

On measures of accretion and dissipation for solutions of the Camassa-Holm equation

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Abstract

We investigate the measures of dissipation and accretion related to the weak solutions of the Camassa-Holm equation. Demonstrating certain properties of nonunique characteristics, we prove a new representation formula for these measures and conclude about their structural features, such as singularity with respect to the Lebesgue measure. We apply these results to gain new insights into the structure of weak solutions, proving in particular that measures of accretion vanish for dissipative solutions of the Camassa-Holm equation.

Keywords: Camassa-Holm, weak solution, dissipative solution, generalized characteristics, measure of accretion, measure of dissipation

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1 Introduction

In this paper we study the Camassa-Holm equation ([6]),

$$u_t - u_{xxt} + 3uu_x = 2u_x u_{xx} + uu_{xxx}, \quad (1)$$

which is a one-dimensional model of unidirectional water wave propagation in shallow canals, where t denotes time, x is a one-dimensional space variable and $u(t, x)$ corresponds to the horizontal velocity of the water surface (see [12]). Equation (1) had been previously derived, though without the physical context of [6], by Fokas and Fuchssteiner [19] as a bihamiltonian generalization of the celebrated Korteweg-de Vries equation (KdV) and is formally integrable [11]. The hallmark of the Camassa-Holm equation, which makes it one of the most relevant models of the shallow water theory, is that it accounts both for peaked solitons ([6]) and, unlike the KdV equation [28], yet similarly as the Whitham equation [31], for wave-breaking (see [10]).

The well-posedness theory for (1) in the case of smooth solutions is due to A. Constantin and J. Escher [8, 9]. The same authors formulated also fairly general conditions precluding wave-breaking [10]. Since then, other criteria have been proposed and recently Brandolese [1] has obtained a general criterion, which nicely encompasses most previous non-breaking criteria, see references therein. Nevertheless, for general initial data the wave-breaking, understood as blow-up of the L^∞ norm of the derivative u_x , is unavoidable. Importantly, however, the Camassa-Holm equation preserves physical relevance also after wave breaking and thus it is worth to study weak (i.e. non-smooth) solutions after the breaking time. This

can be done in the framework provided by (with $*$ denoting convolution)

$$\partial_t u + \partial_x(u^2/2) + P_x = 0, \quad (2)$$

$$P(t, x) = \frac{1}{2}e^{-|x|} * \left(u^2(t, \cdot) + \frac{u_x^2(t, \cdot)}{2} \right), \quad (3)$$

$$u(t = 0, \cdot) = u_0, \quad (4)$$

which for smooth solutions is equivalent to (1), with equivalence provided through the nonlocal operator $(I - \partial_{xx})^{-1}$, which satisfies $(I - \partial_{xx})(\frac{1}{2}e^{-|x|}) = \delta(x)$.

Solutions of (2)-(4) are considered in the space $L^\infty([0, \infty), H^1(\mathbb{R}))$, which corresponds to the maximal physically relevant class of solutions with bounded total energy, given by

$$E(t) = \frac{1}{2} \int_{\mathbb{R}} (u^2(t, x) + u_x^2(t, x)) dx.$$

Let us note that the total energy remains constant for smooth solutions (see e.g. introduction of [3]), yet fails to be so for weak solutions, for which it can both increase and decrease. This allows the weak solutions to account for many phenomena such as soliton interactions, energy dissipation etc.

Definition 1.1 (Weak solutions). *Let $u_0 \in H^1(\mathbb{R})$. We say that a function $u : [0, \infty) \times \mathbb{R} \rightarrow \mathbb{R}$ is a weak solution of (2)-(4) if*

- $u(t, x) \in C([0, \infty) \times \mathbb{R}) \cap L^\infty([0, \infty), H^1(\mathbb{R}))$,
- $u(t = 0, x) = u_0(x)$ for $x \in \mathbb{R}$,
- $u(t, x)$ satisfies (2) in the sense of distributions.

Existence of weak solutions of the Camassa-Holm equation was proven by Xin and Zhang in [32] by the method of vanishing viscosity, which resulted in the so-called *dissipative weak solutions*.

Definition 1.2 (Dissipative weak solutions). *A weak solution of (2)-(4) is called dissipative if*

- $\partial_x u(t, x) \leq \text{const} (1 + \frac{1}{t})$ (Oleinik-type condition)
- $\|u(t, \cdot)\|_{H^1(\mathbb{R})} \leq \|u(0, \cdot)\|_{H^1(\mathbb{R})}$ for every $t > 0$ (weak energy condition).

Dissipative weak solutions are only one of the many classes of (extremely nonunique) weak solutions of (2)-(4). Other classes include the conservative [2] solutions, which satisfy a supplementary conservation law ensuring local conservation of the energy, and intermediate [22] solutions, which interpolate between the conservative and dissipative ones. The uniqueness of conservative solutions was demonstrated by Bressan and Fonte [5] by use of a distance functional related to the optimal transportation problem, see also [24, 20, 21] and [4] for more recent proofs, based on alternative techniques. The question of uniqueness of dissipative solutions, on the other hand, remained for many years one of the eminent unresolved open problems of the theory and, until recently, only constructions of global semigroups of solutions [3, 23] were available. The issue has been finally resolved affirmatively in [25], relying on the framework introduced in [26], the uniqueness result from [3] and ideas similar in spirit to [4].

In this paper we focus on studying in detail the structure of weak solutions of the Camassa-Holm equation, having in mind another open question of the theory: *Is the dissipative weak solution the unique one, which dissipates the energy at the highest possible rate within the class of all weak solutions with the same initial data?*

The relevance of this question (the so called *maximal dissipation criterion*) in evolutionary equations has been highlighted by Dafermos [15, 16] and more specifically in the

context of the related Hunter-Saxton equation by Zhang and Zheng [33]. The hypothesis of Zhang-Zheng was finally resolved in [7], see also references therein, and it is indeed the framework from [7], relying on the ideas of Dafermos regarding generalized characteristics ([13, 14]), and ported to the Camassa-Holm setting in [26], which provides inspiration for the research in the present paper. At any rate, it is clear that to approach the maximal dissipation criterion, one has to understand the behavior of general weak solutions and this paper is a step in this direction.

To fix some ideas let us in the remainder of this introductory part consider a generic example, which presents possible nonuniqueness scenarios in the class of weak solutions of (2)-(4) and represents one of the few known explicit solutions of the Camassa-Holm equation (see however [29, 30]). This example provides also some insight into the definition of dissipative weak solutions; it shows in particular that the Oleinik-type condition in Definition 1.2, which is an analog of the Oleinik entropy criterion, known from the theory of hyperbolic conservation laws (see e.g. [17, Section 3b]), is insufficient for selecting the unique dissipative solution and thus the weak energy condition is indispensable.

Example 1.3 (Peakon-antipeakon interaction, creation of the peakon-antipeakon pair).
Function

$$u(t, x) = p_1(t)e^{-|x-q_1(t)|} - p_1(t)e^{-|x+q_1(t)|} \quad (5)$$

where $q(0) < 0, p_1(0) > 0, p_1(t) = \frac{1}{2}p(t), q_1(t) = \frac{1}{2}q(t)$ and

$$\begin{aligned} p(t) &= H_0 \frac{[p(0) + H_0] + [p(0) - H_0]e^{H_0 t}}{[p(0) + H_0] - [p(0) - H_0]e^{H_0 t}}, \\ q(t) &= q(0) - 2 \log \frac{[p(0) + H_0]e^{-H_0 t/2} + [p(0) - H_0]e^{H_0 t/2}}{2p(0)}, \\ H_0^2 &= p(0)^2(1 - e^{q(0)}) = p(t)^2(1 - e^{q(t)}) \text{ for every } t > 0 \end{aligned} \quad (6)$$

is, see [2], a weak solution of the Camassa-Holm equation with invariant H_0 . It is composed of two solitary waves – peakon $p_1(t)e^{-|x-q_1(t)|}$, centred at $q_1(t)$ and moving to the right, and antipeakon $-p_1(t)e^{-|x+q_1(t)|}$ moving to the left (see Fig. 1). At time $T = \frac{1}{H_0} \log \frac{p(0)+H_0}{p(0)-H_0}$ the two waves annihilate giving rise to a singularity characterized by

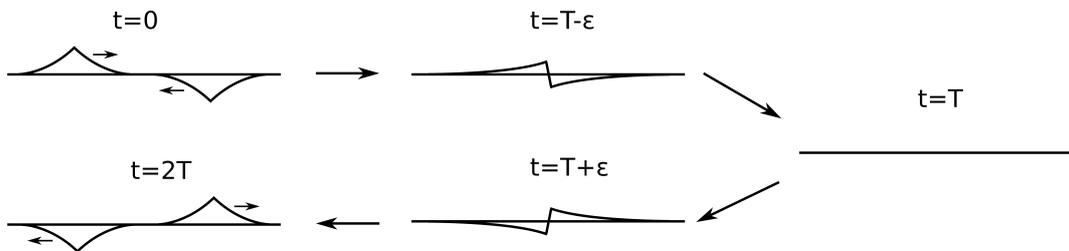


Figure 1: Schematic presentation of a conservative peakon-antipeakon interaction. Two initial waves – peakon (positive amplitude, moving to the right) and antipeakon (negative amplitude, moving to the left) interact, which leads to their annihilation at the critical time T (wave-breaking). Then they reemerge as a peakon and antipeakon moving away from one another. Considering T as initial time we obtain a weak solution of the Camassa-Holm equation corresponding to creation of a peakon-antipeakon pair out of nothing.

$$\lim_{t \rightarrow T^-} \sup_x |u(t, x)| = 0$$

and

$$\lim_{t \rightarrow T^-} e(t, x) = \lim_{t \rightarrow T^-} \frac{1}{2}(u^2(t, x) + u_x^2(t, x)) = \lim_{t \rightarrow T^-} \frac{1}{2}((u_x^-)^2(t, x)) = H_0^2 \delta_0(dx), \quad (7)$$

where the limits are taken in the weak sense, $u_x^- = \max(0, -u_x)$ is the negative part of u_x and $e(t, \cdot)$ is the energy density. To estimate the derivative of the solution we differentiate (5) obtaining

$$u_x = -p_1 \operatorname{sgn}(x - q_1) e^{-|x - q_1|} + p_1 \operatorname{sgn}(x + q_1) e^{-|x + q_1|}$$

and hence $|u_x| \leq 2p_1 = p$. By (6) we have:

$$\begin{aligned} \frac{p(t)}{H_0} &= \frac{[p(0) + H_0] + [p(0) - H_0] e^{H_0 T} e^{H_0(t-T)}}{[p(0) + H_0] - [p(0) - H_0] e^{H_0 T} e^{H_0(t-T)}} = \frac{[p(0) + H_0] + [p(0) + H_0] e^{H_0(t-T)}}{[p(0) + H_0] - [p(0) + H_0] e^{H_0(t-T)}} \\ &= \frac{1 + e^{H_0(t-T)}}{1 - e^{H_0(t-T)}} = \frac{e^{H_0(T-t)} + 1}{e^{H_0(T-t)} - 1} \leq \frac{e^{H_0 T} + 1}{H_0(T-t)}, \end{aligned}$$

and consequently

$$|u_x| \leq C(T)/(T-t). \quad (8)$$

By setting

$$u(t, x) = -u(2T - t, x)$$

for $t > T$ we can prolong the solution beyond the blow-up time. This 'conservative' prolongation can be interpreted as reemergence after interaction of the peakon-antipeakon pair, which are moving now away from one another (see Fig. 1). Note that if we set $u(t, x) = 0$ for $t > T, x \in \mathbb{R}$ we would obtain a 'dissipative' prolongation.

Finally, consider $w(t, x) := u(t+T, x)$ defined for $t \geq 0$. As a simple translation in time of u , w is a weak solution of (2)-(4), which satisfies $w(t=0, \cdot) \equiv 0$. Thus, w represents creation of a peakon-antipeakon pair. Due to estimate (8) we obtain

$$|w_x(t, x)| \leq \tilde{C}/t,$$

Thus the Oleinik-type criterion from Definition 1.2 is satisfied. The weak energy condition in Definition 1.2, on the other hand, is violated. And indeed, given the initial condition $w(t=0, \cdot) \equiv 0$, the 'entropy solution', called in the context of Camassa-Holm 'dissipative weak solution', is a function equal identically 0. The weak energy condition from Definition 1.2 allows us to distinguish between these two scenarios.

