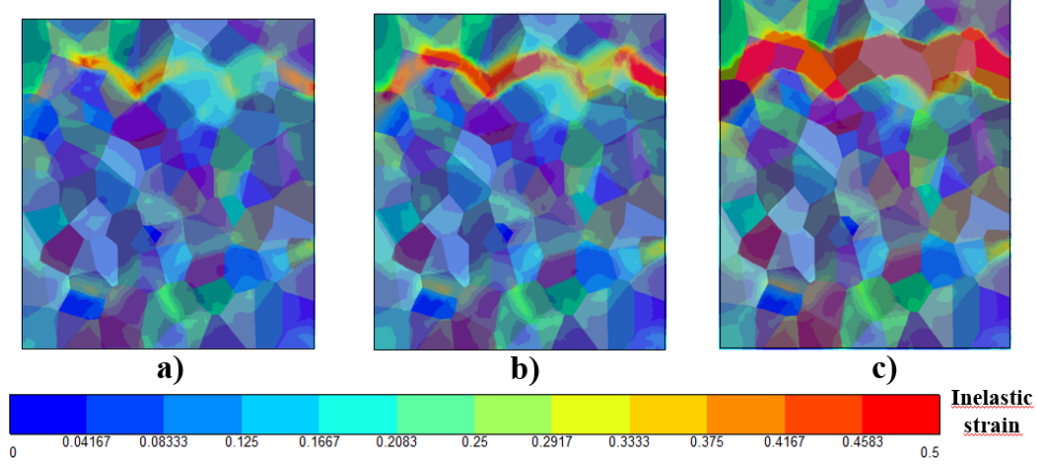


Figure 6: Inelastic strain overlay on the microstructure, at the macroscopic engineering strain of a) 9 %, b) 11 %, c) 16 % during the uniaxial cyclic loading-unloading test.

Utilizing the resulting material model calibration Figure 7 presents the experimental and simulated results of a conventional tensile test until failure. The model agrees well with the measured data with the present parameter set. Table 2 lists the used material parameters. The initial cleavage resistance was set to 1900 MPa after a parametric study. A larger value requires larger softening and coupling factors to better correlate the retarded initiation of softening in the stress-strain curve, while lower value can prematurely initiate damage at lower strains and lead to earlier failure than was experimentally observed. The curve with no active damage shows that the flow stress of the material is higher in the absence of damage and related coupling effects. The isotropic and kinematic viscoplastic response is obtained by way of utilizing cyclic stress-strain curves along with homogenization and self consistent methods, as a part of work presented in [1].

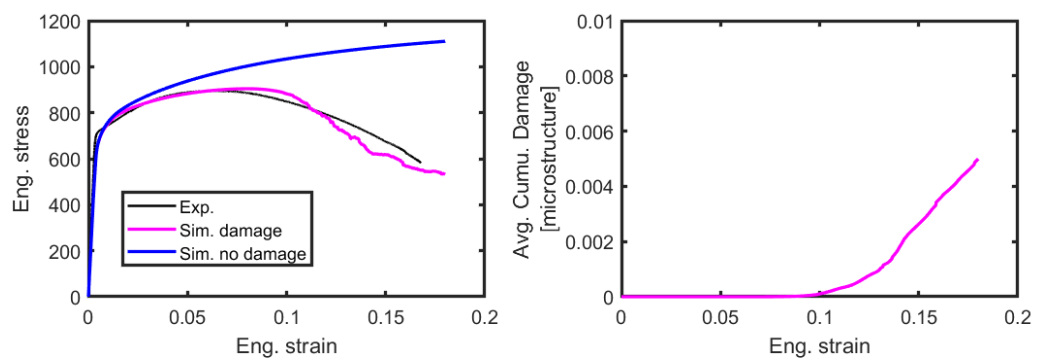


Figure 7: Experimental and simulated tensile test

Table 2: Model parameters for single crystal martensitic Q&T steel.

