

the previous sections can be expressed by

$$\begin{aligned}
g_{00}^0(\Gamma_{15}) &= \mathcal{J}[e_x^2(-)], \\
g_{01}^0(\Gamma_{15}) &= \sqrt{2} \mathcal{J}[e_x(-)e_x(+) \cos 2\pi r_0 q_x], \\
g_{11}^{0\pm}(\Gamma_{15}) &= 2 \mathcal{J}[e_x^2(\pm) \cos^2 2\pi r_0 q_x], \\
g_{11}^0(\Gamma_1) &= 2 \mathcal{J}[e_x^2(+) \sin^2 2\pi r_0 q_x + 2e_x(+)e_y(+) \sin 2\pi r_0 q_x \sin 2\pi r_0 q_y], \\
g_{11}^0(\Gamma_{12}) &= 2 \mathcal{J}[e_x^2(+) \sin^2 2\pi r_0 q_x - e_x(+)e_y(+) \sin 2\pi r_0 q_x \sin 2\pi r_0 q_y], \\
g_{14}^0(\Gamma_1) &\cong g_{14}^0(\Gamma_{12}) \cong g_{01}^0(\Gamma_{15})/\sqrt{2}, \\
g_{44}^0(\Gamma_1) &\cong g_{44}^0(\Gamma_{12}) \cong g_{00}^0(\Gamma_{15}), \\
g_{12}^0(\Gamma_1) &= \frac{2}{3} \mathcal{J}[\{e_x(+) \sin 2\pi r_0 q_x + \text{c.p.}\} \{e_x(-) \sin 2\pi r_0 q_x (\cos 2\pi r_0 q_x + \cos 2\pi r_0 q_z) + \text{c.p.}\}], \\
g_{22}^0(\Gamma_1) &= \frac{2}{3} \mathcal{J}[\{e_x(-) \sin 2\pi r_0 q_x (\cos 2\pi r_0 q_x + \cos 2\pi r_0 q_z) + \text{c.p.}\}^2], \\
g_{12}^0(\Gamma_{12}) &= \frac{1}{3} \mathcal{J}[\{2e_x(+) \sin 2\pi r_0 q_x - e_y(+) \sin 2\pi r_0 q_y - e_z(+) \sin 2\pi r_0 q_z\} \\
&\quad \times \{2[e_y(-) \sin 2\pi r_0 q_y - e_x(-) \sin 2\pi r_0 q_x] \cos 2\pi r_0 q_z - [e_z(-) \sin 2\pi r_0 q_z - e_y(-) \sin 2\pi r_0 q_y] \cos 2\pi r_0 q_x \\
&\quad - [e_x(-) \sin 2\pi r_0 q_x - e_z(-) \sin 2\pi r_0 q_z] \cos 2\pi r_0 q_y\}], \quad (\text{A2}) \\
g_{22}^0(\Gamma_{12}) &= \frac{1}{3} \mathcal{J}[\{2[e_y(-) \sin 2\pi r_0 q_y - e_y(-) \sin 2\pi r_0 q_x] \cos 2\pi r_0 q_z \\
&\quad - [e_x(-) \sin 2\pi r_0 q_x - e_y(-) \sin 2\pi r_0 q_y] \cos 2\pi r_0 q_z - [e_x(-) \sin 2\pi r_0 q_x - e_z(-) \sin 2\pi r_0 q_z] \cos 2\pi r_0 q_x\}^2], \\
g_{33}^0(\Gamma_{25'}) &= 2 \mathcal{J}[e_x^2(+) \sin^2 2\pi r_0 q_y + 2e_y(-)e_x(+) \sin 2\pi r_0 q_x \sin 2\pi r_0 q_y], \\
g_{23}^0(\Gamma_{25'}) &= \sqrt{2} \mathcal{J}[\{e_y(+) \sin 2\pi r_0 q_x + e_z(+) \sin 2\pi r_0 q_y\} \{e_y(-) \sin 2\pi r_0 q_x \cos 2\pi r_0 q_z + e_z(-) \sin 2\pi r_0 q_z \cos 2\pi r_0 q_y\}], \\
g_{22}^0(\Gamma_{25'}) &= 2 \mathcal{J}[\{e_y(-) \sin 2\pi r_0 q_y \cos 2\pi r_0 q_z + e_z(-) \sin 2\pi r_0 q_z \cos 2\pi r_0 q_y\}^2].
\end{aligned}$$

In the above expressions $\mathbf{e}(\pm) \equiv \mathbf{e}(\pm | \mathbf{q}j)$ are the polarization vectors for wave vector $\mathbf{q}$ and branch index j , and $\mathcal{J}$ denotes the complex transform

$$\mathcal{J}[f] = \sum_j \int d\mathbf{q} \frac{f}{\omega_{\mathbf{q}j}^2 - z}, \quad (\text{A3})$$

where $\omega_{\mathbf{q}j}$ are the phonon frequencies, and $z = \omega^2 + i0^+$.

Errata

Energy and Orientation Dependence of Electron-Irradiation-Induced Damage in Undoped GaSb, K. THOMMEN [Phys. Rev. **174**, 938 (1968)]. The ordinate notations of the lower left and right diagrams in Fig. 5 should read $(\eta_3^-)^{1/2} \times 10^{12}$ and $(\eta_4^-)^{1/2} \times 10^{12}$, respectively.

Atomic Migration in Monatomic Crystals, C. P. FLYNN [Phys. Rev. **171**, 682 (1968)]. A conversion

error has been detected in the calculation of E_M for bcc metals. The corrected values in eV for $\delta^2 = 0.67$ are as follows:

	Li	Na	K	Cr	V	Fe	Mo	Ta	Nb	W
E_M	0.05	0.06	0.07	1.5	0.9	1.0	2.2	1.4	0.9	2.8

The values for alkali metals remain small compared to Q , while the values for transition metals are increased.