

Research Article

On the Study of Global Solutions for a Nonlinear Equation

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The well-posedness of global strong solutions for a nonlinear partial differential equation including the Novikov equation is established provided that its initial value $v_0(x)$ satisfies a sign condition and $v_0(x) \in H^s(R)$ with $s > 3/2$. If the initial value $v_0(x) \in H^s(R)$ ($1 \leq s \leq 3/2$) and the mean function of $(1 - \partial_x^2)v_0(x)$ satisfies the sign condition, it is proved that there exists at least one global weak solution to the equation in the space $v(t, x) \in L^2([0, +\infty), H^s(R))$ in the sense of distribution and $v_x \in L^\infty([0, +\infty) \times R)$.

1. Introduction

Recently, Wu [1] obtained the existence of local solutions in the space $C([0, T]; H^s(R)) \cap C^1([0, T]; H^{s-1}(R))$ with $s > 3/2$ for the following nonlinear equation:

$$\begin{aligned} v_t - v_{txx} + kv^m v_x + (m+3)v^{m+1}v_x \\ = (m+2)v^m v_x v_{xx} + v^{m+1}v_{xxx} + \lambda(v - v_{xx}), \end{aligned} \quad (1)$$

where $m \geq 0$ is a natural number, $k \geq 0$, and λ is a constant. Letting $m = 0$ and $\lambda = 0$, (1) becomes the Camassa-Holm equation [2]. If $m = 1$, $k = 0$, and $\lambda = 0$, (1) reduces to the Novikov equation [3].

A lot of works have been carried out to study various dynamic properties for the Camassa-Holm and the Novikov equations. Xin and Zhang [4] proved that there exists a global weak solution for the Camassa-Holm equation in the space $H^1(R)$ without the assumption of sign conditions on the initial value. Coclite et al. [5] investigated the global weak solutions for a generalized hyperelastic rod wave equation or a generalized Camassa-Holm equation. It is shown in Constantin and Escher [6] that the blowup occurs in the form of breaking waves; namely, the solution remains bounded but its slope becomes unbounded in finite time. After wave breaking, the solution can be continued uniquely either as a global conservative weak solution [7] or a global dissipative solution [8–10]. The periodic and the nonperiodic

Cauchy problems for the Novikov equation were discussed by Grayshan [11] in the Sobolev space. Using the Galerkin-type approximation method, Himonas and Holliman [12] established the well-posedness for the Novikov model in the Sobolev space $H^s(R)$ with $s > 3/2$ on both the line and the circle. The scattering theory was employed in Hone et al. [13] to find nonsmooth explicit soliton solutions with multiple peaks for the Novikov equation. Wu and Zhong [14] proved the existence of local strong and weak solutions for a generalized Novikov equation.

The objective of this work is to study (1) with $k = 0$. Namely, we investigate the problem

$$\begin{aligned} v_t - v_{txx} + (m+3)v^{m+1}v_x \\ = (m+2)v^m v_x v_{xx} + v^{m+1}v_{xxx} + \lambda(v - v_{xx}), \quad (2) \\ v(0, x) = v_0(x), \end{aligned}$$

where m , k , and λ are described in (1). Assuming that the initial value $v_0(x)$ satisfies a sign condition and $v_0(x) \in H^s(R)$, $s > 3/2$, we will show that there exists a unique global strong solution in the Sobolev space $C([0, \infty); H^s(R)) \cap C^1([0, \infty); H^{s-1}(R))$. If the initial value $v_0(x) \in H^s(R)$ ($1 \leq s \leq 3/2$) and the mean function of $(1 - \partial_x^2)v_0(x)$ satisfies the sign condition, it is shown that there exists at least one global weak solution to the equation in the space $v(t, x) \in$

$L^2([0, +\infty), H^s(R))$ in the sense of distribution and $v_x \in L^\infty([0, +\infty) \times R)$.

The structure of this paper is as follows. The main results are given in Section 2. Several lemmas are given in Section 3. Section 4 establishes the proof of the main results.