Let us point out that

- creation of a peakon-antipeakon pair, as presented in Example 1.3, can occur at any point of spacetime, which means that weak solutions can, in general, exhibit very complex structures and are highly nonunique,
- creation of a peakon-antipeakon pair involves creation of a finite portion of energy at a given point in spacetime; similarly, annihilation of a peakon-antipeakon pair corresponds to annihilation of a finite portion of energy,
- the creation (and annihilation) of a finite portion of energy at any single timepoint can also be spread over the space, leading to a distribution, which, as we will see in the following, can be accounted for by a measure,
- energy may be also transferred in a continuous fashion (as is the case e.g. for $t < T$ or $t > T$ in Example 1.3).

Our goal in this paper, motivated by the above observations, is to gain new insights into how the energy is accrued/created and dissipated/annihilated in arbitrary weak solutions of (2)-(4). To this end, we define rigorously and study the so-called *accretion* and *dissipation* measures (introduced informally in [26]) and obtain some qualitative results regarding their structure. In particular, we show (see Section 2 for rigorous formulations) that creation of energy at a given timepoint corresponds to a measure which is necessarily singular with respect to the one-dimensional Lebesgue measure and thus cannot be arbitrary. Since, as

we demonstrate, the creation of energy for dissipative solutions cannot occur on a singular set, we conclude that for dissipative weak solutions measures of accretion vanish. In the same vein, we show that if a weak solution is to dissipate the energy at the highest possible rate then measures of accretion have to vanish.

The structure of the paper is the following. In Section 2 we present our main results, recalling first the key results from [26]. In Section 3 we demonstrate some new properties of nonunique characteristics, which we then use in Section 4 to prove a representation formula for measures of accretion/dissipation. Finally, in Section 5 we demonstrate, using the new representation formula, the conclusions regarding the structure of weak and dissipative weak solutions of the Camassa-Holm equation.

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2 Main results

In this section we present our main results. Since they rely heavily on the framework developed in [26], let us first recall the key results of [26], which, although more general, for the purposes of the present paper have been restricted to the Camassa-Holm equation.

The first proposition (Proposition 2.4 from [26], which follows from the Peano existence theorem and [13, Lemma 3.1]) asserts that characteristics, although in general nonunique, exist.

Proposition 2.1. *Let u be a weak solution of (2)-(4). Then for every $\zeta \in \mathbb{R}$ there exists a (nonunique) characteristic of u emanating from ζ , i.e. a function $\zeta : [0, \infty) \rightarrow \mathbb{R}$, which satisfies:*

- $\zeta(0) = \zeta$,
- $\frac{d}{dt}\zeta(t) = u(t, \zeta(t))$,
- $\frac{d}{dt}u(t, \zeta(t)) = -P_x(t, \zeta(t))$.

The second result combines Proposition 2.5, Corollary 6.1 and Lemmas 4.1 and 4.2 from [26] and states that also u_x , when restricted to certain set of initial points, evolves along characteristics.

Proposition 2.2 (Proposition 2.5, Corollary 6.1 from [26]). *Let u be a weak solution of (2)-(4) and fix $t_0 \geq 0$. There exists a family of sets $\{S_{t_0, T}\}_{T > t_0}$ such that for every $\zeta \in S_{t_0, T}$ the characteristic $\zeta(\cdot)$, with $\zeta(t_0) = \zeta$, is unique on $[t_0, T]$, ζ is a Lebesgue point of $u_x(t_0, \cdot)$, and for $t_0 \leq t < T$*

$$\dot{v}(t) = u^2(t) - \frac{1}{2}v^2(t) - P(t), \quad (9)$$

where $u(t) := u(t, \zeta(t))$, $v(t) := u_x(t, \zeta(t))$ and $P(t) := P(t, \zeta(t))$. Moreover, $S_{t_0, T_1} \subset S_{t_0, T_2}$ for $T_1 > T_2$ and $|\mathbb{R} \setminus \bigcup_{T > t_0} S_{t_0, T}| = 0$. Finally, for every $\zeta \in S_{t_0, T}$ there exists $N > 0$ such that $|v(\cdot)| \leq N$ on $[t_0, T]$.

Remark 2.3. *Equation (9) is satisfied in the sense*

$$v(t) - v(t_0) = \int_{t_0}^t \left(u^2(s) - \frac{1}{2}v^2(s) - P(s) \right) ds$$

for almost every $t \in (t_0, T]$, such that $\zeta(t)$ is a Lebesgue point of $u_x(t, \cdot)$ (in other words, difference quotients converge to $u_x(t, \zeta(t))$). Since the set of Lebesgue points is for every fixed t a full measure set, there exists a modification of $u_x(t, x)$ (for every t on a set of x of measure 0) such that equation (9) is satisfied for every $t \in [t_0, T]$.

These two results allowed us in [26] to establish the main theorem on the weak continuity properties of the positive and negative parts of the derivative u_x . Below, by $u_x^\pm := \max(\pm u_x, 0)$ we denote the positive/negative parts of function u_x .

Theorem 2.4 (see Theorem 2.9 in [26] and Remark 2.7). *Let u be a weak solution of the Camassa-Holm equation. Then*

- function $t \mapsto (u_x^+(t, \cdot))^2$ is weakly ladcag (left-continuous with right limits),
- function $t \mapsto (u_x^-(t, \cdot))^2$ is weakly cadlag (right-continuous with left limits),
- the functions $t \mapsto \int_{\mathbb{R}} \phi(x)(u_x^+)^2(t, x)dx$ and $t \mapsto \int_{\mathbb{R}} \phi(x)(u_x^-)^2(t, x)dx$ have locally bounded variation provided ϕ is Lipschitz continuous,
- the limits

$$\phi \mapsto \lim_{t \rightarrow t_0^+} \int_{\mathbb{R}} \phi(x)(u_x^+)^2(t, x)dx$$

and

$$\phi \mapsto \lim_{t \rightarrow t_0^-} \int_{\mathbb{R}} \phi(x)(u_x^-)^2(t, x)dx$$

define, for every fixed $t_0 \geq 0$, bounded linear functionals on $C_c(\mathbb{R})$.

As mentioned in [26], applying Theorem 2.4 we can, using the Riesz representation theorem, define the objects of interest of the present paper – measures of accretion and dissipation. Their rigorous definition is the point of departure of the present paper.

Definition 2.5 (Measures of accretion and dissipation). *Let u be a weak solution of (2)-(4). The measure of accretion μ^+ is defined for every $t_0 \in [0, \infty)$ as*

$$\int_{\mathbb{R}} \phi(x)\mu^+(t_0, dx) := \lim_{t \rightarrow t_0^+} \int_{\mathbb{R}} \phi(x)(u_x^+)^2(t, x)dx - \int_{\mathbb{R}} \phi(x)(u_x^+)^2(t_0, x)dx,$$

for any continuous compactly supported function ϕ . Similarly, for every $t_0 \in (0, \infty)$ the measure of dissipation μ^- is defined as

$$\int_{\mathbb{R}} \phi(x)\mu^-(t_0, dx) := \int_{\mathbb{R}} \phi(x)(u_x^-)^2(t_0, x)dx - \lim_{t \rightarrow t_0^-} \int_{\mathbb{R}} \phi(x)(u_x^-)^2(t, x)dx,$$

for any continuous compactly supported function ϕ .

Remark 2.6. *Measures obtained by the use of the Riesz theorem are in general signed. Due to, however, the last inequality from the proof of [26, Theorem 2.9] (see also Remark 2.7 below)*

$$\lim_{t \rightarrow t_0^+} \int_{\mathbb{R}} \phi(x)(u_x^+)^2(t, x)dx - \int_{\mathbb{R}} \phi(x)(u_x^+)^2(t_0, x)dx$$

is always nonnegative. Thus, $\mu^+(t_0, dx)$ is in fact a nonnegative measure. Similarly, $\mu^-(t_0, dx)$ represents, for any t_0 , a nonpositive measure.

Remark 2.7. *Note the missing factor $\sup(\phi)$ in the last inequality in [26]. The correct form of this inequality is $\int_{\mathbb{R}} \phi u_x^+(t_2, z)^2 dz \geq \int_{\mathbb{R}} \phi u_x^+(t_1, \gamma)^2 d\gamma - (t_2 - t_1)\sup(\phi)(K(a_I(t_1) - a_0(t_1)) + \int_{[a_0(t_1), a_I(t_1)]} u_x(t_1, \gamma)^2 d\gamma - (Lip(\phi)\sup(u)(t_2 - t_1)) \int_{\mathbb{R}} u_x^2(t_2, z)dz$. Moreover, this form is valid for ϕ which are Lipschitz continuous. For general $\phi \in C_c(\mathbb{R})$ one has to replace the term $Lip(\phi)\sup(u)(t_2 - t_1)$ by $MC^\phi(\sup(u)(t_2 - t_1))$, where MC^ϕ is the modulus*

of continuity of ϕ . This means that one obtains, by [26, Lemma 9.1], BV regularity of $\phi \mapsto \int_{\mathbb{R}} \phi(x)(u_x^+)^2(t, x)dx$ only for Lipschitz continuous ϕ , as formulated here in Theorem 2.4. Nevertheless, the right limits of $\phi \mapsto \int_{\mathbb{R}} \phi(x)(u_x^+)^2(t, x)dx$ exist, by [26, Lemma 9.1], for arbitrary $\phi \in C_c(\mathbb{R})$.

Remark 2.8. Measures of accretion and dissipation are defined here for every fixed t_0 separately and thus cannot account for accretion/dissipation of energy spread in time. Nevertheless, it seems to be possible to define such measures (at least in the case of dissipative solutions) as two-dimensional objects dependent on both time and space [27].

Example 2.9. Due to equality (7), the measures of accretion and dissipation for the peakon-antipeakon interaction from Example 1.3 are given by

$$\begin{aligned}\mu^-(T, dx) &= -2H_0^2\delta_0(dx), \\ \mu^+(T, dx) &= 2H_0^2\delta_0(dx), \quad (\text{conservative prolongation}) \\ \mu^+(T, dx) &= 0. \quad (\text{dissipative prolongation})\end{aligned}$$

For $t_0 \neq T$, on the other hand, we have, due to smooth evolution in the neighbourhood of t_0 , that $\mu^+(t_0, dx) = \mu^-(t_0, dx) = 0$. Note also that for conservative solutions $\mu^+(T, dx) + \mu^-(T, dx) = 0$ and measure $\mu^+(T, dx)$ corresponds to the singular part of the measure used to define a global semigroup of solutions in [2].

The main technical result of this paper, based on meticulous studies of characteristics, states that another characterization of measures μ^\pm is possible. Let us begin with a definition.

Definition 2.10 (Thick pushforward). Let u be a weak solution of (2)-(4). For any Borel set $B \subset \mathbb{R}$ the thick pushforward (pushbackward) of B from t_0 to t with $t > t_0$ ($t < t_0$) is defined as

$$B(t) := \{\alpha(t) : \alpha(\cdot) \text{ is any characteristic of } u \text{ satisfying } \alpha(t_0) \in B\}.$$

Theorem 2.11 (Characterization of accretion measure). Let u be a weak solution of (2)-(4). Let $t_0 \geq 0$ and let B be an arbitrary bounded Borel subset of \mathbb{R} . Then

$$\mu^+(t_0, B) = \lim_{t \rightarrow t_0^+} \int_{B(t)} (u_x^+)^2(t, x)dx - \int_B (u_x^+)^2(t_0, x)dx,$$

where $B(t)$ is the thick pushforward of B from t_0 to t (see Definition 2.10).

Let us mention that if u is a weak solution of (2)-(4) then for every $t_0 > 0$ the function $u^{t_0 b}(t, x) := -u(t_0 - t, x)$ is also a weak solution of (2)-(4). Applying Theorem 2.11 to $u^{t_0 b}$, we obtain a dual characterization of the dissipation measure.

Theorem 2.12 (Characterization of dissipation measure). Let u be a weak solution of (2)-(4). Let $t_0 > 0$ and let B be an arbitrary bounded Borel subset of \mathbb{R} . Then

$$\mu^-(t_0, B) = \int_B (u_x^-)^2(t_0, x)dx - \lim_{t \rightarrow t_0^-} \int_{B(t)} (u_x^-)^2(t, x)dx,$$

where $B(t)$ is the thick pushbackward of B from t_0 to t (see Definition 2.10).