Parameter	Value	Unit
Elastic constants [30]		
$C_{11}$	198 000	[MPa]
$C_{12}$	78 125	[MPa]
$C_{44}$	70 700	[MPa]
Slip parameters		
$\tau_0^s$	169.0	[MPa]
$K$	169.0	[MPa.s <sup>1/n</sup> ]
$n$	50.0	-
$b$	19.0	-
$Q$	37.0	[MPa]
$C$	2000.0	[MPa]
$D$	200.0	[MPa]
$h_1$	1.3	-
$h_2$	1.0	-
$h_3$	1.05	-
$h_4$	1.15	-
$h_5$	1.1025	-
$h_6$	1.3	-
$h_7$	1.4950	-
$h_8$	1.0	-
Damage parameters		
$Y_c^s$	1900.0	[MPa]
$Y_1^s$	1900.0	[MPa]
$Y_2^s$	1900.0	[MPa]
$K_d$	5.0	[MPa.s <sup>1/n</sup> ]
$n_d$	2.0	-
$H$	-1750	[MPa]
$\beta$	0.265	-

### 527 3.3. Strain localization and damage in lath martensite

528 Figure 8 shows stress-strain curve with damage accumulation for the lath  
 529 martensitic microstructure in Figure 1c subjected to uniaxial tension. The  
 530 simulation reproduces a similar strain hardening phase up to about 9-10 %  
 531 of strain, that was for observed for the polycrystalline microstructure with  
 532 prior austenite grains only in Figure 6. However, a radical accumulation of  
 533 damage occurs when axial strain increases, contributing strongly on mate-  
 534 rial's softening. Figure 9 represents equivalent stress contours and inelastic  
 535 strain distribution in the microstructure at three strain levels.

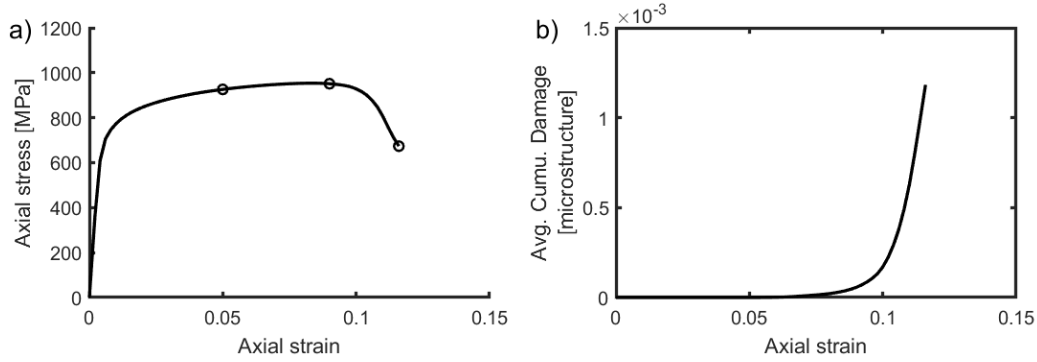


Figure 8: a) Stress-strain curve of lath martensitic microstructure, and b) cumulative damage during applied tensile loading over the entire microstructure.

536 Stress concentrations are observed in respective different hard and soft  
 537 orientations within the polycrystalline fine lath martensite. Strain hardening  
 538 increases the overall strength of the microstructure RVE. At first with 5%  
 539 of axial strain, local strains are accommodated in favorable sites, but the  
 540 intensity of straining remains relatively low. However, when overall strain of  
 541 the microstructure increases, extensive strain localization tends to overcome  
 542 strain hardening effects and revert to localized strain softening. Strain lo-  
 543 calization is a driving force for damage that further accumulates plasticity,  
 544 which finally leads to distinctive local decrease in flow stress and identifiable  
 545 crack path forms, that is visible in Figure 9a,f. The crack path has evolved  
 546 to a point, where further restricting boundaries are required to inhibit its  
 547 growth, and a transition from micro-cracking to short crack growth could be  
 548 interpreted as consequence, i.e., from submicron size to tens of microns size  
 549 in the case of the presently studied microstructure.