2. Main Results

We define

$$\phi(x) = \begin{cases} e^{1/(x^2-1)}, & |x| < 1, \\ 0, & |x| \geq 1, \end{cases} \quad (3)$$

and let $\phi_\varepsilon(x) = \varepsilon^{-1/4} \phi(\varepsilon^{-1/4}x)$ with $0 < \varepsilon < 1/4$. For the convolution $v_{\varepsilon 0} = \phi_\varepsilon \star v_0$, we know that $v_{\varepsilon 0} \in C^\infty$ for any $v_0 \in H^s$ with $s > 0$. Notation $(1 - \partial_x^2)v \in N^+(R)$ (or equivalently $(1 - \partial_x^2)v \in N^-(R)$) means that the mean function of $(1 - \partial_x^2)v$ is nonnegative; namely, $(1 - \partial_x^2)v \star \phi_\varepsilon \geq 0$ (or equivalently $(1 - \partial_x^2)v \star \phi_\varepsilon \leq 0$) for an arbitrary sufficiently small $\varepsilon > 0$. For $T > 0$ and nonnegative number s , we let $C([0, T]; H^s(R))$ denote the Frechet space of all continuous H^s -valued functions on $[0, T)$ and write $\Lambda = (1 - \partial_x^2)^{1/2}$.

We state the result of global strong solutions for problem (2).

Theorem 1. *Let $v_0(x) \in H^s(R)$, $s > 3/2$, and $(1 - \partial_x^2)v_0 \geq 0$ for all $x \in R$ or $(1 - \partial_x^2)v_0 \leq 0$ for all $x \in R$. Then problem (2) has a unique strong solution satisfying*

$$v(t, x) \in C([0, \infty); H^s(R)) \cap C^1([0, \infty); H^{s-1}(R)). \quad (4)$$

Definition 2. A function $v(t, x) \in L^2([0, +\infty), H^s(R))$ is called a global weak solution to problem (2) if for every $T > 0$ and all $\varphi(t, x) \in C_0^\infty([0, T] \times R)$, it holds that

$$\begin{aligned} \int_0^T \int_R \left[v_t - v_{txx} + (m+3)v^{m+1}v_x - (m+2)v^m v_x v_{xx} \right. \\ \left. - v^{m+1}v_{xxx} - \lambda(v - v_{xx}) \right] \varphi(t, x) dx dt = 0 \end{aligned} \quad (5)$$

with $v(0, x) = v_0(x)$.

Now we give the main result of global weak solution for problem (2).

Theorem 3. *Let $v_0(x) \in H^s(R)$, $1 \leq s \leq 3/2$, $(1 - \partial_x^2)v_0 \in N^+(R)$ (or equivalently $(1 - \partial_x^2)v_0 \in N^-(R)$). Then problem (2) has a unique global weak solution $v(t, x) \in L^2([0, +\infty), H^s(R))$ in the sense of distribution and $v_x \in L^\infty([0, +\infty) \times R)$.*

3. Several Lemmas

Lemma 4 (see [1]). *Let $v_0(x) \in H^s(R)$ with $s > 3/2$. Then the Cauchy problem (2) has a unique local solution*

$$v(t, x) \in C([0, T]; H^s(R)) \cap C^1([0, T]; H^{s-1}(R)), \quad (6)$$

where $T > 0$ depends on $\|v_0\|_{H^s(R)}$.

Using the first equation of system (2) derives

$$\frac{d}{dt} \int_R (v^2 + v_x^2) dx = 2\lambda \int_R (v^2 + v_x^2) dx, \quad (7)$$

which yields the conservation law

$$\int_R (v^2 + v_x^2) dx = \int_R (v_0^2 + v_{0x}^2) dx \quad (8)$$

$$+ 2\lambda \int_0^t \int_R (v^2 + v_x^2) dx dt.$$

Lemma 5 (see [1]). *Let $s > 3/2$ and the function $v(t, x)$ is a solution of problem (2) and the initial data $v_0(x) \in H^s$. Then the following inequalities hold:*

$$\begin{aligned} \|v\|_{H^1}^2 &\leq \int_R (v^2 + v_x^2) dx \leq \int_R (v_0^2 + v_{0x}^2) dx, \quad \text{if } \lambda \leq 0. \\ \|v\|_{H^1}^2 &\leq \int_R (v^2 + v_x^2) dx \leq e^{2\lambda t} \int_R (v_0^2 + v_{0x}^2) dx, \quad \text{if } \lambda > 0. \end{aligned} \quad (9)$$