Some of the conclusions from the theory presented above, which give a deeper insight into the structure of weak solutions of the Camassa-Holm equation, are the following.

Theorem 2.13. Let u be a weak solution of (2)-(4). Then for every $t_0 \geq 0$ we have

$$(\mu^+(t_0, dx))^{ac} = 0,$$

where $(\mu^+(t_0, dx))^{ac}$ is the absolutely continuous, with respect to the one-dimensional Lebesgue measure, part of measure $\mu^+(t_0, dx)$. Similarly, for every $t_0 > 0$

$$(\mu^-(t_0, dx))^{ac} = 0.$$

Theorem 2.14. *Let u be a dissipative weak solution of (2)-(4). Then $\mu^+(t, dx) = 0$ for every $t \geq 0$.*

Theorem 2.15. *Let u be a weak solution of (2)-(4). Then $\mu^+(t_0, dx) \neq 0$ for at most countably many $t_0 \geq 0$. Similarly, $\mu^-(t_0, dx) \neq 0$ for at most countably many $t_0 > 0$.*

Theorem 2.16. *Let u be a weak solution of (2)-(4). Suppose $\mu^+(t_0, dx) \neq 0$ for some $t_0 \in [0, \infty)$. Then there exists a weak solution \bar{u} of (2)-(4) such that*

- $\bar{u}(t, \cdot) = u(t, \cdot)$ for $t \in [0, t_0]$,
- $\limsup_{t \rightarrow t_0^+} E(\bar{u}(t, \cdot)) < \liminf_{t \rightarrow t_0^+} E(u(t, \cdot))$,

where $E(u(t, \cdot)) = \int_{\mathbb{R}} \frac{1}{2}(u^2(t, x) + u_x^2(t, x))dx$. Consequently, if $\mu^+(t_0, dx) \neq 0$ for some $t_0 \geq 0$ then u does not dissipate energy at the highest possible rate.

3 Preliminaries on characteristics

In this section we prove some technical results regarding characteristics of solutions of the Camassa-Holm equation, which are crucial in the proof of Theorem 2.11.

Let us begin by recalling a result which asserts that supremum and infimum of a family of characteristics is a characteristic. The formulation below is a slight generalization, to initial points in a bounded set, of [26, Lemma 5.1]. Nevertheless, the proof of Lemma 3.1 below follows exactly the proof of [26, Lemma 5.1] and thus we do not repeat it.

Lemma 3.1. *Let $u : [t_0, T] \times \mathbb{R} \rightarrow \mathbb{R}$ be a locally bounded continuous function. Let $\{x_\alpha\}_{\alpha \in A}$ be a family of functions satisfying, for $t \in [t_0, T]$,*

$$\dot{x}_\alpha(t) = u(t, x_\alpha(t))$$

and such that the set $\{x_\alpha(t_0) : \alpha \in A\}$ is bounded. Then function $y(t) := \sup_{\alpha \in A} x_\alpha(t)$ satisfies $\dot{y}(t) = u(t, y(t))$ and, similarly, function $z(t) := \inf_{\alpha \in A} x_\alpha(t)$ satisfies $\dot{z}(t) = u(t, z(t))$. If u is a weak solution of (2)-(4) then $y(t)$ and $z(t)$ are, by [13, Lemma 3.1], characteristics of u .

Using Lemma 3.1 we can define rightmost and leftmost characteristics. Namely, given a weak solution u of (2)-(4), for every $t_0 \geq 0$ and $\zeta \in \mathbb{R}$ there exist (see [26, Corollary 5.2]) the *rightmost* characteristic $\zeta^r(\cdot)$ and the *leftmost* characteristic $\zeta^l(\cdot)$, which are the unique characteristics defined on $[t_0, \infty)$ satisfying

- $\zeta^r(t_0) = \zeta = \zeta^l(t_0)$,
- $\zeta^l(t) \leq \zeta(t) \leq \zeta^r(t)$ for $t \in [t_0, \infty)$ and every characteristic $\zeta(\cdot)$ with $\zeta(t_0) = \zeta$.

Similarly, for every $\Gamma \in \mathbb{R}$ there exist *rightmost backward* characteristic $\Gamma^{rb}(\cdot)$ and the *leftmost backward* characteristic $\Gamma^{lb}(\cdot)$, which are the unique characteristics defined on $[0, t_0]$ satisfying

- $\Gamma^{rb}(t_0) = \Gamma = \Gamma^{lb}(t_0)$,
- $\Gamma^{lb}(t) \leq \Gamma(t) \leq \Gamma^{rb}(t)$ for $t \in [0, t_0]$ and every characteristic $\Gamma(\cdot)$ with $\Gamma(t_0) = \Gamma$.

We finish the introductory part of this section by presenting a change of variables formula, which is indispensable in the theory developed in this paper. Before, however, let us define a useful family of sets (compare [26, Definition 5.8]).

Definition 3.2.

$$\begin{aligned}
L_{t_0, T}^{unique, N} &:= \{ \zeta \in \mathbb{R} : \zeta(\cdot) \text{ is unique forwards on } [t_0, T], \\
&\quad \zeta(s) \text{ is a Lebesgue point of } u_x(s, \cdot) \text{ for almost every } s \in [t_0, T] \text{ and} \\
&\quad \forall_{\eta \in (\zeta - \frac{1}{N}, \zeta) \cup (\zeta, \zeta + \frac{1}{N}), s \in [t_0, T]} -N \leq \omega(s) \leq N \}, \\
L_{t_0, T}^{unique} &:= \bigcup_{N=1}^{\infty} L_{t_0, T}^{unique, N}, \\
L_T^{unique, N} &:= L_{0, T}^{unique, N}, \\
L_T^{unique} &:= L_{0, T}^{unique},
\end{aligned}$$

where $\omega(s) := \frac{u(s, \eta(s)) - u(s, \zeta(s))}{\eta(s) - \zeta(s)}$, $\eta(\cdot)$ is any characteristic satisfying $\eta(t_0) = \eta$ and unique forwards means that if $\zeta_1(\cdot)$ and $\zeta_2(\cdot)$ are two characteristics satisfying $\zeta_1(t_0) = \zeta_2(t_0) = \zeta$ then $\zeta_1(s) = \zeta_2(s)$ for all $s \geq t_0$.

Remark 3.3. *i) By Fubini theorem the condition that $\zeta(s)$ is a Lebesgue point of $u_x(s, \cdot)$ for almost every $s \in [t_0, T]$ defines a full-measure subset, which however does not have to be Borel. Nevertheless, after removal of a set of measure 0 the sets $L_{t_0, T}^{unique, N}$ become Borel measurable.*

ii) The sets $S_{t_0, T}$ in Proposition 2.2 can be chosen as $S_{t_0, T} = L_{t_0, T}^{unique} \setminus Z_{t_0, T}$ for some sets $Z_{t_0, T}$ of Lebesgue measure 0.

The following proposition is a generalization to arbitrary initial times t_0 of the theory presented in [26, Section 7], with the change of variables formula being a consequence of [18, (6)].

Proposition 3.4 (Change of variables formula). *Fix $t_0 \geq 0$. Let g be a bounded nonnegative Borel measurable function and let $A \subset L_{t_0, t}^{unique}$ be a Borel set. Then*

$$\int_{M_{t-t_0}(A)} g(z) dz = \int_A g(M_{t-t_0}(\zeta)) M'_{t-t_0}(\zeta) d\zeta,$$

where $t \geq t_0$, $M_{t-t_0}(\zeta) := \zeta^l(t)$ and $\zeta^l(\cdot)$ is the leftmost characteristic satisfying $\zeta^l(t_0) = \zeta$. Moreover, for $\zeta \in L_{t_0, t}^{unique}$

$$M'_{t-t_0}(\zeta) = e^{\int_{t_0}^t v(s) ds},$$

where $v(s) = u_x(s, \zeta(s))$.

Remark 3.5. *By considering sets $A \cap L_{t_0, t}^{unique, N}$ instead of A and passing to the limit $N \rightarrow \infty$, it suffices to assume that g in Proposition 3.4 is bounded on $L_{t_0, t}^{unique, N}$ for every $N = 1, 2, \dots$ with a bound possibly dependent on N .*

Now, let us present our new results regarding characteristics. The first one shows that rightmost characteristics converge from the right to a rightmost characteristic.

Lemma 3.6. *Let u be a weak solution of (2)-(4) and fix $t_0 \geq 0$. For every $\beta \in \mathbb{R}$ and $t_1 \in [t_0, \infty)$ we have*

$$\lim_{\delta \rightarrow 0^+} (\beta + \delta)^r(t_1) = \beta^r(t_1),$$

where $\beta^r(\cdot)$ and $(\beta + \delta)^r(\cdot)$ are the rightmost characteristics with $\beta^r(t_0) = \beta$ and $(\beta + \delta)^r(t_0) = \beta + \delta$, respectively.

Proof. Let $\Gamma := \lim_{\delta \rightarrow 0^+} (\beta + \delta)^r(t_1)$. Clearly, $\Gamma \geq \beta^r(t_1)$ since $(\beta + \delta)^r(t_1) \geq \beta^r(t_1)$ for every $\delta > 0$. Suppose $\Gamma > \beta^r(t_1)$. Let $\Gamma^{lb}(t)$ be the leftmost backward characteristic with $\Gamma^{lb}(t_1) = \Gamma$. Suppose $\Gamma^{lb}(t_0) \leq \beta$. Then there would exist a characteristic emanating from β

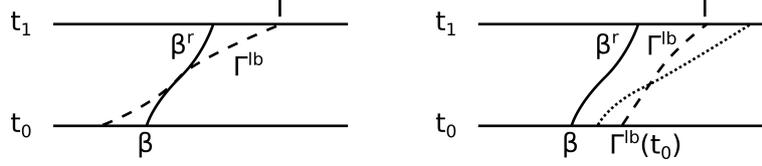


Figure 2: Illustration of the two cases leading to a contradiction in the proof of Lemma 3.6 for $\Gamma > \beta^r(t_1)$. Left: case $\Gamma^{lb}(t_0) \leq \beta$. Then there exists a characteristic emanating from β which is more to the right than β^r . Right: case $\Gamma^{lb}(t_0) > \beta$. Then the rightmost characteristic emanating from a point between β and $\Gamma^{lb}(t_0)$ (dotted line) has to finish at a point to the right of Γ at time t_1 (due to definition of Γ). Thus it has to cross Γ^{lb} , which allows us to find a backward characteristic finishing at Γ , which lies more to the left than Γ^{lb} .

which is more to the right than β^r , obtained as a concatenation of $\beta^r(\cdot)$ (until the crossing time) and $\Gamma^{lb}(\cdot)$ (from the crossing time on) (see Figure 2left).

Hence, $\Gamma^{lb}(t_0) > \beta$. Then, however, the characteristic $[(\Gamma^{lb}(t_0) + \beta)/2]^r(\cdot)$ has to cross $\Gamma^{lb}(\cdot)$ on time interval $[t_0, t_1]$ due to definition of Γ (see Fig. 2right). Hence, there exists a characteristic $\Gamma^b(\cdot)$ such that $\Gamma^b(t_1) = \Gamma$ and $\Gamma^b(\cdot)$ is more to the left than Γ^{lb} (one needs to take a concatenation of $[(\Gamma^{lb}(t_0) + \beta)/2]^r(\cdot)$ and $\Gamma^{lb}(\cdot)$). This gives contradiction. Hence, $\Gamma = \beta^r(t_1)$. \square

The next technical result shows an important convergence property of an integral of $(u_x^+)^2$ in the case when characteristics of a weak solution of the Camassa-Holm equation exhibit behaviour presented in Lemma 3.6.