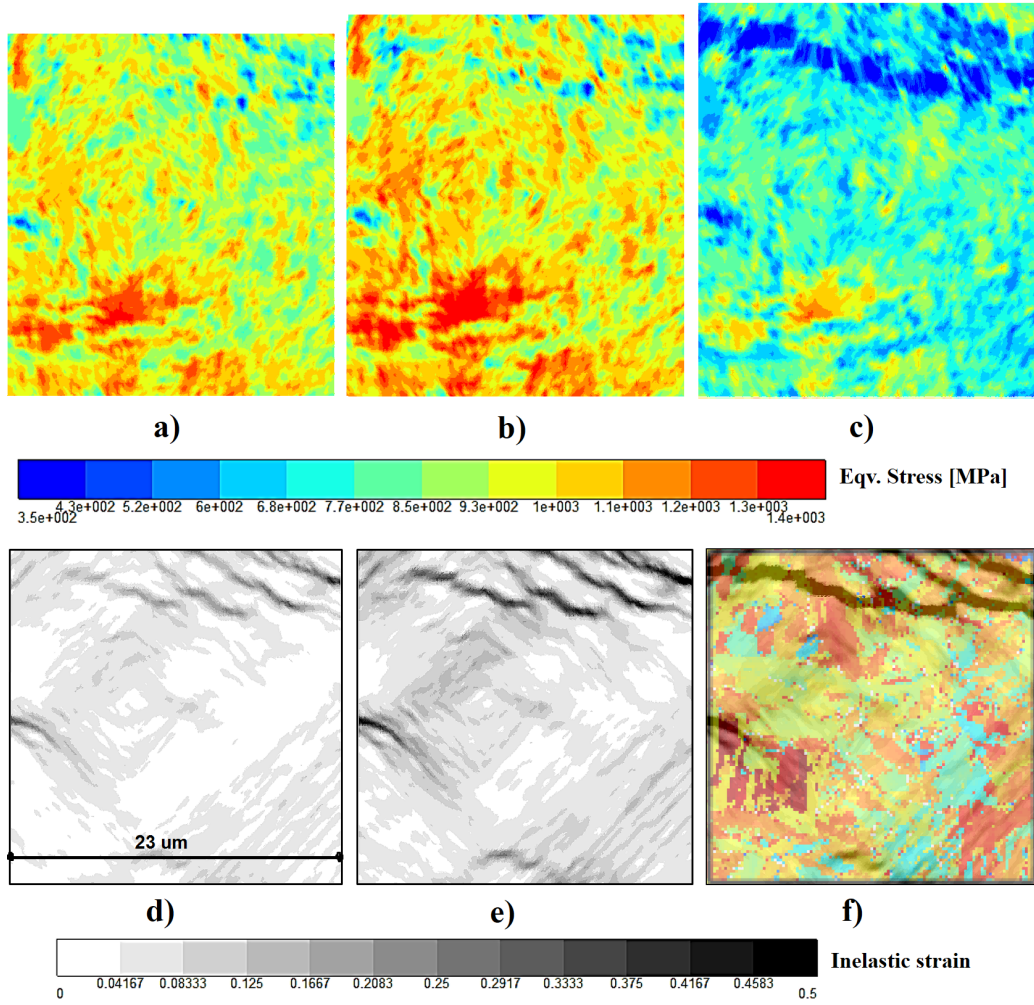


Figure 9: Equivalent stress contours a)-c), and inelastic strain plotted over undeformed shape d)-f) of lath martensite microstructure at different axial strains, 0.05, 0.08 and 0.116, respectively.

### 550 3.4. Fatigue application

551 The effect of damage in cyclic conditions is key aspect in order to capture  
552 damage nucleation and crack growth under low, high and very high cycle  
553 fatigue. Figure 10a,b shows cyclic stress-strain loops for plasticity-only and  
554 plasticity-damage cases. The stress-strain curves are generated by averaging  
555 over all of the elements in the full field microstructural mesh, which thus  
556 involves elements beyond damaged sites only. A large absolute strain am-  
557 plitude of  $\Delta\varepsilon = 2\%$  was chosen to accelerate strain localization and damage  
558 evolution.

559 A typical cyclic isotropic-kinematic hardening behavior is observed in the  
560 absence of damage and related self-softening of the material. Figure 10b  
561 shows the softening and effect of damage in the hysteresis. During a phase  
562 when plasticity prevails, i.e., damage is not present, the material experi-  
563 ences strain hardening. This phase is similar to the Lemaitre-type loading-  
564 unloading experiments and simulations in Figure 5 before cyclic softening  
565 and damage effects become relevant. The following damage phase agitates  
566 notable softening in the stress-strain loops. This is further pronounced by  
567 the tensile-compression asymmetry of the model, with lower flow stress in  
568 tension and higher absolute flow stress in compression with a character to  
569 introduce apparent hardening behavior towards higher compressive strain.  
570 The softening itself is more limited when micro-cracking (damage initiation)  
571 is dominant. However, the high strain amplitude of the simulation quickly  
572 increases the effect of damage, resulting in the observed dramatic softening  
573 of the material within two complete cycles.