For $q \in (0, s-1]$, there is a constant c such that

$$\begin{aligned} \int_R (\Lambda^{q+1}v)^2 dx \\ \leq \int_R (\Lambda^{q+1}v_0)^2 dx \\ + c \int_0^t \|v\|_{H^{q+1}}^2 (|\lambda| + (\|v\|_{L^\infty}^{m-1} + \|v\|_{L^\infty}^m) \|v_x\|_{L^\infty} \\ + \|v\|_{L^\infty}^{m-1} \|v_x\|_{L^\infty}^2) d\tau. \end{aligned} \quad (10)$$

For $q \in [0, s-1]$, there is a constant c such that

$$\begin{aligned} \|v\|_{H^q} \leq c \|v\|_{H^{q+1}} (|\lambda| + (\|v\|_{L^\infty}^{m-1} + \|v\|_{L^\infty}^m) \|v\|_{H^1} \\ + \|v\|_{L^\infty}^m \|v_x\|_{L^\infty} + \|v\|_{L^\infty}^{m-1} \|v_x\|_{L^\infty}^2). \end{aligned} \quad (11)$$

Consider the differential equation

$$\begin{aligned} p_t &= v^{m+1}(t, p), \quad t \in [0, T], \\ p(0, x) &= x, \end{aligned} \quad (12)$$

where $v(t, x)$ is the solution of problem (2) and T is the maximal existence time of the solution.

Lemma 6. *Let $v_0 \in H^s(R)$, $s \geq 3$, and let $T > 0$ be the maximal existence time of the solution to problem (2). Then system (12) has a unique solution $p(t, x) \in C^1([0, T] \times R)$. Moreover, the map $p(t, \cdot)$ is an increasing diffeomorphism of R with $p_x(t, x) > 0$ for $(t, x) \in [0, T] \times R$.*

Proof. From Lemma 4, we know that there exists a unique solution

$$v(t, x) \in C([0, T]; H^s(R)) \cap C^1([0, T]; H^{s-1}(R)). \quad (13)$$

The Sobolev imbedding theorem derives $H^s(R) \in C^1(R)$. This means that two functions $v(t, x)$ and $v_x(t, x)$ are bounded, Lipschitz in space and C^1 in time. Using the existence and uniqueness theorem of ordinary differential equations, we derive that problem (12) has a unique solution $p(t, x) \in C^1([0, T) \times R)$.

Differentiating (12) with respect to x gives rise to

$$\begin{aligned} \frac{d}{dt} p_x &= (m+1) v^m v_x(t, p) p_x, \quad t \in [0, T], \\ p_x(0, x) &= 1, \end{aligned} \quad (14)$$

from which we obtain

$$p_x(t, x) = \exp \left(\int_0^t (m+1) v^m v_x(\tau, p(\tau, x)) d\tau \right). \quad (15)$$

For every $T' < T$, applying the Sobolev imbedding theorem results in

$$\sup_{(\tau, x) \in [0, T'] \times R} |v_x(\tau, x)| < \infty. \quad (16)$$

Therefore, we know that there exists a constant $M > 0$ such that $p_x(t, x) \geq e^{-Mt}$ for $(t, x) \in [0, T) \times R$. The proof is completed. \square

Lemma 7. Let $v_0 \in H^s$ with $s \geq 3$, and let $T > 0$ be the maximal existence time of the problem (2); it holds that

$$y(t, p(t, x)) p_x^2(t, x) = y_0(x) e^{\int_0^t (mv^m v_x + \lambda) d\tau}, \quad (17)$$

where $(t, x) \in [0, T) \times R$ and $y := v - v_{xx}$.