Lemma 3.7. *Let u be a weak solution of (2)-(4) and fix $0 \leq t_0 < T$. Let $\beta(\cdot)$ be a characteristic of u and let $\{\beta^\delta\}_{\delta \in \Delta}$ be a collection of characteristics of u such that*

- $\Delta \subset (0, \infty)$ has an accumulation point 0,
- for every $t_1 \in [t_0, T)$ the function $\delta \mapsto \beta^\delta(t_1)$ is nondecreasing and

$$\lim_{\delta \rightarrow 0^+} \beta^\delta(t_1) = \beta(t_1). \quad (10)$$

Then

$$\lim_{\delta \rightarrow 0^+} \left(\limsup_{t \rightarrow t_0^+} \int_{(\beta(t), \beta^\delta(t))} (u_x^+)^2(t, x) dx \right) = 0. \quad (11)$$

Similarly, if $\alpha(\cdot)$ is a characteristic of u and $\{\alpha^\delta\}_{\delta \in \Delta}$ is a family of characteristics such that

- $\Delta \subset (0, \infty)$ has an accumulation point 0,
- for every $t_1 \in [t_0, T)$ the function $\delta \mapsto \alpha^\delta(t_1)$ is nonincreasing and

$$\lim_{\delta \rightarrow 0^+} \alpha^\delta(t_1) = \alpha(t_1) \quad (12)$$

then

$$\lim_{\delta \rightarrow 0^+} \left(\limsup_{t \rightarrow t_0^+} \int_{(\alpha^\delta(t), \alpha(t))} (u_x^+)^2(t, x) dx \right) = 0. \quad (13)$$

Proof. We only prove (11), the proof of (13) being analogous.

Take arbitrary $\epsilon > 0$ and $t_1 \in (t_0, T)$ such that

- $t_1 \in \left(t_0, t_0 + \frac{\epsilon}{100(1+2K)} \right)$ where $K := \sup(|u|)$ is a bound on the propagation speed of characteristics,

- $$\Omega(t_1 - t_0) \leq 1, \quad (14)$$

where function $\Omega(\cdot) := \sqrt{LC} \tan(\cdot\sqrt{LC} - \frac{\pi}{2})$ is discussed in [26, Definition 5.5],

- $$\sqrt{LC} \tan(-\sqrt{LC}(t_1 - t_0) + \arctan(1/\sqrt{LC})) > 1/2, \quad (15)$$

where L, C are certain constants dependent only on the $L^\infty([0, \infty), H^1(\mathbb{R}))$ norm of the solution u (see [26]),

- $t_1 - t_0 < \frac{\epsilon}{4(\sup(u^2) + \sup(P))}.$

Let us comment that the third technical condition means that, due to [26, (26)], if $u_x(t, \zeta) > 1$ for some $t \in [t_0, t_1]$ and if $\zeta \in L_{t, t_1}^{unique} \cap S_{t, t_1}$ then $u_x(s, \zeta(s)) > 1/2$ for every $s \in [t, t_1]$ such that $u_x(s, \zeta(s))$ exists. Let us also mention that due $L^\infty([0, \infty), H^1(\mathbb{R}))$ regularity of u , functions u and P are globally bounded and so the quantities $\sup(u^2)$, $\sup(P)$ are well defined.

Let now δ be so small that $\beta^\delta(t_0) - \beta(t_0) < \epsilon/(100(1 + 2K))$ and

$$\int_{(\beta(t_1), \beta^\delta(t_1))} (u_x^+)^2(t_1, x) dx < \epsilon/100. \quad (16)$$

We will show that

$$\int_{(\beta(t), \beta^\delta(t))} (u_x^+)^2(t, x) dx < \epsilon$$

for every $t \in (t_0, t_1)$, see Fig. 3.

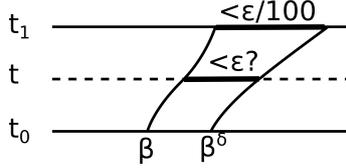


Figure 3: Illustration of the proof of Lemma 3.7. We prove that if t_1 is sufficiently close to t_0 and δ is small enough then $\int_{(\beta(t_1), \beta^\delta(t_1))} (u_x^+)^2(t_1, x) dx < \epsilon/100$ implies $\int_{(\beta(t), \beta^\delta(t))} (u_x^+)^2(t, x) dx < \epsilon$ for every $t \in [t_0, t_1]$.

Suppose the contrary. Then there exists $t \in (t_0, t_1)$ such that

$$\int_{(\beta(t), \beta^\delta(t))} (u_x^+)^2(t, x) dx \geq \epsilon.$$

Due to finite propagation speed of characteristics ($K = \sup(|u|)$) we have

$$\beta^\delta(t) - \beta(t) < \beta^\delta(t_0) - \beta(t_0) + 2K\epsilon/(100(1 + 2K)) < \epsilon/100.$$

Hence,

$$\int_{(\beta(t), \beta^\delta(t)) \cap \{u_x^+(t, \cdot) \geq 1\}} (u_x^+)^2(t, x) dx \geq (99/100)\epsilon.$$

Moreover, $|\{u_x^+(t, \cdot) \geq 1\} \setminus I_{t_1-t}| = 0$ due to (14), where I_{t_1-t} is defined in [26, Definition 5.5], and hence by [26, Lemma 5.7] and [26, Lemma 5.9] we can find N so big that

$$\int_{(\beta(t), \beta^\delta(t)) \cap \{u_x^+(t, \cdot) \geq 1\} \cap L_{t, t_1}^{unique, N} \cap S_{t, t_1}} (u_x^+)^2(t, x) dx \geq (98/100)\epsilon.$$

Using now (15) we obtain that if $\zeta \in (\beta(t), \beta^\delta(t)) \cap \{u_x^+(t, \cdot) \geq 1\} \cap L_{t, t_1}^{unique, N} \cap S_{t, t_1}$ then $u_x(t_1, \zeta(t_1)) > 1/2$ and hence, by [26, Proposition 7.1],

$$e^{-\epsilon} \leq e^{-\frac{2(t_1-t_0)}{1/2}((\sup u)^2 + \sup(P))} \leq \frac{v^2(t_1)M'_{t_1-t}(\zeta)}{v^2(t)},$$

where $v(t) = u_x(t, \zeta(t))$ and $M_{t_1-t}(\zeta) = \zeta^l(t_1)$ is the leftmost characteristic satisfying $\zeta^l(t) = \zeta$, which for the range of ζ that are of interest to us are in fact unique characteristics.

Using now the change of variables formula from Proposition 3.4 we obtain

$$\begin{aligned} \int_{(\beta(t_1), \beta^\delta(t_1))} (u_x^+)^2(t_1, z) dz &\geq \int_{M_{t_1-t}((\beta(t), \beta^\delta(t)) \cap \{u_x^+(t, \cdot) \geq 1\} \cap L_{t, t_1}^{unique, N} \cap S_{t, t_1})} (u_x^+)^2(t_1, z) dz \\ &= \int_{(\beta(t), \beta^\delta(t)) \cap \{u_x^+(t, \cdot) \geq 1\} \cap L_{t, t_1}^{unique, N} \cap S_{t, t_1}} (u_x^+)^2(t_1, M_{t_1-t}(\zeta)) M'_{t_1-t}(\zeta) d\zeta \\ &\geq \int_{(\beta(t), \beta^\delta(t)) \cap \{u_x^+(t, \cdot) \geq 1\} \cap L_{t, t_1}^{unique, N} \cap S_{t, t_1}} (u_x^+)^2(t, \zeta) e^{-\epsilon} d\zeta \\ &\geq 98/100\epsilon e^{-\epsilon} \geq 1/100\epsilon, \end{aligned}$$

which contradicts (16). \square

Corollary 3.8. *Let u be a weak solution of (2)-(4). Then for every $t_0 \geq 0$ and every $\alpha, \beta \in \mathbb{R}$ we have*

$$\lim_{\delta \rightarrow 0^+} \left(\limsup_{t \rightarrow t_0^+} \int_{(\beta^r(t), (\beta+\delta)^r(t))} (u_x^+)^2(t, x) dx \right) = 0, \quad (17)$$

$$\lim_{\delta \rightarrow 0^+} \left(\limsup_{t \rightarrow t_0^+} \int_{[(\alpha-\delta)^l(t), \alpha^l(t)]} (u_x^+)^2(t, x) dx \right) = 0. \quad (18)$$

Proof. To prove (17) take $\beta(\cdot) \equiv \beta^r(\cdot)$ and $\beta^\delta(\cdot) \equiv (\beta + \delta)^r(\cdot)$ and, using Lemma 3.6 apply Lemma 3.7. (18) follows in a similar fashion. \square

Lemma 3.7 is a direct consequence of an even more general result contained in Lemma 3.9, which we now formulate and prove. Let us note that the proof of Lemma 3.9 follows the lines of the proof of Lemma 3.7 with a few minor modifications. Since, however, Lemma 3.9 is much more abstract, we decided to keep both proofs for clarity of the exposition. Importantly, the formulation of Lemma 3.7 will not always be sufficient for our purposes, and on one occasion we will have to resort to the general formulation given in Lemma 3.9.

Lemma 3.9. *Let u be a weak solution of (2)-(4) and fix t_0, T , which satisfy $0 \leq t_0 < T$. Let $\{\mathcal{B}^\delta\}_{\delta \in \Delta}$ be a family of collections of characteristics on $[t_0, T]$ satisfying*

- $\Delta \subset (0, \infty)$ has an accumulation point 0,
- $\mathcal{B}^{\delta_1} \subset \mathcal{B}^{\delta_2}$ for $\delta_1 < \delta_2$,
- for every $t \in [t_0, T]$ the set $B^\delta(t) := \{\beta(t) : \beta(\cdot) \in \mathcal{B}^\delta\}$ is Borel measurable,
- for every $t \in [t_0, T]$

$$\lim_{\delta \rightarrow 0^+} \int_{B^\delta(t)} (u_x^+)^2(t, x) = 0,$$

- there exists $D > 0$ such that for every $t \in [t_0, T]$ we have

$$\text{diam}(B^\delta(t)) < D. \quad (19)$$

Then

$$\lim_{\delta \rightarrow 0^+} \left(\limsup_{t \rightarrow t_0^+} \int_{B^\delta(t)} (u_x^+)^2(t, x) dx \right) = 0.$$

Proof. Take arbitrary $\epsilon > 0$ and $t_1 \in (t_0, T)$ such that

- $t_1 \in \left(t_0, t_0 + \frac{\epsilon}{2(\sup(u^2) + \sup(P))} \left(\frac{\epsilon}{200D} \right)^{1/2} \right),$
- $$\Omega(t_1 - t_0) \leq \left(\frac{\epsilon}{100D} \right)^{1/2}, \quad (20)$$

with Ω defined in [26, Definition 5.5]

- $\sqrt{LC} \tan \left(-\sqrt{LC}(t_1 - t_0) + \arctan((\epsilon/(100D))^{1/2}/\sqrt{LC}) \right) > (\epsilon/(200D))^{1/2},$ where L, C are constants dependent only on the $L^\infty([0, \infty), H^1(\mathbb{R}))$ norm of the solution u (see [26]).