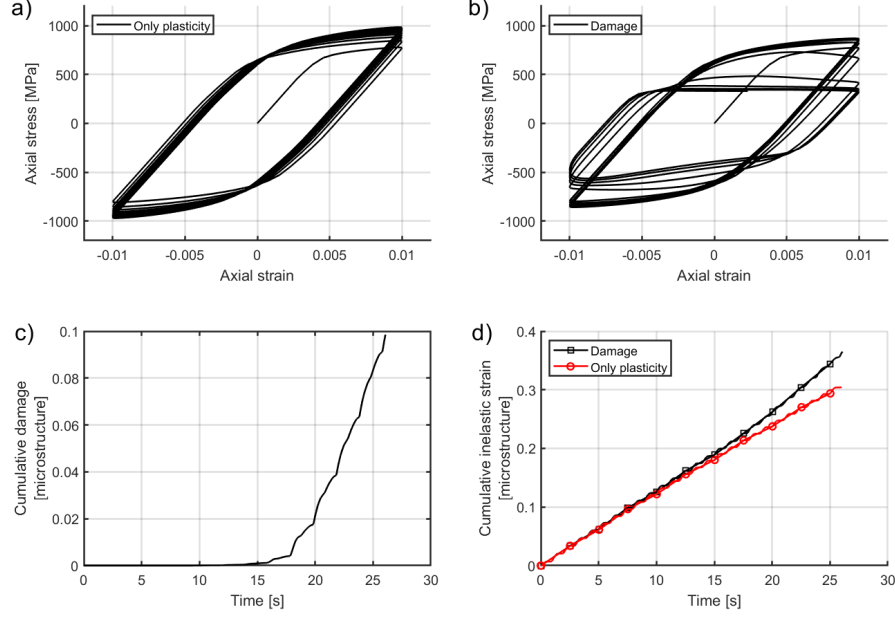


Figure 10: Simulated cyclic mechanical response of the model, stress-strain loops for a) plasticity-only case, b) damage case, and c) cumulative damage for damage case, d) cumulative inelastic strain for both plasticity-only and damage cases. Model is strain controlled with a loading ratio (R) of -1 (fully reversed loading).

574 The material's flow stress in tension after effective crack growth (dam-  
 575 age) in the microstructure is largely restricted by the strain localization and  
 576 crack opening flow of the damaged regions. After repeated opening-closing  
 577 sequences, the material has diminishing small resistance against crack open-  
 578 ing as already visualized in Figure 2. However, the regions that are only  
 579 partially damaged and can further inhabit plasticity and damage, or regions  
 580 with non-opened cracks, still carry out resistance against deformation. These  
 581 contributions lead to a stable 300-350 MPa flow resistance also at tensile axial  
 582 strains over the whole microstructure. In addition, a residual strength for slip  
 583 resistance of 5 % of the initial slip resistance in highly damaged conditions  
 584 was assumed for numerical stability, i.e.,  $\tau_{res}^s = \max(0.05 \cdot \tau_0, \tau^s)$ .

585 Figure 11b shows how compressive axial strains promote crack closure.  
 586 However, a crucial observation can be made, the crack closure does not either  
 587 take place immediately or completely after load reversal to compression. At

588 local scale, some regions of the microstructure can undergo crack closure, but  
 589 other regions remain partially open and operate as regions with very small  
 590 flow stress until cracks become closed. This phenomenon causes a hardening  
 591 effect in the flow stress observed in Figure 10b when more and more cracks  
 592 close with the increasing compressive axial strain towards 1 %.

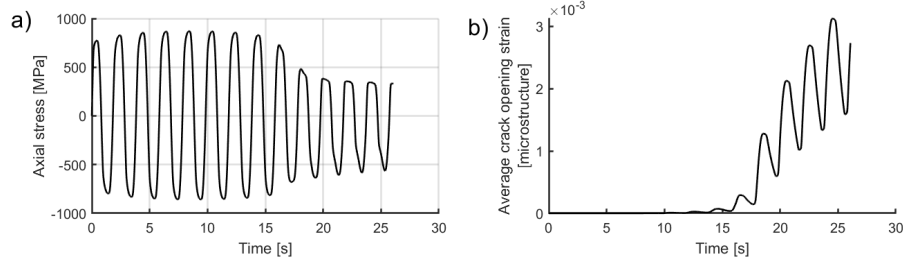


Figure 11: a) Cyclic stress response of the model, and b) average opening strain over the whole microstructure of all cleavage systems as a function of time.