Proof. We have

$$\begin{aligned} & \frac{d}{dt} [y(t, p(t, x)) p_x^2(t, x)] \\ &= y_t p_x^2 + 2y p_{xt} p_x + y_x p_t p_x^2 \\ &= y_t p_x^2 + 2y(m+1) v^m v_x p_x^2 + v^{m+1} y_x p_x^2 \\ &= [y_t + (m+2) v^m v_x y + y_x v^{m+1}] p_x^2 + mv^m v_x y p_x^2 \\ &= [v_t - v_{txx} + (m+2) v^m v_x (v - v_{xx}) \\ & \quad + v^{m+1} (v_x - v_{xxx}) - \lambda (v - v_{xx})] p_x^2 \\ & \quad + (mv^m v_x + \lambda) y p_x^2 \\ &= [v_t - v_{txx} + (m+3) v^{m+1} v_x - (m+2) v^m v_x v_{xx} \\ & \quad - v^{m+1} v_{xxx} - \lambda (v - v_{xx})] p_x^2 \\ & \quad + (mv^m v_x + \lambda) y p_x^2 \\ &= (mv^m v_x + \lambda) y p_x^2, \end{aligned} \quad (18)$$

from which we have

$$y(t, p(t, x)) p_x^2(t, x) = p_x(0, x) y_0(x) e^{\int_0^t (mv^m v_x + \lambda) d\tau}. \quad (19)$$

Using $p_x(0, x) = 1$ completes the proof. \square

Lemma 8. If $v_0 \in H^s(R)$, $s \geq 3/2$, $(1 - \partial_x^2)v_0 \geq 0$ or $(1 - \partial_x^2)v_0 \leq 0$, then the solution of problem (2) satisfies

$$\|v_x\|_{L^\infty} \leq \|v\|_{L^\infty}. \quad (20)$$

Proof. We only need to prove this lemma for the case $v_0 - v_{0xx} \geq 0$ since the proof of the other case $(1 - \partial_x^2)v_0 \leq 0$ is similar. It follows from Lemmas 6 and 7 that $v - v_{xx} \geq 0$. Letting $\xi(t, x) = v - v_{xx}$, we have

$$v = \frac{1}{2} e^{-x} \int_{-\infty}^x e^{\eta} \xi(t, \eta) d\eta + \frac{1}{2} e^x \int_x^{\infty} e^{-\eta} \xi(t, \eta) d\eta, \quad (21)$$

which derives

$$\begin{aligned} \partial_x v(t, x) &= -\frac{1}{2} \left(e^{-x} \int_{-\infty}^x e^{\eta} \xi(t, \eta) d\eta + e^x \int_x^{\infty} e^{-\eta} \xi(t, \eta) d\eta \right) \\ & \quad + e^x \int_x^{\infty} e^{-\eta} \xi(t, \eta) d\eta \\ &= -v(t, x) + e^x \int_x^{\infty} e^{-\eta} \xi(t, \eta) d\eta \\ &\geq -v(t, x). \end{aligned} \quad (22)$$

On the other hand, we have

$$\begin{aligned} \partial_x v(t, x) &= \frac{1}{2} \left(e^{-x} \int_{-\infty}^x e^{\eta} \xi(t, \eta) d\eta + e^x \int_x^{\infty} e^{-\eta} \xi(t, \eta) d\eta \right) \\ & \quad - e^{-x} \int_{-\infty}^x e^{\eta} \xi(t, \eta) d\eta \\ &= v(t, x) - e^{-x} \int_{-\infty}^x e^{\eta} \xi(t, \eta) d\eta \\ &\leq v(t, x). \end{aligned} \quad (23)$$

The inequalities (22) and (23) derive that inequality (20) is valid. \square

Lemma 9. For $s > 0$, $u \in H^s(R)$, and $u_\varepsilon = \phi_\varepsilon \star u$, it holds that

$$\begin{aligned} \|u_{\varepsilon x}\|_{L^\infty} &\leq c \|u_x\|_{L^\infty}, \\ \|u_\varepsilon\|_{H^q} &\leq c, \quad \text{if } q \leq s, \\ \|u_\varepsilon\|_{H^q} &\leq c \varepsilon^{(s-q)/4}, \quad \text{if } q > s, \\ \|u_\varepsilon - u\|_{H^q} &\leq c \varepsilon^{(s-q)/4}, \quad \text{if } q \leq s, \\ \|u_\varepsilon - u\|_{H^s} &= o(1), \end{aligned} \quad (24)$$

where c is a constant independent of ε .