Finally, take δ so small that

$$\int_{B^\delta(t_1)} (u_x^+)^2(t_1, x) dx < \epsilon/100.$$

We will show that

$$\int_{B^\delta(t)} (u_x^+)^2(t, x) dx < \epsilon$$

for every $t \in (t_0, t_1)$. Suppose the contrary. Then there exists $t \in (t_0, t_1)$ such that

$$\int_{B^\delta(t)} (u_x^+)^2(t, x) dx \geq \epsilon.$$

Since, however, $|B^\delta(t)| < D$ we conclude

$$\int_{B^\delta(t) \cap \{(u_x^+)^2(t, \cdot) \geq \epsilon/(100D)\}} (u_x^+)^2(t, x) dx \geq (99/100)\epsilon.$$

Moreover, $|\{(u_x^+)^2(t, \cdot) \geq \epsilon/(100D)\} \setminus I_{t_1-t}| = 0$ due to (20) and hence by [26, Lemma 5.7] we can find N so big that

$$\int_{B^\delta(t) \cap \{(u_x^+)^2(t, \cdot) \geq \epsilon/(100D)\} \cap L_{t, t_1}^{unique, N} \cap S_{t, t_1}} (u_x^+)^2(t, x) dx \geq (98/100)\epsilon.$$

Using now the fact that if $\zeta \in B^\delta(t) \cap \{(u_x^+)^2(t, \cdot) \geq \epsilon/(100D)\} \cap L_{t, t_1}^{unique, N} \cap S_{t, t_1}$ then $u_x(t_1, \zeta(t_1)) > (\epsilon/(200D))^{1/2}$ (for a.e. ζ) we obtain, by [26, Proposition 7.1],

$$e^{-\epsilon} \leq e^{-\frac{2(t_1-t_0)}{(\epsilon/(200D))^{1/2}(\sup u^2 + \sup(P))}} \leq \frac{v^2(t_1)M'_{t_1-t}(\zeta)}{v^2(t)}.$$

Using now the change of variables formula from Proposition 3.4 we obtain

$$\begin{aligned} \int_{B^\delta(t_1)} (u_x^+)^2(t_1, z) dz &\geq \int_{M_{t_1-t}(B^\delta(t) \cap \{(u_x^+)^2(t, \cdot) \geq \epsilon/(100D)\} \cap L_{t, t_1}^{unique, N} \cap S_{t, t_1})} (u_x^+)^2(t_1, z) dz \\ &= \int_{B^\delta(t) \cap \{(u_x^+)^2(t, \cdot) \geq \epsilon/(100D)\} \cap L_{t, t_1}^{unique, N} \cap S_{t, t_1}} (u_x^+)^2(t_1, M_{t_1-t}(\zeta)) M'_{t_1-t}(\zeta) d\zeta \\ &\geq \int_{B^\delta(t) \cap \{(u_x^+)^2(t, \cdot) \geq \epsilon/(100D)\} \cap L_{t, t_1}^{unique, N} \cap S_{t, t_1}} (u_x^+)^2(t, \zeta) e^{-\epsilon} d\zeta \\ &\geq (98/100)\epsilon e^{-\epsilon} \geq \epsilon/100, \end{aligned}$$

which gives contradiction. \square

Remark 3.10. Lemma 3.7 follows from Lemma 3.9 by setting $\mathcal{B}^\delta := \{\zeta(\cdot) : \beta(t) \leq \zeta(t) \leq \beta^\delta(t) \text{ for every } t \in [t_0, T]\}$ and observing that by finite propagation speed of characteristics (10) implies (19).

4 Proof of Theorem 2.11

To prove Theorem 2.11 we show it for the following sequence of classes of $B \subset \mathbb{R}$:

- B – closed interval,
- B – open interval,
- B – arbitrary open set,
- B – arbitrary compact set,
- B – arbitrary Borel set.

B - closed interval

Let $B = [\alpha, \beta]$. Then the thick pushforward of B is given by $B(t) = [\alpha^l(t), \beta^r(t)]$, where $\alpha^l(\cdot)$ is the leftmost and $\beta^r(\cdot)$ the rightmost characteristic emanating from α and β , respectively. To prove Theorem 2.11 for B in this form, it suffices thus to show the following result.

Proposition 4.1. *Let u be a weak solution of (2)-(4). The measure of accretion μ^+ associated to u satisfies*

$$\mu^+(t_0, [\alpha, \beta]) = \lim_{t \rightarrow t_0^+} \int_{[\alpha^l(t), \beta^r(t)]} (u_x^+)^2(t, x) dx - \int_{[\alpha, \beta]} (u_x^+)^2(t_0, x) dx \quad (21)$$

for every $\alpha, \beta \in \mathbb{R}$, $\alpha \leq \beta$, and $t_0 \in [0, \infty)$.

Proof. By regularity of the measure μ^+ , it suffices to show that

$$\lim_{t \rightarrow t_0^+} \int_{[\alpha^l(t), \beta^r(t)]} (u_x^+)^2(t, x) dx = \lim_{\epsilon \rightarrow 0^+} \lim_{t \rightarrow t_0^+} \int_{\mathbb{R}} \phi^\epsilon(x) (u_x^+)^2(t, x) dx, \quad (22)$$

where (see Fig. 4left)

$$\phi^\epsilon(x) = \begin{cases} 1 & \text{if } x \in [\alpha - \epsilon, \beta + \epsilon], \\ 0 & \text{if } x \leq \alpha - 2\epsilon \text{ or } x \geq \beta + 2\epsilon, \\ \frac{1}{\epsilon}(x - (\alpha - 2\epsilon)) & \text{if } x \in (\alpha - 2\epsilon, \alpha - \epsilon), \\ -\frac{1}{\epsilon}(x - (\beta + 2\epsilon)) & \text{if } x \in (\beta + \epsilon, \beta + 2\epsilon). \end{cases}$$

Indeed, by definition of measure μ^+ (Definition 2.5) we have

$$\int_{\mathbb{R}} \phi^\epsilon(x) d\mu^+(t_0, dx) := \lim_{t \rightarrow t_0^+} \int_{\mathbb{R}} \phi^\epsilon(x) (u_x^+)^2(t, x) dx - \int_{\mathbb{R}} \phi^\epsilon(x) (u_x^+)^2(t_0, x) dx.$$

Passage to the limit $\epsilon \rightarrow 0$ in combination with outer regularity of the measure μ^+ and (22) leads to (21).

To prove (22) we observe that by finite propagation speed of characteristics we have $\phi^\epsilon \geq \mathbf{1}_{[\alpha^l(t), \beta^r(t)]}$ for t sufficiently close to t_0 (Fig. 4middle). Hence, for every $\epsilon > 0$,

$$\lim_{t \rightarrow t_0^+} \int_{[\alpha^l(t), \beta^r(t)]} (u_x^+)^2(t, x) dx \leq \lim_{t \rightarrow t_0^+} \int_{\mathbb{R}} \phi^\epsilon(x) (u_x^+)^2(t, x) dx.$$

and thus

$$\lim_{t \rightarrow t_0^+} \int_{[\alpha^l(t), \beta^r(t)]} (u_x^+)^2(t, x) dx \leq \lim_{\epsilon \rightarrow 0^+} \lim_{t \rightarrow t_0^+} \int_{\mathbb{R}} \phi^\epsilon(x) (u_x^+)^2(t, x) dx.$$

On the other hand, for every $\delta > 0$ we have (see Fig. 4right)

$$\lim_{\epsilon \rightarrow 0^+} \lim_{t \rightarrow t_0^+} \int_{\mathbb{R}} \phi^\epsilon(x) (u_x^+)^2(t, x) dx \leq \lim_{t \rightarrow t_0^+} \int_{[(\alpha - \delta)^l(t), (\beta + \delta)^r(t)]} (u_x^+)^2(t, x) dx.$$

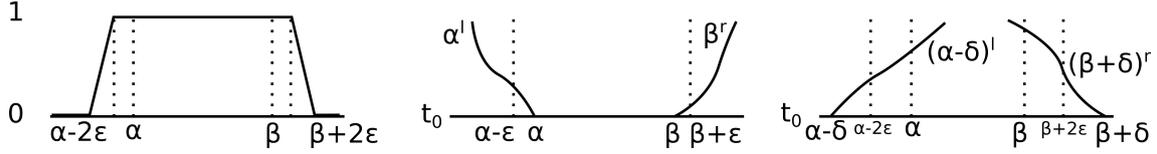


Figure 4: Graph of the function ϕ^ϵ (left). Illustration of the inequality $\phi^\epsilon \geq \mathbf{1}_{[\alpha^l(t), \beta^r(t)]}$ for t sufficiently close to t_0 (middle). Illustration of the fact that for every $\delta > 0$ we have $\phi^\epsilon \leq \mathbf{1}_{[(\alpha-\delta)^l(t), (\beta+\delta)^r(t)]}$ for ϵ small enough and t sufficiently close to t_0 (right).

By Corollary 3.8 we can pass to the limit $\delta \rightarrow 0^+$ to obtain

$$\lim_{\epsilon \rightarrow 0^+} \lim_{t \rightarrow t_0^+} \int_{\mathbb{R}} \phi^\epsilon(x) (u_x^+)^2(t, x) dx \leq \lim_{t \rightarrow t_0^+} \int_{[\alpha^l(t), \beta^r(t)]} (u_x^+)^2(t, x) dx.$$

Let us comment that the limits on the right-hand side of the last two inequalities exist by [26, Theorem 2.7]. \square

Remark 4.2. Proposition 4.1 shows that the 'canonical' choice of characteristics in [26, Theorem 2.8], resulting in a measure is the leftmost characteristic emanating from α and the rightmost characteristic emanating from β .

Corollary 4.3 (of Proposition 4.1). *Let u be a weak solution of (2)-(4). The measure μ^+ satisfies for every $\alpha, \beta \in \mathbb{R}$, $\alpha < \beta$:*

$$\begin{aligned} \mu^+(t_0, \{\alpha\}) &= \lim_{t \rightarrow t_0^+} \int_{[\alpha^l(t), \alpha^r(t)]} (u_x^+)^2(t, x) dx, \\ \mu^+(t_0, (\alpha, \beta)) &= \lim_{t \rightarrow t_0^+} \int_{(\alpha^r(t), \beta^l(t))} (u_x^+)^2(t, x) dx - \int_{(\alpha, \beta)} (u_x^+)^2(t_0, x) dx. \end{aligned}$$

B - open interval

Let us begin by introducing the notions of *left-rightmost characteristic* and *right-leftmost characteristic*.

Definition 4.4. *Let A_ζ be a family of characteristics given by*

$$A_\zeta := \{\zeta(\cdot) : \zeta(t_0) = \zeta \text{ and there exists } t_1 > t_0 \text{ such that } \zeta(t) = \zeta^r(t) \text{ for every } t \in [t_0, t_1]\}.$$

The left-rightmost characteristic, ζ^{rl} , is defined, using Lemma 3.1, by (see Fig. 5)

$$\zeta^{rl}(t) := \inf\{\zeta(t) : \zeta(\cdot) \in A_\zeta\}.$$

Similarly, let \hat{A}_ζ be a family of characteristics given by

$$\hat{A}_\zeta := \{\zeta(\cdot) : \zeta(t_0) = \zeta \text{ and there exists } t_1 > t_0 \text{ such that } \zeta(t) = \zeta^l(t) \text{ for every } t \in [t_0, t_1]\}$$

Then right-leftmost characteristic, ζ^{lr} , is defined by

$$\zeta^{lr}(t) := \sup\{\zeta(t) : \zeta(\cdot) \in \hat{A}_\zeta\}.$$

Lemma 4.5. *Let u be a weak solution of (2)-(4) and fix $t_0 \geq 0$. Then for every $\zeta \in \mathbb{R}$ we have*

$$\lim_{t \rightarrow t_0^+} \int_{[\zeta^{rl}(t), \zeta^r(t)]} (u_x^+)^2(t, x) dx = 0. \quad (23)$$

and, similarly,

$$\lim_{t \rightarrow t_0^+} \int_{[\zeta^l(t), \zeta^{lr}(t)]} (u_x^+)^2(t, x) dx = 0. \quad (24)$$

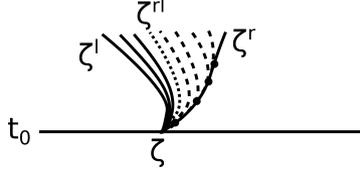


Figure 5: Illustration of the concept of the left-rightmost characteristic. The left-rightmost characteristic $\zeta^{rl}(\cdot)$, dotted line, is the infimum of the family (dashed lines) of characteristics which coincide with the rightmost characteristic $\zeta^r(\cdot)$ up to some timepoint $t > t_0$. Since there might exist characteristics emanating from ζ which are more to the left than ζ^{rl} , the left-rightmost characteristic does not in general coincide with the leftmost characteristic.

Remark 4.6. *The convergences from Lemma 4.5 hold in stark contrast to*

$$\lim_{t \rightarrow t_0^+} \int_{[\zeta^l(t), \zeta^r(t)]} (u_x^+)^2(t, x) dx = \mu^+(\{\zeta\}),$$

which in general does not vanish.