593 Figure 12a compares accumulation of inelastic strain for plasticity-only  
 594 and plasticity-damage cases. Strain localization precedes the damage ini-  
 595 tiation and damage dominant behavior of the steel microstructure. It is  
 596 expected that the microstructural features drive the strain localization, typi-  
 597 cally triple points and loading wise soft orientations in the present case. The  
 598 extent of plastic localization depends much on whether material model can  
 599 promote local softening through cyclic softening and microdamage.

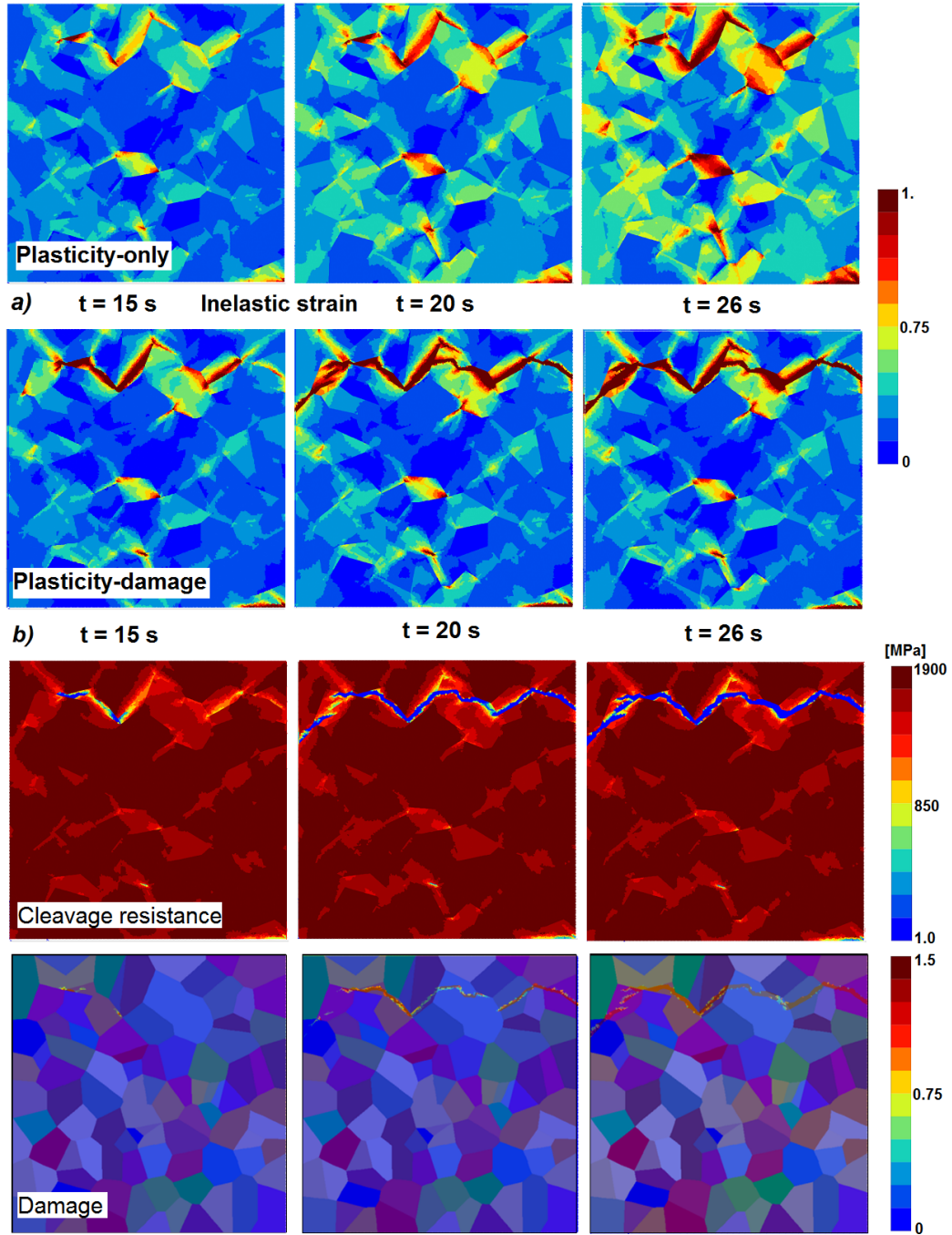


Figure 12: a) Inelastic strain accumulation at different time steps for plasticity-only and plasticity-damage cases, b) cleavage resistance and cumulative damage strain at different time steps.