The proof of this lemma can be found in [15, 16].
From Lemma 4, it derives that the Cauchy problem

$$\begin{aligned}
 v_t - v_{txx} &= -(m+3)v^{m+1}v_x + (m+2)v^m v_x v_{xx} \\
 &\quad + v^{m+1}v_{xxx} + \lambda(v - v_{xx}) \\
 &= -\frac{m+3}{m+2}(v^{m+2})_x + \frac{1}{m+2}\partial_x^3(v^{m+2}) \\
 &\quad - (m+1)\partial_x(v^m v_x^2) + v^m v_x v_{xx} + \lambda(v - v_{xx}), \\
 v(0, x) &= v_{\varepsilon 0}(x),
 \end{aligned} \tag{25}$$

has a unique solution v depending on the parameter ε . We write $v_\varepsilon(t, x)$ to represent the solution of problem (25). Using Lemma 4 derives that $v_\varepsilon(t, x) \in C^\infty([0, T], H^\infty(R))$ since $v_{\varepsilon 0}(x) \in C_0^\infty(R)$.

Lemma 10. *Provided that $v_0 \in H^s(R)$, $1 \leq s \leq 3/2$, and $(1 - \partial_x^2)v_0 \in N^+(R)$ (or equivalently $(1 - \partial_x^2)v_0 \in N^-(R)$), then there exists a constant $c > 0$ independent of ε and t such that the solution of problem (25) satisfies*

$$\|v_{\varepsilon x}\|_{L^\infty} \leq ce^{ct}. \tag{26}$$

Proof. Using Lemmas 5 and 9, if $v_0 \in H^s(R)$ with $1 \leq s \leq 3/2$, we have

$$\|v_\varepsilon\|_{L^\infty(R)} \leq c\|v_\varepsilon\|_{H^1(R)} \leq ce^{ct}\|v_{\varepsilon 0}\|_{H^1(R)} \leq ce^{ct}, \tag{27}$$

where c is independent of ε and t .

From Lemma 8, we have

$$\|v_{\varepsilon x}\|_{L^\infty(R)} \leq \|v_\varepsilon\|_{L^\infty(R)}, \tag{28}$$

which completes the proof. \square

4. Proof of Main Results

Proof of Theorem 1. Since $\|v\|_{L^\infty(R)} \leq c\|v\|_{H^1(R)} \leq ce^{ct}$ and taking $q + 1 = s$ in inequality (10), we have

$$\|v\|_{H^s}^2 \leq \|v_0\|_{H^s}^2 + c \int_0^t e^{c\tau} \|v\|_{H^s}^2 (\|v_x\|_{L^\infty} + \|v_x\|_{L^\infty}^2) d\tau, \tag{29}$$

from which we obtain

$$\|v\|_{H^s} \leq \|v_0\|_{H^s} e^{c \int_0^t e^{c\tau} (\|v_x\|_{L^\infty} + \|v_x\|_{L^\infty}^2) d\tau}. \tag{30}$$