Proof of Lemma 4.5. Define

$$I := \inf\{t : t > t_0 \text{ and there exists a characteristic } \zeta(\cdot) \in A_\zeta \text{ such that } \zeta(t) < \zeta^r(t)\},$$

where A_ζ is given by Definition 4.4. If $I > t_0$ then the proof is trivial, since $(\zeta^{rl}(t), \zeta^r(t)) = \emptyset$ for all $t < I$. So, suppose $I = t_0$. Then there exists a sequence δ_n such that $\delta_n \rightarrow 0$ as $n \rightarrow \infty$ and for every $n > 0$ there exists a characteristic $\zeta(\cdot) \in A_\zeta$ such that

$$\delta_n = \inf\{t - t_0 : \zeta(t) < \zeta^r(t)\}.$$

In other words, timepoint $t_0 + \delta_n$ is a branching point for some characteristic $\zeta(\cdot) \in A_\zeta$. Let β^{δ_n} be the characteristic defined by

$$\beta^{\delta_n}(\cdot) := \inf\{\zeta(\cdot) : \zeta(t) = \zeta^r(t) \text{ for every } t \in [t_0, t_0 + \delta_n]\}.$$

We interpret β^{δ_n} as the leftmost characteristic from the point of branching. With such choice of β^{δ_n} and $\Delta := \{\delta_n\}_{n=1}^\infty$ as well as $\beta(\cdot) \equiv \zeta^{rl}(\cdot)$ the assumptions of Lemma 3.7 are satisfied. Using Lemma 3.7 and the property

$$\int_{[\zeta^{rl}(t), \zeta^r(t)]} (u_x^+)^2(t, x) dx = \int_{[\zeta^{rl}(t), \zeta^{\delta_n}(t)]} (u_x^+)^2(t, x) dx$$

for $t \in [t_0, t_0 + \delta_n]$ we infer (23). The proof of (24) is analogous. \square

Lemma 4.7. *Let u be a weak solution of (2)-(4). Let $B = (\alpha, \beta)$ be a bounded open interval. Then for t sufficiently close to t_0 (so that $\alpha^r(\cdot)$ and $\beta^l(\cdot)$ do not cross)*

$$(\alpha^r(t), \beta^l(t)) \subset B(t) \subset [\alpha^{rl}(t), \beta^{lr}(t)],$$

where $B(t)$ is the thick pushforward, α^r is the rightmost characteristic with $\alpha^r(t_0) = \alpha$, β^l is the leftmost characteristic with $\beta^l(t_0) = \beta$, α^{rl} is the left-rightmost characteristic with $\alpha^{rl}(t_0) = \alpha$ and β^{lr} is the right-leftmost characteristic satisfying $\beta^{lr}(t_0) = \beta$.

Proof. If a characteristic emanating from $B = (\alpha, \beta)$ crosses α^r or β^l , it does so at a timepoint strictly bigger than t_0 . This implies the second inclusion. The first inclusion can be proven as follows. Take $\zeta \in (\alpha^r(t), \beta^l(t))$ and consider the leftmost backward characteristic ζ^{lb} with $\zeta^{lb}(t) = \zeta$. Clearly, $\zeta^{lb}(\cdot)$ does not cross $\alpha^r(\cdot)$ on $[t_0, t]$ due to the rightmost property of $\alpha^r(\cdot)$. By the same token, $\zeta^{lb}(\cdot)$ does not cross $\beta^l(\cdot)$ on $[t_0, t]$. Hence, $\alpha < \zeta^{lb}(t_0) < \beta$ and thus $\zeta = \zeta^{lb}(t) \in B(t)$. \square

Proposition 4.8. *Let u be a weak solution of (2)-(4) and let $B = (\alpha, \beta)$ be an open interval. Fix $t_0 \geq 0$. Then*

$$\lim_{t \rightarrow t_0^+} \int_{(\alpha^r(t), \beta^l(t))} (u_x^+)^2(t, x) dx = \lim_{t \rightarrow t_0^+} \int_{B(t)} (u_x^+)^2(t, x) dx, \quad (25)$$

where $B(t)$ is the thick pushforward. Consequently,

$$\mu^+(t_0, (\alpha, \beta)) = \lim_{t \rightarrow t_0^+} \int_{B(t)} (u_x^+)^2(t, x) dx - \int_{(\alpha, \beta)} (u_x^+)^2(t_0, x) dx. \quad (26)$$

Proof. We first prove that

$$\lim_{t \rightarrow t_0^+} \int_{(\alpha^r(t), \beta^l(t))} (u_x^+)^2(t, x) dx \geq \limsup_{t \rightarrow t_0^+} \int_{B(t)} (u_x^+)^2(t, x) dx.$$

Indeed, $B(t) \subset [\alpha^{rl}(t), \beta^{lr}(t)]$ by Lemma 4.7. Hence,

$$\limsup_{t \rightarrow t_0^+} \int_{B(t)} (u_x^+)^2(t, x) dx \leq \limsup_{t \rightarrow t_0^+} \int_{[\alpha^{rl}(t), \beta^{lr}(t)]} (u_x^+)^2(t, x) dx = \lim_{t \rightarrow t_0^+} \int_{(\alpha^r(t), \beta^l(t))} (u_x^+)^2(t, x) dx,$$

where the last equality follows by Lemma 4.5. The reverse inequality

$$\lim_{t \rightarrow t_0^+} \int_{(\alpha^r(t), \beta^l(t))} (u_x^+)^2(t, x) dx \leq \liminf_{t \rightarrow t_0^+} \int_{B(t)} (u_x^+)^2(t, x) dx$$

is clear since $(\alpha^r(t), \beta^l(t)) \subset B(t)$ by Lemma 4.7. This proves (25). Formula (26) follows by Corollary 4.3. \square

B - arbitrary open set

Proposition 4.9. *Let u be a weak solution of (2)-(4) and let $B = \bigcup_{n=1}^{\infty} (\alpha_n, \beta_n)$ be an arbitrary bounded open subset in \mathbb{R} . Then*

$$\mu^+(t_0, B) = \lim_{t \rightarrow t_0^+} \int_{B(t)} (u_x^+)^2(t, x) dx - \int_B (u_x^+)^2(t_0, x) dx, \quad (27)$$

where $B(t)$ is the thick pushforward of the set B . Similarly,

$$\mu^+(t_0, B) = \lim_{t \rightarrow t_0^+} \int_{\bigcup_{n=1}^{\infty} (\alpha_n^r(t), \beta_n^l(t))} (u_x^+)^2(t, x) dx - \int_B (u_x^+)^2(t_0, x) dx. \quad (28)$$

Proof. Recall that by Lemma 3.1 the suprema and infima of families of characteristics are again characteristics. Using this property we define characteristics $\alpha_n(t)$ and $\beta_n(t)$ recursively as follows.

1. $\alpha_1(t) := \inf_{\alpha \in (\alpha_1, \beta_1)} \{\alpha^l(t) : \alpha^l(t_0) = \alpha\}$.
2. $\beta_1(t) := \sup_{\beta \in (\alpha_1, \beta_1)} \{\beta^r(t) : \beta^r(t_0) = \beta\}$.
3. To define $\alpha_n(\cdot)$ let

$$\alpha_n^*(t) := \inf_{\alpha \in (\alpha_n, \beta_n)} \{\alpha^l(t) : \alpha^l(t_0) = \alpha\}$$

Suppose there exist $n_*, n_{**} < n$ such that $\beta_{n_*} \leq \alpha_n < \alpha_{n_{**}}$ and

$$\begin{aligned} \beta_{n_*} &= \max\{\beta_{n'} : n' < n, \beta_{n'} \leq \alpha_n\}, \\ \alpha_{n_{**}} &= \min\{\alpha_{n'} : n' < n, \alpha_{n'} > \alpha_n\}. \end{aligned}$$

Let

$$\begin{aligned} t_* &:= \inf\{t > t_0 : \alpha_n^*(t) < \beta_{n_*}(t)\}, \\ t_{**} &:= \inf\{t > t_0 : \alpha_n^*(t) = \alpha_{n_{**}}(t)\}, \end{aligned}$$

where both t_* and t_{**} can assume the value $+\infty$. Define (see Fig. 6)

$$\alpha_n(t) := \begin{cases} \alpha_n^*(t) & \text{for } t \leq \min(t_*, t_{**}), \\ \begin{cases} \beta_{n_*}(t) & \text{if } t_* \leq t_{**} \\ \alpha_{n_{**}}(t) & \text{if } t_* > t_{**} \end{cases} & \text{for } t > \min(t_*, t_{**}). \end{cases}$$

The same defining formula holds if n_* does not exist (then we put $t_* = \infty$ and β_{n_*} is irrelevant) or n_{**} does not exist (then we put $t_{**} = \infty$ and $\alpha_{n_{**}}$ is irrelevant).

4. To define $\beta_n(\cdot)$ let

$$\beta_n^*(t) := \sup_{\beta \in (\alpha_n, \beta_n)} \{\beta^r(t) : \beta^r(t_0) = \beta\}.$$

Suppose there exists $n_{**} < n$ such that $\beta_n \leq \alpha_{n_{**}}$ and

$$\alpha_{n_{**}} = \min\{\alpha_{n'} : n' < n, \alpha_{n'} \geq \beta_n\}.$$

Let

$$t_{**} := \inf\{t > t_0 : \beta_n^*(t) > \alpha_{n_{**}}(t)\},$$

where $\alpha_n(t)$ has been defined in the previous step and both t_* and t_{**} can assume the value $+\infty$. Define (see Fig. 6)

$$\beta_n(t) := \begin{cases} \beta_n^*(t) & \text{for } t \leq t_{**}, \\ \alpha_{n_{**}}(t) & \text{for } t > t_{**}. \end{cases}$$

The same defining formula holds if n_{**} does not exist (then we put $t_{**} = \infty$ and $\alpha_{n_{**}}$ is irrelevant).

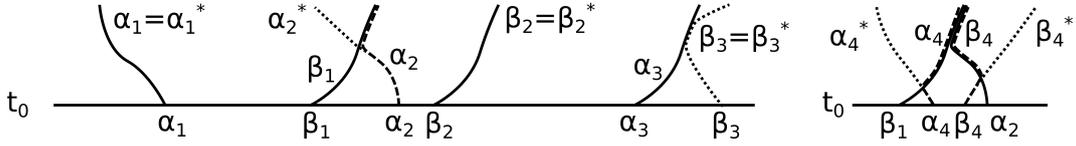


Figure 6: Illustration of the process of definition of characteristics α_n, β_n . Characteristics α_1 and β_1 are defined without any intersection and thus coincide with α_1^* and β_1^* , respectively. Characteristic α_2^* (dotted line) crosses β_1 and thus α_2 follows, from the timepoint of intersection on, characteristic β_1 . Characteristic β_n can also cross α_n , which is the case for β_3 . In this case, however β_n does not follow α_n , yet remains β_n^* . An example of a more complex situation when both α_n and β_n cross other characteristics is shown in the right panel for α_4 and β_4 . Here α_4, β_4 follow, beginning from some timepoint, the characteristics β_1 and α_2 , respectively. As β_1 and α_2 coincide from some timepoint on, so do α_4 and β_4 . For clarity of the picture, the intersections of characteristics are depicted as transversal. In fact, however, the characteristics can only intersect tangentially.