600 Plasticity-only case distributes inelastic strain in various regions in the  
601 studied RVE, while plasticity-damage case has the capability to drastically  
602 sharpen localization, resulting in shear band type of failure zone. Some mesh  
603 dependency is expected in the width of the band because of the absence of  
604 regularization damage term, as already pointed out with respect to model  
605 derivation. Nevertheless, the probability to nucleate crack opening in the  
606 plasticity rich zones, whether being narrow or wide, is increasing quickly as  
607 a function of decreasing crack opening cleavage resistance shown in Figure  
608 12b. One can use the model to indicate that the material has completely  
609 failed at the point where resistance against damage is vanishing small, that  
610 was presently chosen to a threshold value of 1 MPa.

#### 611 4. Discussion

612 The extension of the introduced micromechanical damage concept to crys-  
613 tal plasticity modeling of plastic slip enabled the coupling between deforma-  
614 tion and cleavage type measure of microstructure scale damage. The im-  
615 portance of strain localization preceding material failure in martensitic steels  
616 can be investigated in arbitrary loading conditions because of the finite strain  
617 implementation of the present model, including the concept's capabilities to  
618 capture deformation and failure processes quite generally in polycrystalline  
619 microstructures. It follows that the damage path is much dictated by the  
620 following strain localization bands, as shown in Figure 12, as is expected.  
621 The susceptibility of the material to invoke strain localization, thus, has a  
622 significant role in terms of developing damage in the first place, and even  
623 more importantly in the transition from micro-cracking to microstructurally  
624 short crack state and finally to long more microstructure insensitive crack  
625 formation. The progression of this micromechanism ultimately defining the  
626 material's performance against failure and also links micromechanics and  
627 conventional fracture mechanics.

628 The model behavior can be adjusted to represent brittle failure character  
629 such as could be expected from untempered martensite or even to direct  
630 fracture mechanical analyses. Alternatively, large local plastic strains can be  
631 allowed for pronounced plasticity driven damage process, e.g., such as those  
632 of tempered martensite steel grades. The effectiveness of plasticity-damage  
633 coupling and representation of the relationship between micro-cracking and  
634 local softening behavior of the material expand the usability of the model,  
635 as shown in Figures 3 and 4.

636 In the view of understanding fundamental reasons for failure during large  
637 deformations, complexity of the microstructure has a relevance in the stress  
638 concentrations and strain localization in polycrystalline materials. Present  
639 work focused on two different cases: a simplified prior austenite based mi-  
640 crostructure and SEM imaging based microstructure. Deformation behavior  
641 suggests that the strain build-up does not always lead to formation of large in-  
642 dividual (short) cracks before further strain localization takes place. Several  
643 strain localization zones and thus several susceptible micro-cracking sites can  
644 contribute to the formation of crack network. This was the case and typical  
645 for both of the investigated microstructure types, as also observed in [1, 10]  
646 which also focus on evolution of damage in Q&T steels under comparable  
647 loading conditions. Long range effects were observed in Figure 6, where sev-  
648 eral intra-grain or grain boundary character (near grain boundaries) related  
649 micro-defects joined together to form a microstructurally short crack and  
650 transition towards a longer crack several PAGs in size. A short range variant  
651 of the crack nucleation process was presented in Figure 9. Multiple strain  
652 localization sites within  $2\text{-}5\text{ }\mu\text{m}$  of each other first evolved individually before  
653 large macroscopic strain promoted coincidence of the maturing microcracks  
654 to form a more extensive crack spanning the entire computational domain.  
655 Recently, Chatterjee et al. [29] observed that shear localization has a major  
656 role in preceding crack formation. They also noted that the growth of the  
657 cracks are inhibited most effectively by block boundaries, while prior austen-  
658 ite and packet boundaries have slightly lesser role in the studied martensitic  
659 steel grade. The numerical observations on martensitic microstructure in  
660 the current study agrees well with the failure process of their experiments.  
661 Further simulations of a number of different discrete martensitic microstruc-  
662 tures at different loading conditions could even better validate the model  
663 capabilities as well as allow the virtual design of martensitic microstructure  
664 to combat damage at loading conditions.