Applying Lemma 8 yields

$$\|v\|_{H^s} \leq \|v_0\|_{H^s} ce^{e^{ct}}, \tag{31}$$

from which we complete the proof of Theorem 1. \square

Provided that $1 \leq s \leq 3/2$, for problem (25), applying Lemmas 5, 8, and 10, and the Gronwall's inequality, we obtain the inequalities

$$\begin{aligned}
 \|v_\varepsilon\|_{H^1} &\leq \|v_{\varepsilon 0}\|_{H^1} \leq ce^{ct}, \\
 \|v_\varepsilon\|_{H^q} &\leq c\|v_{\varepsilon 0}\|_{H^q} \exp \left[\int_0^t (\|v_{\varepsilon x}\| + \|v_{\varepsilon x}\|_{L^\infty}^2) d\tau \right] \leq ce^{e^{ct}}, \\
 \|u_{\varepsilon t}\|_{H^r} &\leq c\|u_\varepsilon\|_{H^{r+1}} (1 + e^{ct}) \leq c(1 + e^{ct}),
 \end{aligned} \tag{32}$$

where $q \in (0, s]$, $r \in [0, s-1]$, and c is a constant independent of t and ε . Using the Aubin compactness theorem, we know that there is a subsequence $\{v_{\varepsilon_n}\}$ of $\{v_\varepsilon\}$ such that $\{v_{\varepsilon_n}\}$ and their temporal derivatives $\{v_{\varepsilon_n t}\}$ converge weakly to a function $v(t, x)$ and its derivative v_t in the space $L^2([0, T], H^s(R))$ and $L^2([0, T], H^{s-1}(R))$, respectively, where T is an arbitrary fixed positive number. In addition, for any real number $M_1 > 0$, $\{v_{\varepsilon_n}\}$ converges strongly to the function v in the space $L^2([0, T], H^q(-M_1, M_1))$ for $q \in (0, s]$ and $\{v_{\varepsilon_n t}\}$ converges strongly to v_t in the space $L^2([0, T], H^r(-M_1, M_1))$ for $r \in [0, s-1]$.

Proof of Theorem 3. For an arbitrary fixed $T > 0$, using Lemma 10, we know that $\{v_{\varepsilon_n x}\}$ ($\varepsilon_n \rightarrow 0$) is bounded in the space L^∞ . Therefore, we derive that the sequences $\{v_{\varepsilon_n}\}$, $\{v_{\varepsilon_n x}\}$, $\{v_{\varepsilon_n x}^2\}$, and $\{v_{\varepsilon_n x}^3\}$ converge weakly to v , v_x , v_x^2 , and v_x^3 in $L^2([0, T], H^r(-R_1, R_1))$ for any $r \in [0, s-1]$, separately. Applying the identity $v^m(v_x^2)_x = (v^m v_x^2)_x - (v^m)_x v_x^2$, we conclude that v satisfies the equation

$$\begin{aligned}
 & - \int_0^T \int_R v(\varphi_t - \varphi_{xxt}) dx dt \\
 &= \int_0^T \int_R \left[\left(\frac{m+3}{m+2} v^{m+2} + (m+1)v^m v_x^2 \right) \varphi_x \right. \\
 &\quad \left. - \frac{1}{m+2} v^{m+2} \varphi_{xxx} - \frac{1}{2} v^m v_x^2 \varphi_x \right. \\
 &\quad \left. - \frac{m}{2} v^{m-1} v_x^3 \varphi + \lambda v(\varphi - \varphi_{xx}) \right] dx dt,
 \end{aligned} \tag{33}$$

where $\varphi(t, x) \in C_0^\infty([0, T] \times R)$. We know that $Y = L^1([0, T] \times R)$ is a separable Banach space and $\{v_{\varepsilon_n x}\}$ is a bounded sequence in the dual space $Y^* = L^\infty([0, T] \times R)$ of Y . Thus, there exists a subsequence of $\{v_{\varepsilon_n x}\}$, still denoted by $\{v_{\varepsilon_n x}\}$, weakly star convergent to a function u in $L^\infty([0, T] \times R)$. Since $\{v_{\varepsilon_n}\}$ weakly converges to v_x in $L^2([0, T] \times R)$, it derives that $v_x = u$ almost everywhere. Therefore, we obtain $v_x \in L^\infty([0, T] \times R)$. Since $T > 0$ is an arbitrary number, we complete the proof of existence of global weak solutions to problem (2). \square

Conflict of Interests

The authors declare that there is no conflict of interests regarding the publication of this paper.

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