The definition is designed in such a way that if $\alpha_{n_1} < \beta_{n_1} \leq \alpha_{n_2} < \beta_{n_2}$ for some $1 \leq n_1, n_2 \leq N$ then

$$\alpha_{n_1}(t) < \beta_{n_1}(t) \leq \alpha_{n_2}(t) < \beta_{n_2}(t)$$

for every $t \geq t_0$ and

$$\bigcup_{n=1}^{\infty} (\alpha_n(t), \beta_n(t)) \subset B(t) \subset \bigcup_{n=1}^{\infty} [\alpha_n(t), \beta_n(t)]$$

(compare Lemma 4.7). Hence, up to a set of measure 0 we have

$$B(t) = \bigcup_{n=1}^{\infty} (\alpha_n(t), \beta_n(t)).$$

and the union above is disjoint. Take now $B^{1/N} := \bigcup_{n=N}^{\infty} \{\gamma(\cdot) : \gamma(t_0) \in (\alpha_n, \beta_n) \text{ and } \gamma(t) \in [\alpha_n(t), \beta_n(t)] \text{ for every } t \in (t_0, \infty)\}$. Observe that the family $\{B^\delta\}_{\delta \in \Delta}$ with $\Delta = \{1/N : N = 1, 2, \dots\}$ fulfils the assumptions of Lemma 3.9. Hence,

$$\lim_{N \rightarrow \infty} \left(\limsup_{t \rightarrow t_0^+} \int_{B^{1/N}(t)} (u_x^+)^2(t, x) dx \right) = 0. \quad (29)$$

Let now

$$B_N := \bigcup_{n=1}^N (\alpha_n, \beta_n).$$

For every $n \in \{1, 2, \dots, N\}$ and every $t \geq t_0$ we have

$$(\alpha_n^r(t), \beta_n^l(t)) \subset (\alpha_n(t), \beta_n(t)) \subset (\alpha_n, \beta_n)(t) \subset [\alpha_n^{rl}(t), \beta_n^{lr}(t)],$$

where the first inclusion follows from the definition of rightmost and leftmost characteristics, the second inclusion follows by the definition of $\alpha_n(t), \beta_n(t)$ and the last inclusion is due to Lemma 4.7. Consequently, by Proposition 4.8 and Lemma 4.5,

$$\mu^+(t_0, (\alpha_n, \beta_n)) = \lim_{t \rightarrow t_0^+} \int_{(\alpha_n(t), \beta_n(t))} (u_x^+)^2(t, x) dx - \int_{(\alpha_n, \beta_n)} (u_x^+)^2(t_0, x) dx.$$

Hence, for every $N < \infty$

$$\mu^+(t_0, B_N) = \lim_{t \rightarrow t_0^+} \int_{\bigcup_{n=1}^N (\alpha_n(t), \beta_n(t))} (u_x^+)^2(t, x) dx - \int_{B_N} (u_x^+)^2(t_0, x) dx.$$

Passing to the limit $N \rightarrow \infty$ we obtain:

$$\lim_{N \rightarrow \infty} \lim_{t \rightarrow t_0^+} \int_{\bigcup_{n=1}^N (\alpha_n(t), \beta_n(t))} (u_x^+)^2(t, x) dx = \mu^+(t_0, B) + \int_B (u_x^+)^2(t_0, x) dx.$$

Moreover, by (29) we have

$$\lim_{N \rightarrow \infty} \left(\limsup_{t \rightarrow t_0^+} \int_{\bigcup_{n=N+1}^{\infty} (\alpha_n(t), \beta_n(t))} (u_x^+)^2(t, x) dx \right) = 0.$$

Fix $\epsilon > 0$. By the two above convergences there exists N_ϵ so large that there exists $t_\epsilon > t_0$ such that for $t \in [t_0, t_\epsilon]$

$$\left| \int_{\bigcup_{n=1}^{N_\epsilon} (\alpha_n(t), \beta_n(t))} (u_x^+)^2(t, x) dx - \left[\mu^+(t_0, B) + \int_B (u_x^+)^2(t_0, x) dx \right] \right| < \epsilon/2,$$

$$\int_{\bigcup_{n=N_\epsilon+1}^{\infty} (\alpha_n(t), \beta_n(t))} (u_x^+)^2(t, x) dx < \epsilon/2.$$

Hence,

$$\left| \int_{\bigcup_{n=1}^{\infty} (\alpha_n(t), \beta_n(t))} (u_x^+)^2(t, x) dx - \left[\mu^+(t_0, B) + \int_B (u_x^+)^2(t_0, x) dx \right] \right| < \epsilon,$$

which means that for every $\epsilon > 0$ there exists $t_\epsilon > t_0$ such that for $t \in [t_0, t_\epsilon]$

$$\left| \int_{B(t)} (u_x^+)^2(t, x) dx - \left[\mu^+(t_0, B) + \int_B (u_x^+)^2(t_0, x) dx \right] \right| < \epsilon.$$

This concludes the proof of (27). The proof of (28) is similar and in fact simpler since the sets $(\alpha_n^r(t), \beta_n^l(t))$ are pairwise disjoint for different n . \square

Remark 4.10. Characteristic α_n^* can be interpreted as follows. Let

$$\tau := \inf\{t > t_0 : \text{there exists } \alpha \in (\alpha_n, \beta_n) \text{ such that } \alpha^l(t) = \alpha_n^r(t)\}.$$

There are now two possibilities. Either there exists $\alpha \in (\alpha_n, \beta_n)$ such that $\alpha(\tau) = \alpha_n^r(\tau)$ and then

$$\alpha_n^*(t) = \begin{cases} \alpha_n^r(t) & \text{for } t_0 \leq t \leq \tau, \\ (\alpha_n^r(\tau))^l(t) & \text{for } t > \tau \end{cases}$$

or, otherwise,

$$\alpha_n^*(t) = \begin{cases} \alpha_n^r(t) & \text{for } t_0 \leq t \leq \tau, \\ (\alpha_n^r(\tau))^{rl}(t) & \text{for } t > \tau. \end{cases}$$

B - arbitrary compact set

Lemma 4.11. Let u be a weak solution of (2)-(4) and let $K \subset \mathbb{R}$ be an arbitrary compact set. Fix t_0, T such that $0 \leq t_0 < T$. Let $U = (\underline{U}, \overline{U})$ be a bounded open interval satisfying $K \subset U$ and so large that

$$\begin{cases} [\sup(K) + 2(T - t_0) \sup(|u|)] \in U, \\ [\inf(K) - 2(T - t_0) \sup(|u|)] \in U. \end{cases} \quad (30)$$

Let the open set $U \setminus K$ have the representation

$$U \setminus K = \bigcup_{n=1}^{\infty} (\alpha_n, \beta_n),$$

where the open intervals (α_n, β_n) are pairwise disjoint. Then for $t \in [t_0, T]$

$$K(t) = (\underline{U}^r(t), \overline{U}^l(t)) \setminus \bigcup_{n=1}^{\infty} (\alpha_n^r(t), \beta_n^l(t)), \quad (31)$$

where we adopt the convention $(\alpha_n^r(t), \beta_n^l(t)) = \emptyset$ if $\alpha_n^r(t) > \beta_n^l(t)$.

Proof. First we prove the inclusion \subset in (31). Let $x \in K(t)$. Condition (30) guarantees that

$$K(t) \subset (\underline{U}^r(t), \overline{U}^l(t))$$

for $t \in [t_0, T]$ and hence $x \in (\underline{U}^r(t), \overline{U}^l(t))$. Fix now n . We will show that $x \notin (\alpha_n^r(t), \beta_n^l(t))$. Indeed, since $x \in K(t)$ there exists $\zeta \in K$ such that $\zeta(t) = x$. As $(\alpha_n, \beta_n) \cap K = \emptyset$ then either $\beta_n \leq \zeta$ or $\alpha_n \geq \zeta$. In the former case we have $\beta_n^l(t) \leq \zeta^l(t) \leq \zeta(t) = x$ since β_n^l is the leftmost characteristic. In the latter case, similarly, we have $\alpha_n^r(t) \geq \zeta^r(t) \geq \zeta(t) = x$.

To prove the reverse inclusion, let $x \in (\underline{U}^r(t), \overline{U}^l(t)) \setminus \bigcup_{n=1}^{\infty} (\alpha_n^r(t), \beta_n^l(t))$. Then there exists $\zeta \in U$ and a characteristic $\zeta(\cdot)$ such that $\zeta(t) = x$. If $\zeta \in K$ then $x \in K(t)$ and the proof is finished. If $\zeta \notin K$ then $\zeta \in (\alpha_n, \beta_n)$ for some n . Since, however, $\zeta(t) \notin (\alpha_n^r(t), \beta_n^l(t))$, there exists time $\tau \in (t_0, t]$ such that either $\zeta(\tau) = \alpha_n^r(\tau)$ (see Fig. 7) or $\zeta(\tau) = \beta_n^l(\tau)$. In the former case we observe that due to (30), $\alpha_n \neq \underline{U}$ and hence $\alpha_n \in K$. The characteristic defined by

$$\alpha_n(s) = \begin{cases} \alpha_n^r(s) & \text{if } s \in [t_0, \tau], \\ \zeta(t) & \text{if } s \in (\tau, t] \end{cases}$$

satisfies $\alpha_n(t_0) \in K$ and $\alpha_n(t) = x$. Hence, $x \in K(t)$. The proof in the case $\zeta(\tau) = \beta_n^l(\tau)$ is analogous. \square

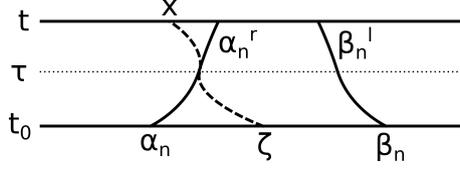


Figure 7: Illustration of the proof of Lemma 4.11. If $\zeta \in (\alpha_n, \beta_n)$ and $x = \zeta(t) \notin (\alpha_n^r(t), \beta_n^l(t))$ then, due to geometrical reasons, $\alpha_n^r(\cdot)$ and $\zeta(\cdot)$ have to cross. Hence, the concatenation of $\alpha_n^r(\cdot)$ (until the crossing time τ) and $\zeta(\cdot)$ (from τ on) is an example of characteristic emanating from K and ending in x .

Proposition 4.12. *Let u be a weak solution of (2)-(4) and let K be an arbitrary compact subset of \mathbb{R} . Fix $t_0 \geq 0$. Then*

$$\mu^+(t_0, K) = \lim_{t \rightarrow t_0^+} \int_{K(t)} (u_x^+)^2 dx - \int_K (u_x^+)^2(t_0, x) dx, \quad (32)$$

where $K(t)$ is the thick pushforward.

Proof. Let $U = (\underline{U}, \bar{U})$ be a bounded open interval such that $K \subset U$ and so large that the assumptions of Lemma 4.11 are satisfied. By Corollary 4.3

$$\mu^+(t_0, U) = \lim_{t \rightarrow t_0^+} \int_{(\underline{U}^r(t), \bar{U}^l(t))} (u_x^+)^2(t, x) dx - \int_U (u_x^+)^2(t_0, x) dx.$$

On the other hand, $U = K \cup A$ where $A = U \setminus K = \bigcup_{n=1}^{\infty} (\alpha_n, \beta_n)$. Hence, by (28)

$$\mu^+(t_0, A) = \lim_{t \rightarrow t_0^+} \int_{\bigcup_{n=1}^{\infty} (\alpha_n^r(t), \beta_n^l(t))} (u_x^+)^2(t, x) dx - \int_A (u_x^+)^2(t_0, x) dx.$$

Subtracting these two equalities and using Lemma 4.11 we conclude. \square

B - arbitrary Borel set

Proof of Theorem 2.11. Fix $\epsilon > 0$ By regularity of the measure μ^+ there exists a compact set K^ϵ and an open set A^ϵ such that

$$K^\epsilon \subset B \subset A^\epsilon$$

and

$$\begin{aligned} \int_{A^\epsilon \setminus B} (u_x^+(t_0, x))^2 dx &< \epsilon, \\ \mu^+(t_0, A^\epsilon \setminus B) &< \epsilon, \\ \int_{B \setminus K^\epsilon} (u_x^+(t_0, x))^2 dx &< \epsilon, \\ \mu^+(t_0, B \setminus K^\epsilon) &< \epsilon. \end{aligned}$$

We calculate

$$\begin{aligned} \limsup_{t \rightarrow t_0^+} \int_{B(t)} (u_x^+)^2(t, x) dx &\leq \lim_{t \rightarrow t_0^+} \int_{A^\epsilon(t)} (u_x^+)^2(t, x) dx \\ &= \mu^+(t_0, A^\epsilon) + \int_{A^\epsilon} (u_x^+)^2(t_0, x) dx \\ &\leq \mu^+(t_0, B) + \epsilon + \int_B (u_x^+)^2(t_0, x) dx + \epsilon, \end{aligned}$$

where we used (27) to pass to the limit with t . Similarly,

$$\begin{aligned} \liminf_{t \rightarrow t_0^+} \int_{B(t)} (u_x^+)^2(t, x) dx &\geq \lim_{t \rightarrow t_0^+} \int_{K^\epsilon(t)} (u_x^+)^2(t, x) dx \\ &= \mu^+(t_0, K^\epsilon) + \int_{K^\epsilon} (u_x^+)^2(t_0, x) dx \\ &\geq \mu^+(t_0, B) - \epsilon + \int_B (u_x^+)^2(t_0, x) dx - \epsilon, \end{aligned}$$

where we used (32) to pass to the limit with t . Passing $\epsilon \rightarrow 0$ we conclude. \square

5 Proofs of Theorems 2.13, 2.14, 2.15 and 2.16

Proof of Theorem 2.13. Decompose

$$\mu^+(t_0, dx) = g\mathcal{L}^1(dx) + (\mu^+)^{sing}(t_0, dx)$$

where g is the density of $(\mu^+(t_0, dx))^{ac}$ with respect to the one-dimensional Lebesgue measure and $(\mu^+)^{sing}$ denotes the singular, with respect to \mathcal{L}^1 , part of measure $\mu^+(t_0, dx)$. Suppose $g \neq 0$. Then there exists a Borel set B and $\epsilon > 0$ such that

$$\begin{aligned} (\mu^+)^{sing}(t_0, B) &= 0, \\ \mathcal{L}^1(B) &\geq \epsilon, \\ g &\geq \epsilon \text{ on } B. \end{aligned}$$

Furthermore, due to Proposition 2.2 there exists a subset $B_1 \subset B$ and $\tau > 0$ such that

- $\mathcal{L}^1(B_1) > \epsilon/2$,
- the characteristics originating in B_1 are unique on $[t_0, t_0 + \tau]$

and for every $\zeta \in B_1$ and $t \in [t_0, t_0 + \tau]$

- $|u_x(t, \zeta(t))| < 1/\tau$, where $\zeta(\cdot)$ is the unique characteristic satisfying $\zeta(t_0) = \zeta$,
- the differential equation

$$\dot{v} = u^2 - \frac{1}{2}v^2 - P,$$

is satisfied, where $v = v(t) = u_x(t, \zeta(t))$, $u = u(t) = u(t, \zeta(t))$, $P = P(t) = P(t, \zeta(t))$.