665 The modeling of fatigue crack growth is a special topic that deserves much  
666 attention due to the associated engineering importance. The present work  
667 contributes to this field by employing the introduced and calibrated crystal  
668 plasticity and damage models as well as the crack closure concept to high  
669 strength Q&T steels. An analysis of cyclic loading fatigue case, in Figures  
670 10-12, showed that the model has not only promising features in nucleat-  
671 ing and growing cracks (damage) within discrete microstructures, but also  
672 the capability to suppress crack growth under tension-compression reverted  
673 deformation conditions. Furthermore, with increasing amount of cycles, plas-

674 ticity and damage, the asymmetric flow stress behavior was demonstrated by  
675 the model, a feature considered relevant for assessing defect evolution under  
676 cyclic loading in complex microstructures. One origin for this behavior arises  
677 from the fact that complete crack closure does not necessarily take place even  
678 in fully inverted strain controlled deformation conditions. A present short-  
679 coming of the model is that shear damage was not allowed further take place  
680 when the crack is in closed state due to uncertainties in the frictional behav-  
681 ior and contact conditions. In fatigue conditions considering more multiaxial  
682 modes of loading, it may have a non-diminishable role and remains an open  
683 further point of study.

684 Size scaling of plasticity and regularization of the damage are not ad-  
685 dressed in the present context. Previous papers [25, 24] have suggested for-  
686 mulation of microdamage variable in small deformation framework to control  
687 the diffusion of damage in finite element based approaches, with a successful  
688 almost complete removal of mesh dependency. An alternative model based  
689 on reduced micromorphic formulation has been suggested for ductile damage  
690 for FCC materials utilizing Gurson-type of a damage model in [26]. Its finite  
691 strain basis enlarges the usability of the model to extreme loading conditions,  
692 whereas size scaling is achieved with a cumulative microplasticity extension.  
693 They noted that despite microdamage regularization is not directly imple-  
694 mented, as in ref. [24], the plasticity driven damage is also regularized by size  
695 dependent plasticity model. In more brittle behavior case, the microdamage  
696 approach could provide even better regularization due to less interaction with  
697 plasticity. However, as the present model and its related model parameteriza-  
698 tion are strongly coupled with plasticity, it could become more attractive to  
699 simply involve microplasticity driven length scale dependency in the present  
700 plasticity model, as suggested in ref. [26]. This strategy may sufficiently  
701 provide damage regularization and could remain computationally efficient  
702 because of only one added degree of freedom to the formulation. This, how-  
703 ever, is considered beyond the scope of this work and a topic for following  
704 studies. Another topic of great importance from the material perspective is  
705 to address microstructure scale defects such as nonmetallic inclusions, the  
706 work which was initiated by present authors in Frondelius et al. [1].

## 707 5. Conclusions

708 A micromechanical approach based on crystal plasticity was used to in-  
709 vestigate the strain localization and damage behavior of high strength Q&T



710 steel under monotonic and cyclic loading conditions. A model coupling plas-  
711 ticity and damage was formulated in a finite strains framework for crystalline  
712 BCC materials. The results of this work can be concluded as follows:

- 713 • A parametric study on the effect of model parameters was performed,  
714 producing a view on the range of feasible material types that could  
715 be undertaken by the model. It revealed that the plasticity-damage  
716 model can represent characteristics found on brittle and more ductile  
717 BCC materials. A model parametrization technique was suggested and  
718 utilized for a martensitic Q&T steel grade with a Lemaitre-type ten-  
719 sile loading-unloading experiments using a full field polycrystalline mi-  
720 crostructural model. The model could well reproduce the stress-strain  
721 behavior and the failure process of the material.
- 722 • Model capabilities were investigated with polycrystalline microstruc-  
723 tures under severe deformation conditions, including cyclic deforma-  
724 tion conditions relevant in fatigue. Plasticity driven damage growth  
725 can be simulated with the model from micro-crack evolution over to  
726 short crack state progressing from intra-grain to inter-grain and fur-  
727 ther, irrespective of the complexity of the computational microstruc-  
728 ture. This feature opens many opportunities for investigating failure  
729 performance of martensitic steels, whenever short crack nucleation and  
730 growth control the long crack based failure process.
- 731 • Restrictions and improvements of the model related to size dependen-  
732 cies of plasticity and damage regularization are possible to overcome  
733 and achievable through modifications in the modeling framework by  
734 making use of, e.g., the reduced micromorphic extension.

## 735 Acknowledgement

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741 **Data availability**

742 The raw/processed data required to reproduce these findings cannot be  
743 shared at this time as the data also forms part of an ongoing study.

744 **References**

745 **References**

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