The last condition implies, in particular, that for every $\zeta \in B_1$ we have

$$\lim_{t \rightarrow t_0^+} u_x^+(t, \zeta(t)) = u_x^+(t_0, \zeta(t_0)).$$

Consequently, using the representation formula from Theorem 2.11 and the change of variables formula from Proposition 3.4 we obtain

$$\begin{aligned} \mu^+(t_0, B_1) &= \lim_{t \rightarrow t_0^+} \int_{B_1(t)} (u_x^+)^2(t, x) dx - \int_{B_1} (u_x^+)^2(t_0, x) dx \\ &= \lim_{t \rightarrow t_0^+} \int_{B_1} (u_x^+)^2(t, \zeta(t)) e^{\int_{t_0}^t u_x(s, \zeta(s)) ds} d\zeta - \int_{B_1} (u_x^+)^2(t_0, \zeta) d\zeta \\ &= 0, \end{aligned}$$

where the last equality follows by the Lebesgue dominated convergence theorem and the bound

$$(u_x^+)^2(t, \zeta(t)) e^{\int_{t_0}^t u_x(s, \zeta(s)) ds} \leq \tau^{-2} e^{\tau/\tau} = \tau^{-2} e,$$

which holds for every $\zeta \in B_1$ and $t \in [t_0, t_0 + \tau]$. On the other hand,

$$\mu^+(t_0, B_1) \geq \int_{B_1} g(x) dx \geq \epsilon^2/2,$$

which gives contradiction and concludes the proof for μ^+ . The proof for μ^- follows dually by considering the 'backward' solution $u^{t_0 b}(t, x) := -u(t_0 - t, x)$ and using the result for μ^+ . \square

Proof of Theorem 2.14. We distinguish two cases, $t_0 = 0$ and $t_0 > 0$ as the proofs in those cases are different. Both proceed by obtaining a contradiction if we assume that $\mu^+(t_0, \cdot) \neq 0$, however in the former we obtain a contradiction with the weak energy condition and in the latter with the Oleinik-type criterion from Definition 1.2.

Case $t_0 = 0$. Suppose $\mu^+(t_0, \mathbb{R}) = S > 0$. For a test function $\phi \in C_c(\mathbb{R})$ we calculate, using Theorem 2.4 and Definition 2.5,

$$\begin{aligned} \lim_{t \rightarrow t_0^+} \int_{\mathbb{R}} \phi(x)(u_x)^2(t, x) dx &= \lim_{t \rightarrow t_0^+} \int_{\mathbb{R}} \phi(x)(u_x^+)^2(t, x) dx + \lim_{t \rightarrow t_0^+} \int_{\mathbb{R}} \phi(x)(u_x^-)^2(t, x) dx \\ &= \int_{\mathbb{R}} \phi(x)(u_x)^2(t_0, x) dx + \int_{\mathbb{R}} \phi(x) \mu^+(t_0, dx). \end{aligned}$$

Take $\phi = \phi^K$ of the following form:

$$\phi^K(x) := \begin{cases} 0 & \text{if } x \in (-\infty, -K-1), \\ x + K + 1 & \text{if } x \in [-K-1, -K), \\ 1 & \text{if } x \in [-K, K), \\ K + 1 - x & \text{if } x \in [K, K+1), \\ 0 & \text{if } x \in [K+1, \infty), \end{cases}$$

where K is so big that

$$\begin{aligned} \int_{\mathbb{R}} \phi(x) \mu^+(t_0, dx) &\geq S - S/4, \\ \int_{\mathbb{R}} \phi(x)(u_x)^2(t_0, x) dx &\geq \int_{\mathbb{R}} (u_x)^2(t_0, x) dx - S/4, \\ \int_{\mathbb{R}} \phi(x) u^2(t_0, x) dx &\geq \int_{\mathbb{R}} u^2(t_0, x) dx - S/4. \end{aligned}$$

Then

$$\liminf_{t \rightarrow t_0^+} \int_{\mathbb{R}} (u_x)^2(t, x) dx \geq \lim_{t \rightarrow t_0^+} \int_{\mathbb{R}} \phi(x)(u_x)^2(t, x) dx \geq \int_{\mathbb{R}} (u_x)^2(t_0, x) dx - S/4 + (S - S/4)$$

and

$$\liminf_{t \rightarrow t_0^+} \int_{\mathbb{R}} u^2(t, x) dx \geq \lim_{t \rightarrow t_0^+} \int_{\mathbb{R}} \phi(x) u^2(t, x) dx = \int_{\mathbb{R}} \phi(x) u^2(t_0, x) dx \geq \int_{\mathbb{R}} u^2(t_0, x) dx - S/4.$$

Combination of these two convergences leads to

$$\liminf_{t \rightarrow t_0^+} \int_{\mathbb{R}} [(u_x)^2(t, x) + u^2(t, x)] dx \geq \int_{\mathbb{R}} [(u_x)^2(t_0, x) + u^2(t_0, x)] dx + S/4,$$

which contradicts the weak energy condition from Definition 1.2 and hence implies that u is not dissipative.

Case $t_0 > 0$. Since $\mu^+(t_0, x) \perp \mathcal{L}^1$ by Theorem 2.13, there exists a set C (without loss of generality bounded) satisfying

- $\mathcal{L}^1(C) = 0$,
- $\mu^+(t_0, \mathbb{R} \setminus C) = 0$.

Suppose $\mu^+(t_0, C) = S > 0$ and fix $\epsilon > 0$. By regularity of the measure \mathcal{L}^1 there exists an open set $U = \bigcup_{n=1}^{\infty} (\alpha_n, \beta_n)$ such that

- $C \subset U$,
- $\mathcal{L}^1(U) < \epsilon$,
- (α_n, β_n) are pairwise disjoint.

Let N be so large that

$$\mu^+ \left(t_0, \bigcup_{n=1}^N (\alpha_n, \beta_n) \right) > S/2.$$

Next, take δ so small that for $U_N^\delta := \bigcup_{n=1}^N (\alpha_n + \delta, \beta_n - \delta)$ we have

$$\mu^+(t_0, U_N^\delta) > S/4.$$

Taking

$$\phi^\epsilon(x) := \begin{cases} 1 & \text{if } x \in U_N^\delta, \\ (x - \alpha_n)/\delta & \text{if } x \in (\alpha_n, \alpha_n + \delta], \\ (\beta_n - x)/\delta & \text{if } x \in [\beta_n - \delta, \beta_n), \\ 0 & \text{otherwise} \end{cases}$$

we obtain

$$\begin{aligned} \lim_{t \rightarrow t_0^+} \int_{\mathbb{R}} (u_x^+)^2(t_0, x) \phi^\epsilon(x) dx &= \int_{\mathbb{R}} \phi^\epsilon(x) d\mu^+(t_0, x) + \int_{\mathbb{R}} (u_x^+)^2(t_0, x) \phi^\epsilon(x) dx \\ &\geq S/4 + \int_{\mathbb{R}} (u_x^+)^2(t_0, x) \phi^\epsilon(x) dx \geq S/4. \end{aligned} \quad (33)$$

Noting that $0 \leq \phi^\epsilon(x) \leq 1$ for every $x \in \mathbb{R}$ and $\mathcal{L}^1(\{x \in \mathbb{R} : \phi^\epsilon(x) \neq 0\}) < \epsilon$ we pass to the limit $\epsilon \rightarrow 0$ in (33) and obtain that u_x^+ is unbounded in the neighborhood of $\{t_0\} \times C$, which contradicts dissipativity. \square

Proof of Theorem 2.15. The function

$$t \mapsto \int_{[\alpha^l(t), \beta^r(t)]} (u_x^+)^2(t, x) dx$$

belongs to BV_{loc} by [26, Theorem 2.7]. Hence, it has only at most countably many discontinuities and thus $\mu^+([\alpha^l(t), \beta^r(t)]) \neq 0$ for at most countably many t . Considering $\alpha_n := -n$ and $\beta_n := n$ we obtain, using the $\sup(|u|)$ bound on the speed of characteristics, that $\mu^+([-n + \sup(u)t, n - \sup(u)t]) \neq 0$ for at most countably many t . Since this holds for every $n = 1, 2, \dots$, we obtain $\mu^+(t, \mathbb{R}) \neq 0$ for at most countably many t . \square

Proof of Theorem 2.16. By the theory of Bressan-Constantin [2] there exists a 'conservative' solution \hat{u} of (2)-(4), defined on $[t_0, \infty) \times \mathbb{R}$, which satisfies:

- $\hat{u}(t_0, \cdot) = u(t_0, \cdot)$,
- $E(\hat{u}(t, \cdot)) \leq E(\hat{u}(t_0, \cdot))$ for every $t \geq t_0$,
- $E(\hat{u}(t, \cdot)) = E(\hat{u}(t_0, \cdot))$ for almost every $t \geq t_0$.

Let \bar{u} be defined by

$$\bar{u}(t, \cdot) = \begin{cases} u(t, \cdot) & \text{for } t \in [0, t_0], \\ \hat{u}(t, \cdot) & \text{for } t \in (t_0, \infty). \end{cases}$$

and let $D > 0$ be so large that for $B := [-D, D]$ we have

$$\int_B \frac{1}{2}(u^2(t_0, x) + u_x^2(t_0, x))dx \geq E(u(t_0, \cdot)) - \frac{\mu^+(t_0, B)}{4}.$$

Then, observing that

$$\begin{aligned} \lim_{t \rightarrow t_0^+} \int_{B(t)} u^2(t, x)dx &= \int_B u^2(t_0, x)dx, \\ \lim_{t \rightarrow t_0^+} \int_{B(t)} (u_x^-)^2(t, x)dx &= \int_B (u_x^-)^2(t_0, x)dx, \\ \lim_{t \rightarrow t_0^+} \int_{B(t)} (u_x^+)^2(t, x)dx &= \int_B (u_x^+)^2(t_0, x)dx + \mu^+(t_0, B), \end{aligned}$$

where the first convergence follows by continuity of u , the second by [26, Theorem 2.6] and the third by Theorem 2.11, we obtain

$$\begin{aligned} \limsup_{t \rightarrow t_0^+} E(\bar{u}(t, \cdot)) &= \limsup_{t \rightarrow t_0^+} E(\hat{u}(t, \cdot)) \leq E(u(t_0, \cdot)) \\ &\leq \int_B \frac{1}{2}(u^2(t_0, x) + u_x^2(t_0, x))dx + \frac{\mu^+(t_0, B)}{4} \\ &= \lim_{t \rightarrow t_0^+} \int_{B(t)} \frac{1}{2}(u^2(t, x) + u_x^2(t, x))dx - \frac{\mu^+(t_0, B)}{2} + \frac{\mu^+(t_0, B)}{4} \\ &\leq \liminf_{t \rightarrow t_0^+} E(u(t, \cdot)) - \frac{\mu^+(t_0, B)}{4}. \end{aligned}$$

□

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